



An exploration of strongly interacting dark sectors in ATLAS and beyond

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RAL PPD seminar
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1824

The University of Manchester



Who am I?

Member of ATLAS
experiment



PhD from University of Witwatersrand, South Africa

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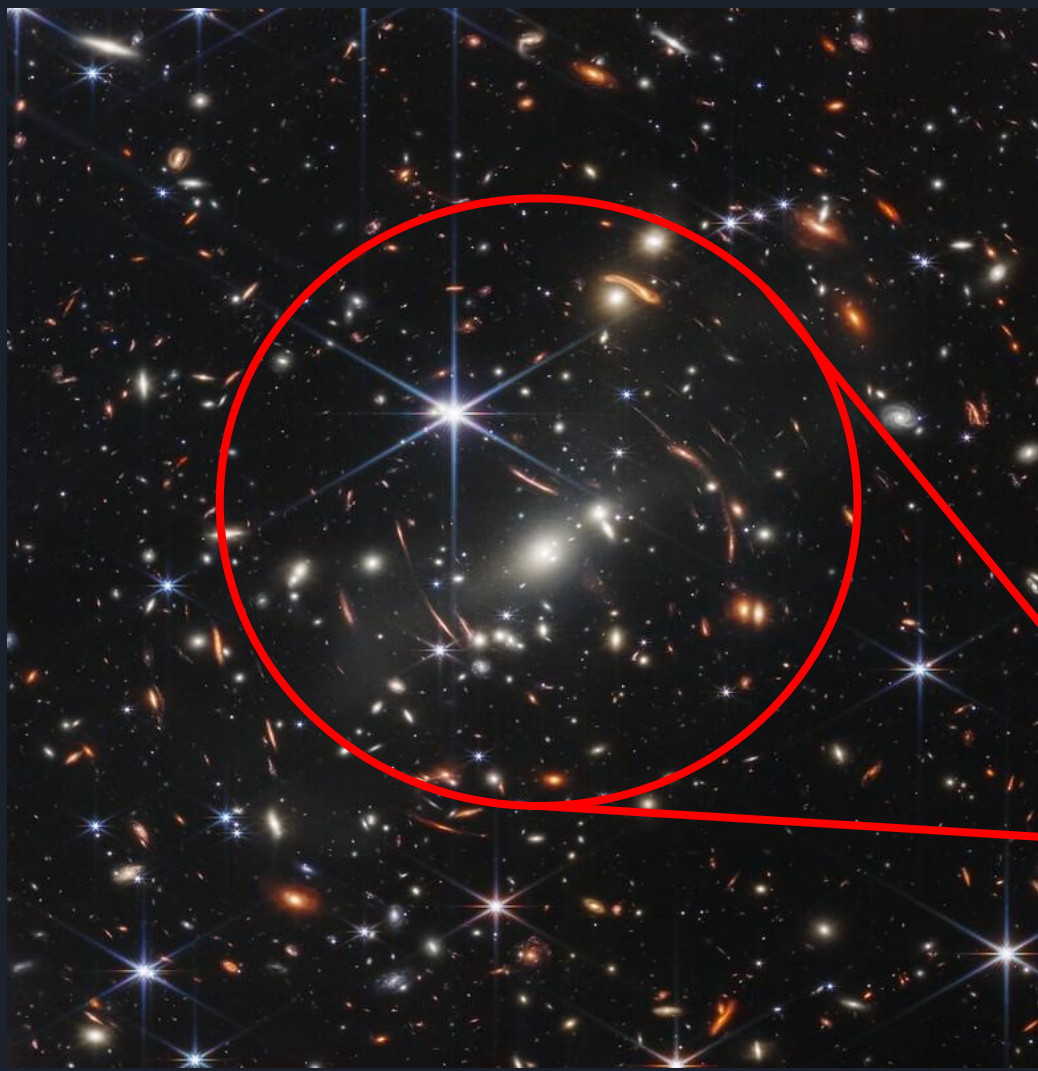
Marie Curie Early Stage researcher
@ Glasgow University, Scotland





First image from the brand new James Webb Space Telescope, designed by NASA - 11th July 2022

Produced the deepest and sharpest infrared image of the distant universe to date!



First image from the brand new James Webb Space Telescope, designed by NASA - 11th July 2022

Produce sharp the d

Why does the light from distant galaxies look all squiggly and bending?



First image from the brand new James Webb Space Telescope, designed by NASA – 11th July 2022

Produces sharp images of the distant universe

Why does the light from distant galaxies look all squiggly and bending?

Gravitational lensing occurs when a massive celestial body — such as a galaxy cluster — causes a sufficient curvature of spacetime for the path of light around it to be visibly bent, as if by a lens

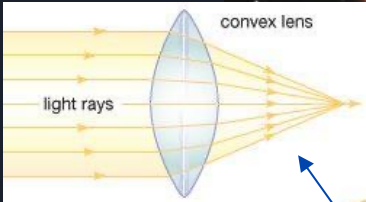
Gravitational lensing!



First image from the brand new James Webb Space Telescope, designed by NASA – 11th July 2022

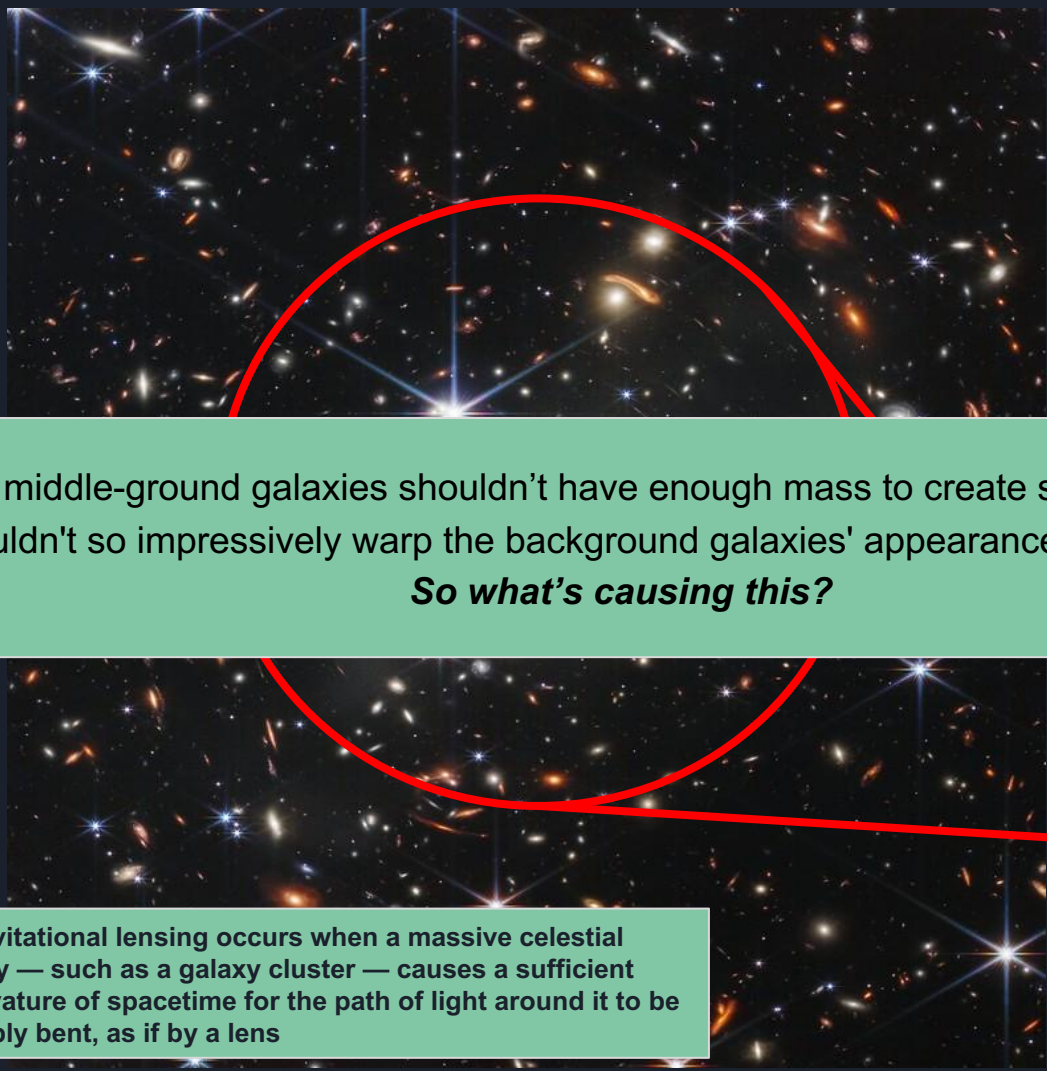
Produces sharp images of the distant universe

Why does the light from distant galaxies look all squiggly and bending?



Gravitational lensing occurs when a massive celestial body — such as a galaxy cluster — causes a sufficient curvature of spacetime for the path of light around it to be visibly bent, as if by a lens

Gravitational lensing!



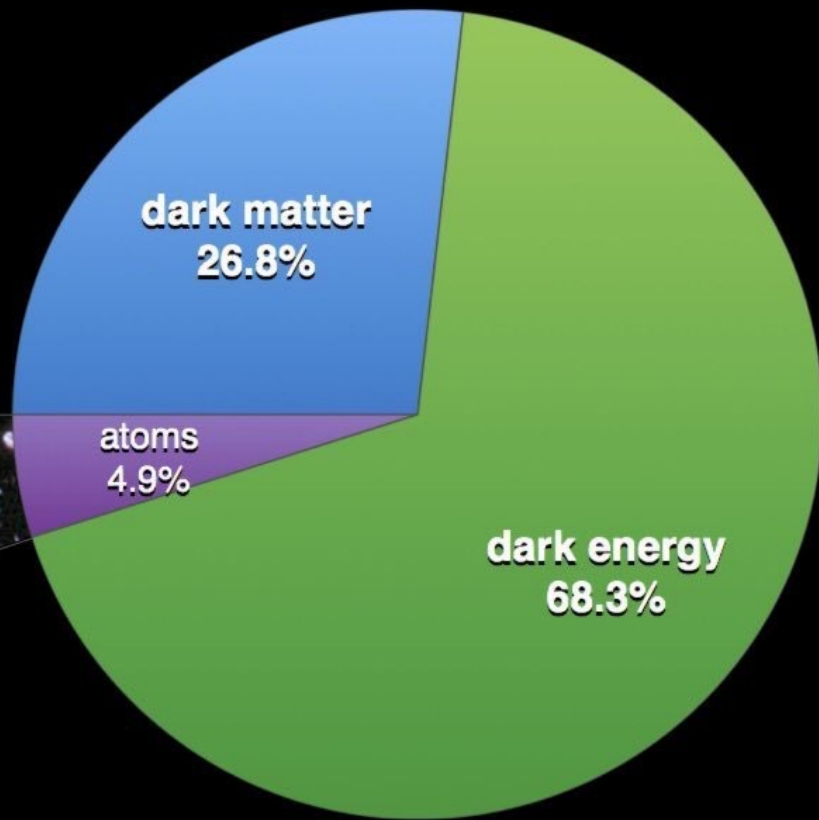
First image from the brand new James Webb Space Telescope, designed by NASA – 11th July 2022

The middle-ground galaxies shouldn't have enough mass to create such a powerful lens, and shouldn't so impressively warp the background galaxies' appearances.
So what's causing this?

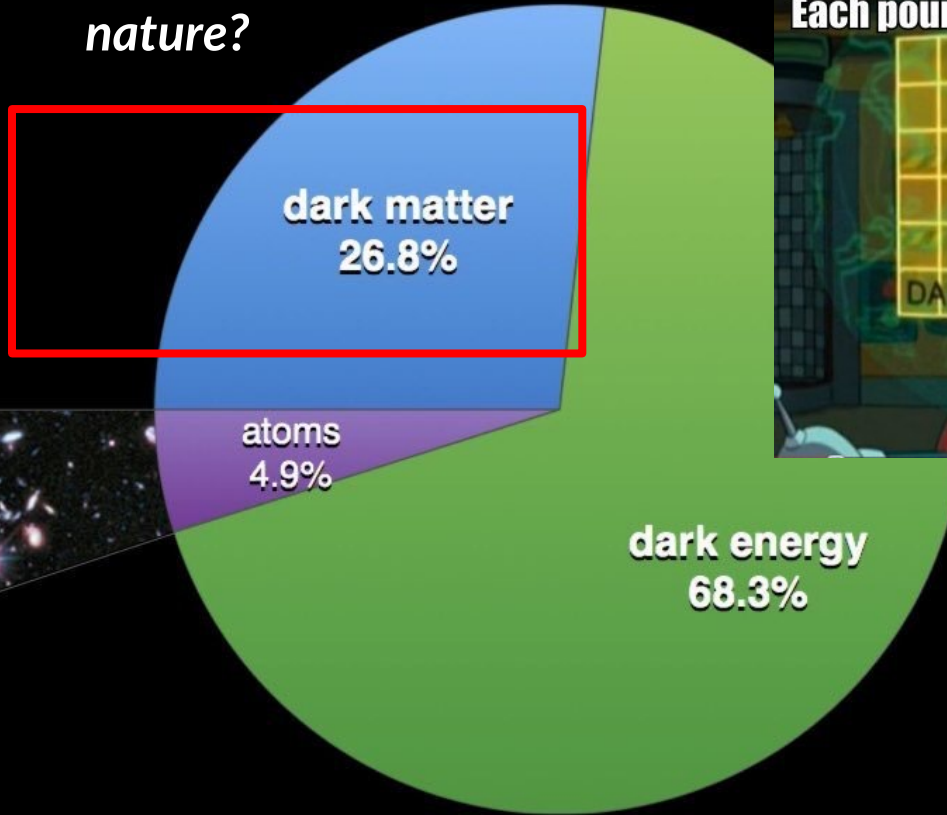
from
look all
squiggly and bending?

Gravitational lensing occurs when a massive celestial body — such as a galaxy cluster — causes a sufficient curvature of spacetime for the path of light around it to be visibly bent, as if by a lens

Gravitational lensing!



*What is its
nature?*

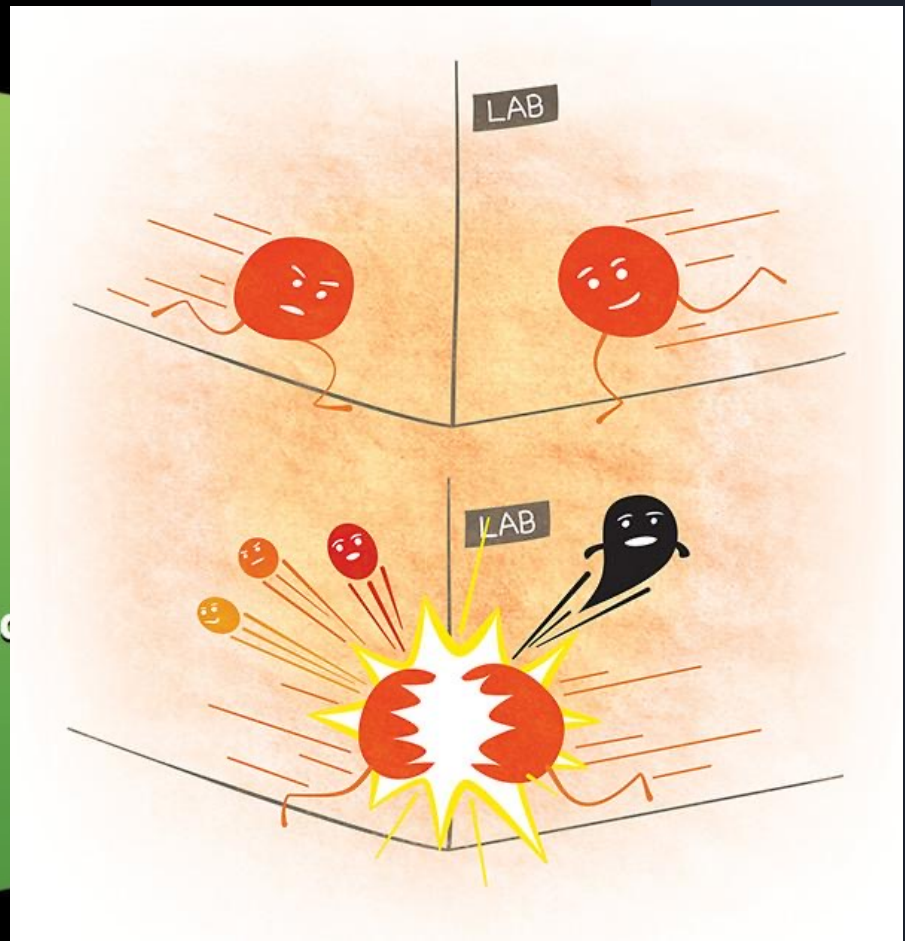
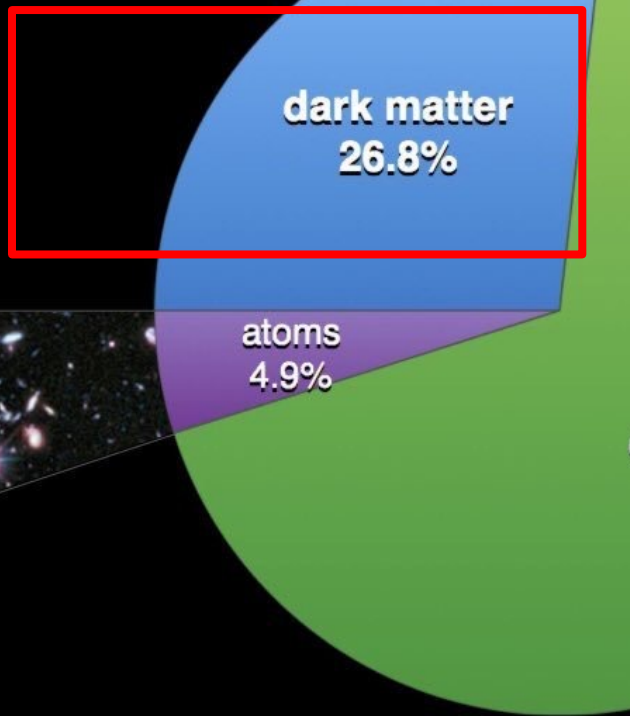


Dark Matter

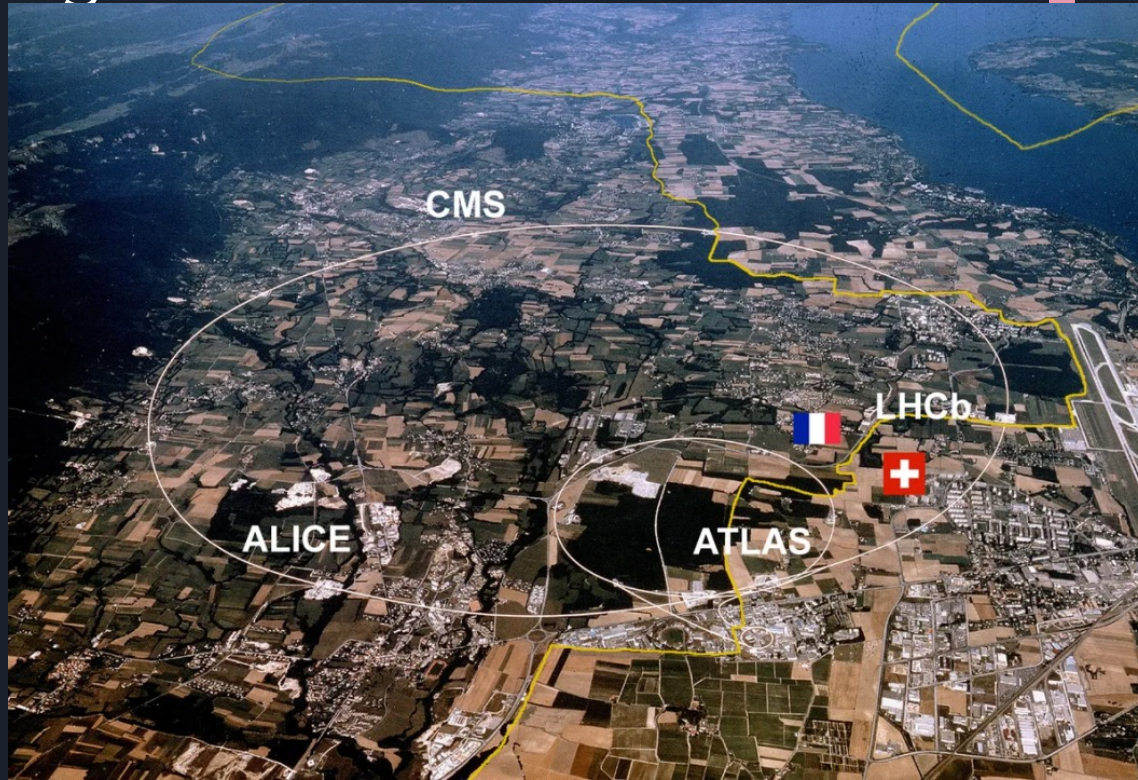
Each pound weighs over 10 thousand pounds!



What is its nature?

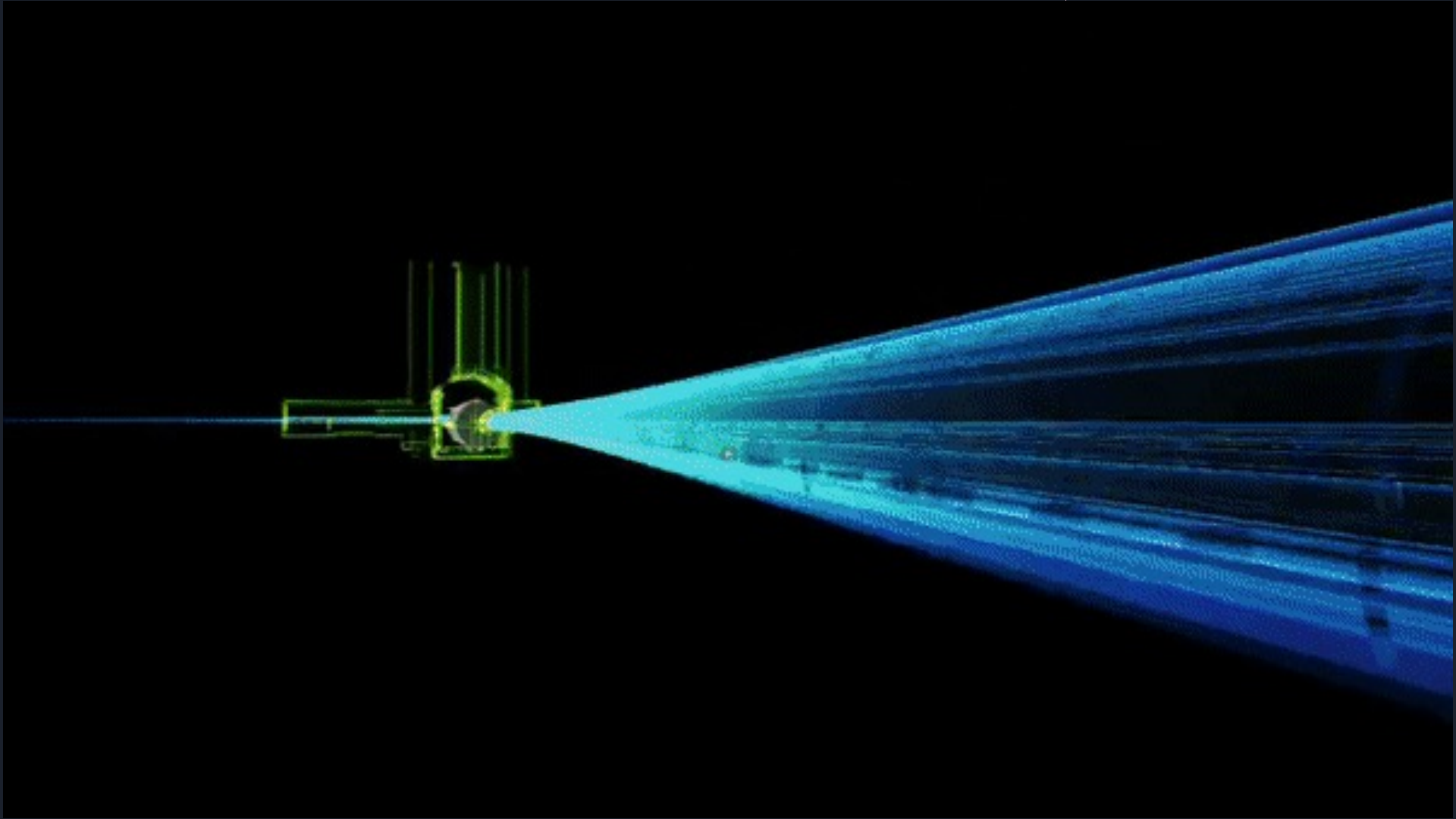


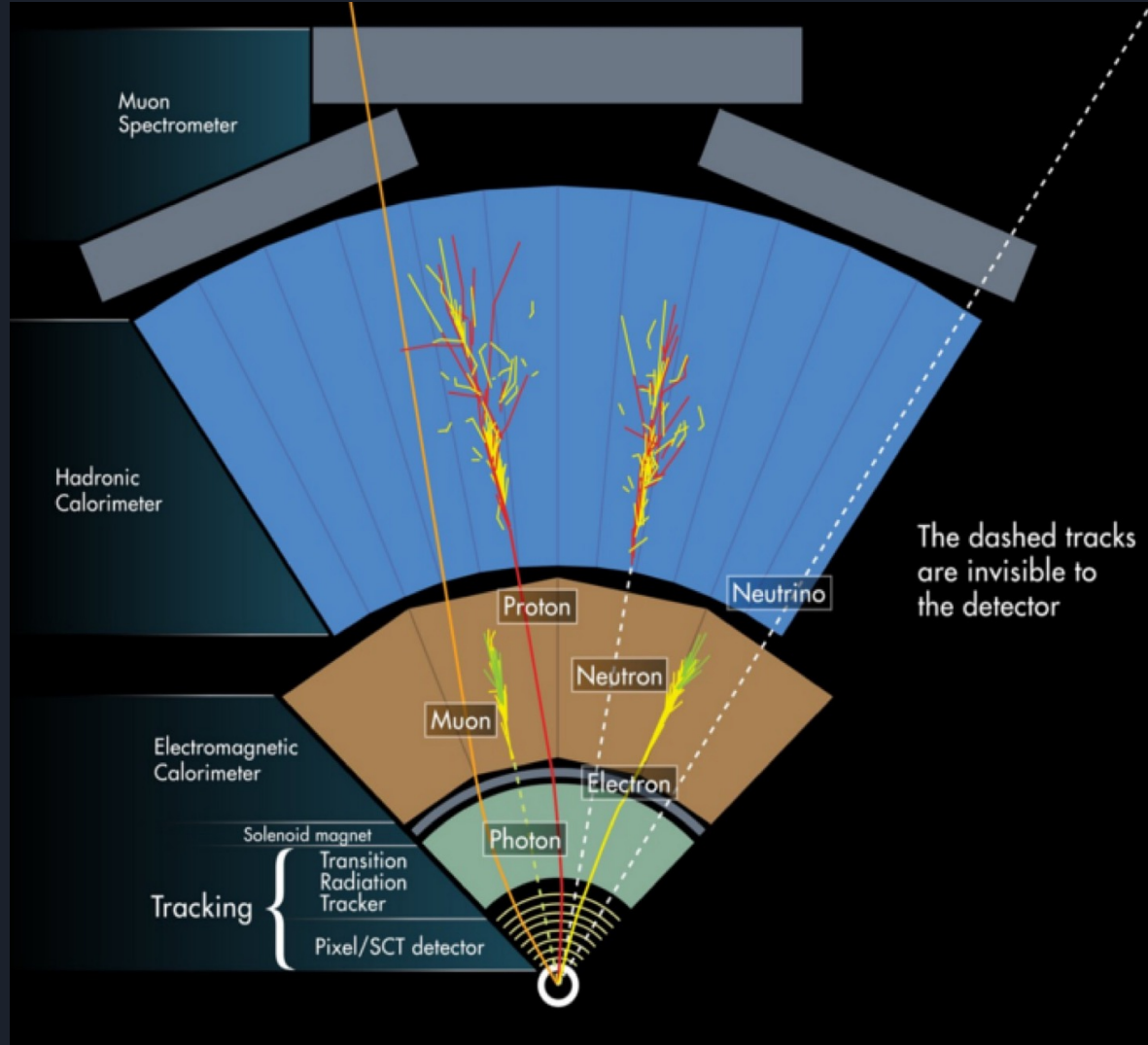
... at the Large Hadron Collider



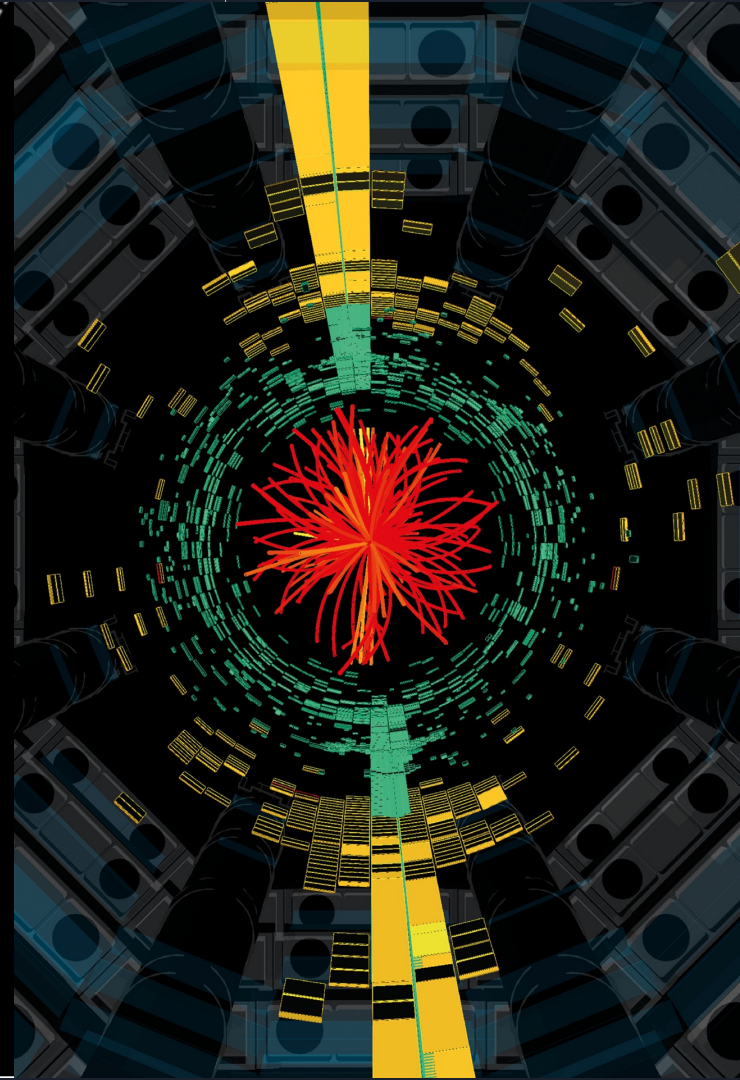
World's largest and most powerful particle accelerator.

Circumference of 27 km, running underground across the borders of Switzerland and France.





The dashed tracks are invisible to the detector



Hidden Sectors

SM Sector

Connectors
/
Portals

Hidden
Sectors

Z' , SUSY particles,
Higgs, Extra Dim,
Leptoquarks,
CP-odd...

can be strongly
or weakly coupled
i.e., dark Higgs,
dark photon,
dark SU(N),
Asym DM...



When a Hidden Sector particle is (quasi-)stable, a dark matter candidate can potentially exist

We have not found any concrete sign of new physics ... yet!

Looking at unusual topologies and hidden corners of the phase space
→ signature based searches, using benchmark models.

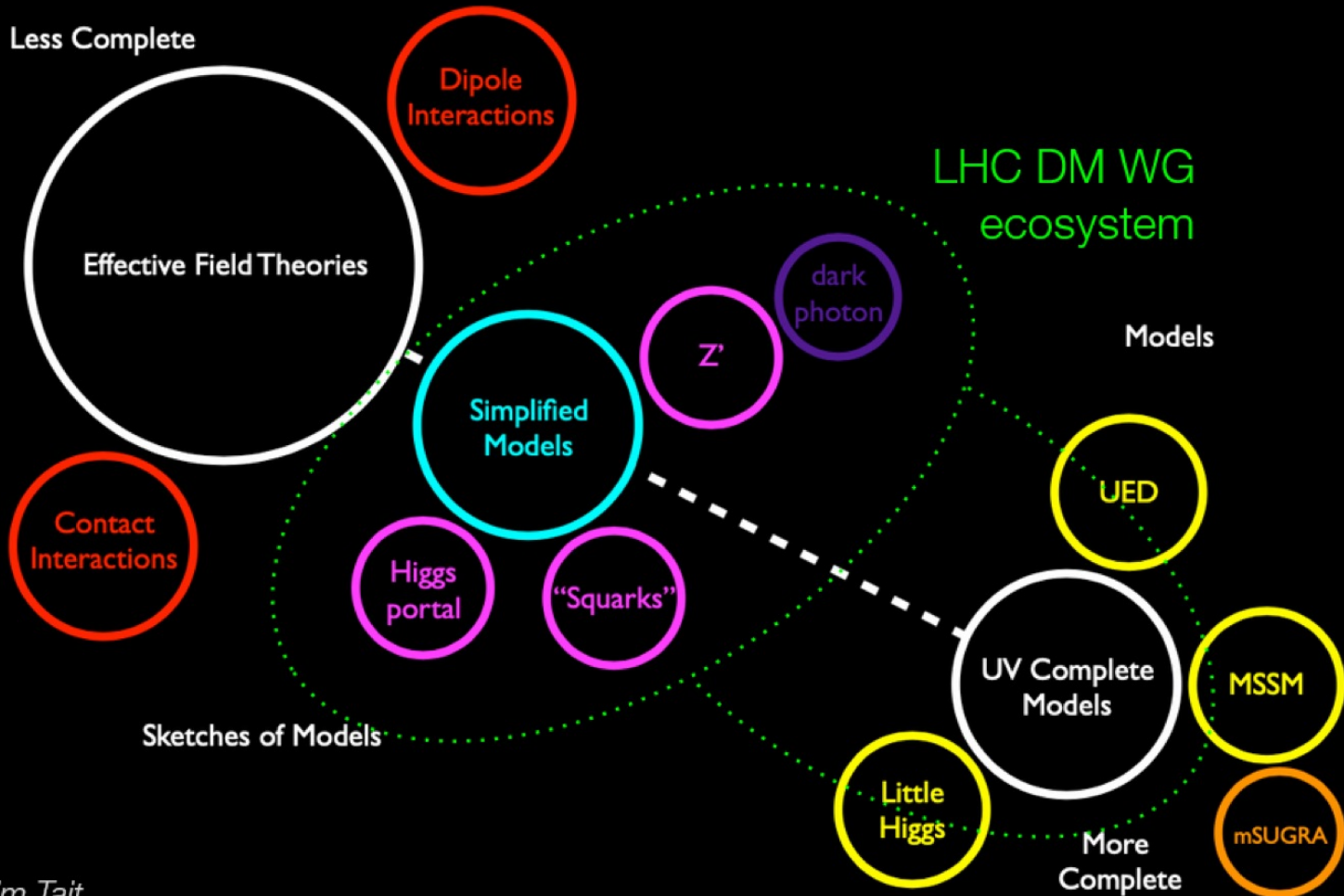
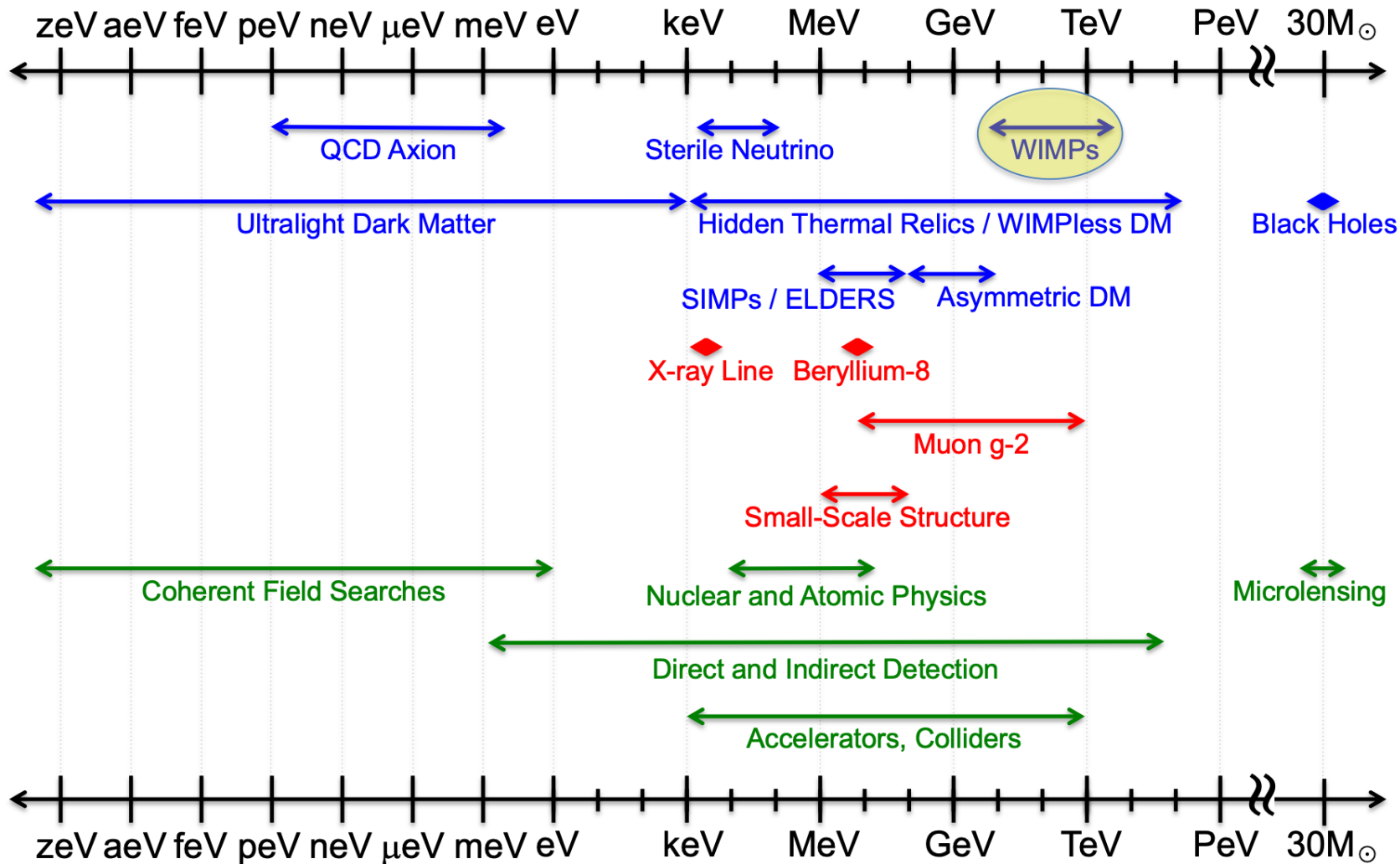


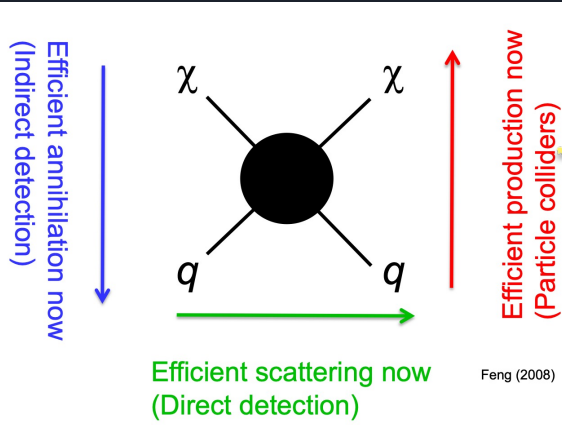
Figure: Tim Tait



WIMPs

The miracle... WIMPs motivated by

- cosmology (production mechanism of thermal freeze-out, expected to have right relic density)
- particle theory (i.e. present in many BSM models)
- particle experiment (accessible in current and near-future energy scales)

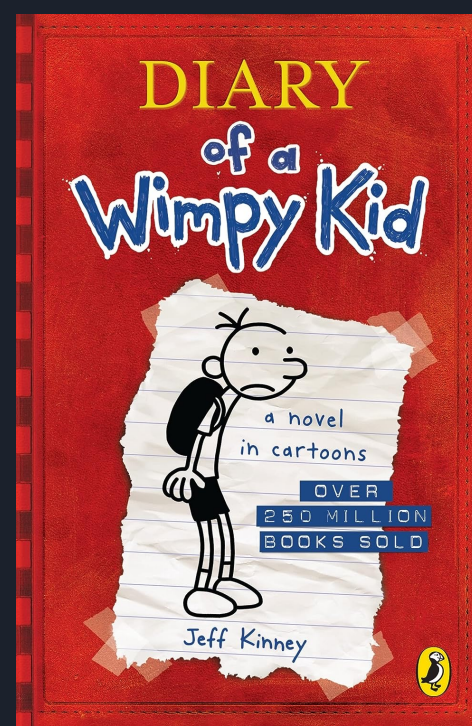


Complementarity of various WIMP dark matter detection methods

Direct WIMP production of $\chi\chi$ pairs is invisible
→ must look for signatures of WIMPs produced in conjunction with other particles.

If SUSY: pair of squarks/gluinos → neutralino WIMP (i.e. MET)
(not discussed in detail)

Simplified models: DM + few other particles → few defining parameters



Over the years, hundreds of dark matter models and search strategies have been proposed....



ATLAS Collaboration, G. Aad et al., *Search for dark matter in events with a hadronically decaying W or Z boson and missing transverse momentum in pp collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector*, *Phys.Rev.Lett.* **112** (2014), no. 4 041802, [[arXiv:1309.4017](https://arxiv.org/abs/1309.4017)].

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ATLAS Collaboration, G. Aad et al., *Search for dark matter in events with a hadronically decaying W or Z boson and missing transverse momentum in pp collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector*, *Phys.Rev.Lett.* **112** (2014), no. 4 041802, [[arXiv:1309.4017](#)].

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ATLAS Collaboration, G. Aad et al., *Search for new phenomena in events with a photon and missing transverse momentum in pp collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector*, *Phys. Rev. D* **91** (2015), no. 1 012008, [[arXiv:1411.1559](#)]. [Erratum: *Phys. Rev. D* **92**, no. 5, 059903 (2015)].

ATLAS Collaboration, G. Aad et al., *Search for dark matter in events with a Z boson and missing transverse momentum in pp collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector*, *Phys. Rev. D* **90** (2014) 012004, [[arXiv:1404.0051](#)].

ATLAS Collaboration, G. Aad et al., *Search for dark matter in events with a hadronically decaying W or Z boson and missing transverse momentum in pp collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector*, *Phys. Rev. Lett.* **112** (2014), no. 4 041802, [[arXiv:1309.4017](#)].

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CMS Collaboration, V. Khachatryan et al. *Search for new phenomena with top quark pairs in final states with one lepton, unparticles in monojet events in proton-proton collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector* (2015), no. 5 235, [[arXiv:1408.3583](https://arxiv.org/abs/1408.3583)]

Search for new phenomena with top quark pairs in final states with one lepton, jets, and missing transverse momentum in pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector

ATLAS Collaboration, G. Aad et al., *Search for direct pair production of supersymmetric partners to the τ lepton in proton-proton collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector*

ATLAS uses the Higgs boson as a tool to search for Dark Matter

29th October 2020 | By [ATLAS Collaboration](#)

Phys. Rev. D **90** (2014) 012004 [[arXiv:1307.3536](https://arxiv.org/abs/1307.3536)]
ATLAS Collaboration, G. Aad et al., *Search for*

hadronic tt plus missing transverse momentum final state at $\sqrt{s} = 13$ TeV with the ATLAS detector

energetic jets and large missing transverse momentum in pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector
Search for direct pair production of supersymmetric partners to the τ lepton in proton-proton collisions at $\sqrt{s} = 13$ TeV

nically
} TeV
.7].

with the CMS Collaboration • [Albert M Sirunyan](#) (Yerevan Phys. Inst.) et al. (Jul 30, 2019)

Published in: *Eur.Phys.J.C* 80 (2020) 3, 189 • e-Print: [1907.13179](https://arxiv.org/abs/1907.13179) [hep-ex]

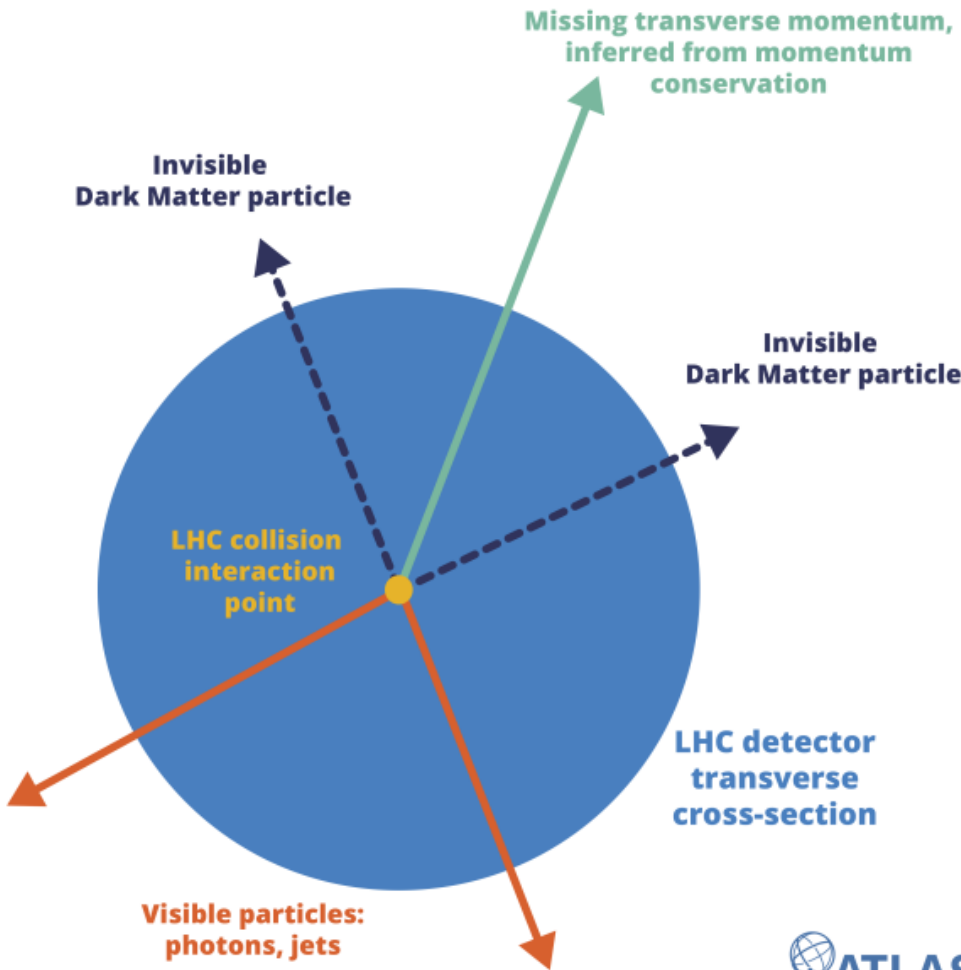
Over the years, hundreds of experiments have been proposed....

CMS Collaboration, V. Khachatryan et al. *Search for invisible particles in monojet events* *Phys. Rev. Lett.* **C75** (2015), no. 5 235, [arXiv:1406.0064].
ATLAS Collaboration, G. Aad et al. *Search for invisible particles in monojet events* *Phys. Rev. Lett.* **C75** (2015), no. 5 235, [arXiv:1406.0064].

ATLAS uses Missing Energy to Search for Dark Matter

29th October 2020 | By ATLAS

Phys. Rev. Lett. **D90** (2013) 151802
ATLAS Collaboration, G. Aad et al. *Search for direct production of dark matter particles in proton-proton collisions with the ATLAS detector* *Phys. Rev. Lett.* **D107** (2011) 251802
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Published in: *Eur. Phys. J.*



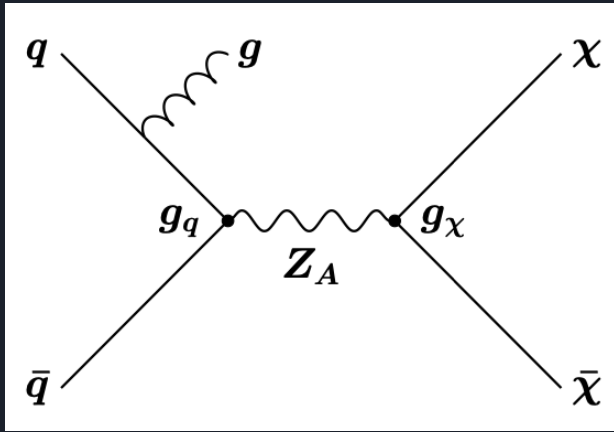
Final states with one lepton, ...
is at $\sqrt{s} = 13$ TeV with the

Search for Dark Matter

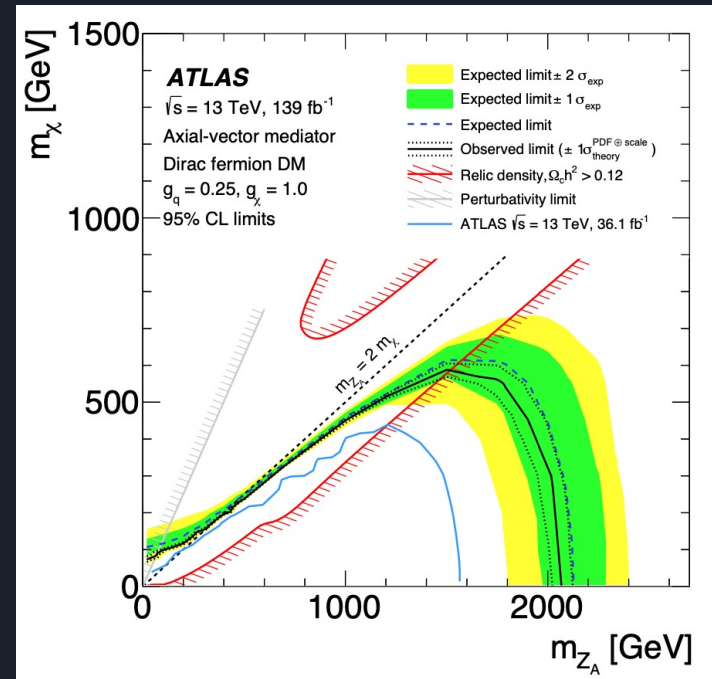
the momentum that ...
ATLAS detector

in proton- ...
nically ...
} TeV ...
[7].

WIMPs



Mono-jet: WIMP pair production with ISR gluon

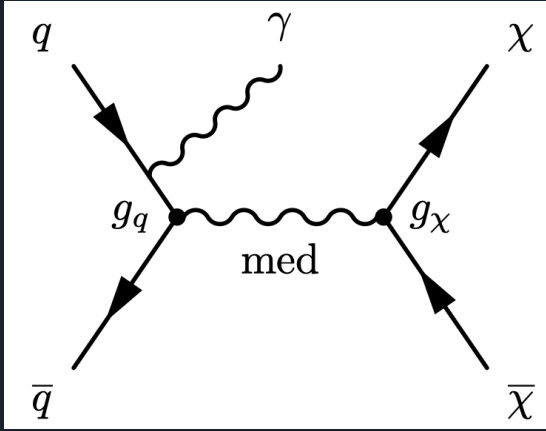


In the region $m_{Z_A} > 2m_\chi$, mediator masses up to about 2.1 TeV are excluded for $m_\chi = 1$ GeV

Masses corresponding to relic density determined by the Planck and WMAP satellites
 → line that crosses the excluded region at $m_{Z_A} \sim 1500$ GeV and $m_\chi \sim 585$ GeV

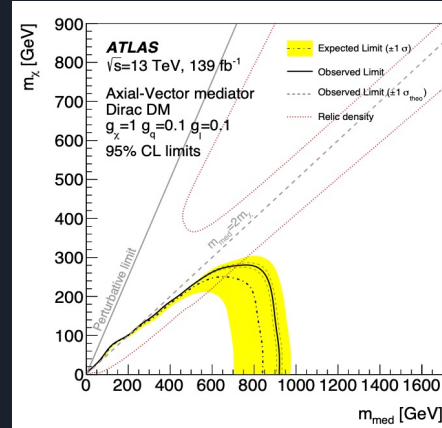
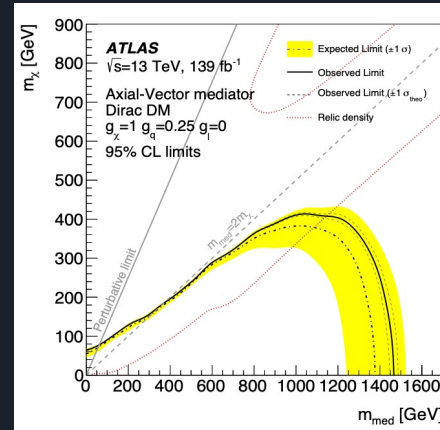
Exclusive (EM)	EM0	EM1	EM2	EM3	EM4	EM5	EM6
E_T^{miss} [GeV]	200–250	250–300	300–350	350–400	400–500	500–600	600–700
	EM7	EM8	EM9	EM10	EM11	EM12	
	700–800	800–900	900–1000	1000–1100	1100–1200	> 1200	
Inclusive (IM)	IM0	IM1	IM2	IM3	IM4	IM5	IM6
E_T^{miss} [GeV]	> 200	> 250	> 300	> 350	> 400	> 500	> 600
	IM7	IM8	IM9	IM10	IM11	IM12	
	> 700	> 800	> 900	> 1000	> 1100	> 1200	

WIMPs



Mono-photon: WIMP pair production with ISR photon

E_T^{miss} [GeV]	SRI1	SRI2	SRI3	SRI4	SRE1	SRE2	SRE3
	> 200	> 250	> 300	> 375	200–250	250–300	300–375



JHEP 02 (2021) 226

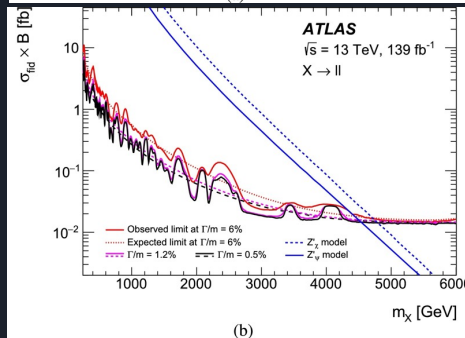
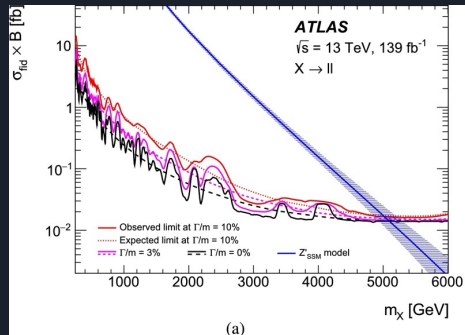
Masses corresponding to relic density determined by the Planck and WMAP satellites

Mediator	g_q	g_χ	g_l	m_{med} [GeV]	m_χ [GeV]
Axial-vector	0.25	1	0	1460	415
Axial-vector	0.1	1	0.1	920	280
Vector	0.25	1	0	1470	580
Vector	0.1	1	0.01	950	400

WIMPs

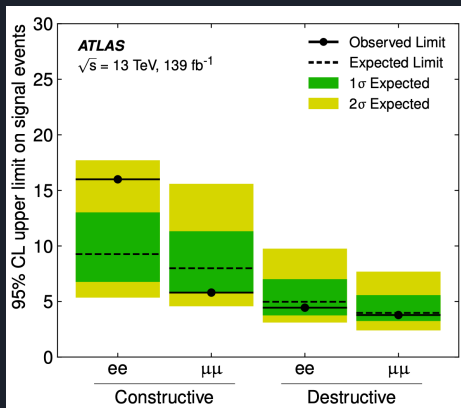
di-jet and dilepton final states

Phys. Lett. B 796 (2019) 68



Observed limit on fiducial $\sigma \times BR$ ranges from 3.6 (13.1) fb at 250 GeV to about 0.014 (0.018) fb at 6 TeV for the zero (10%) relative width signal in the combined dilepton channel.

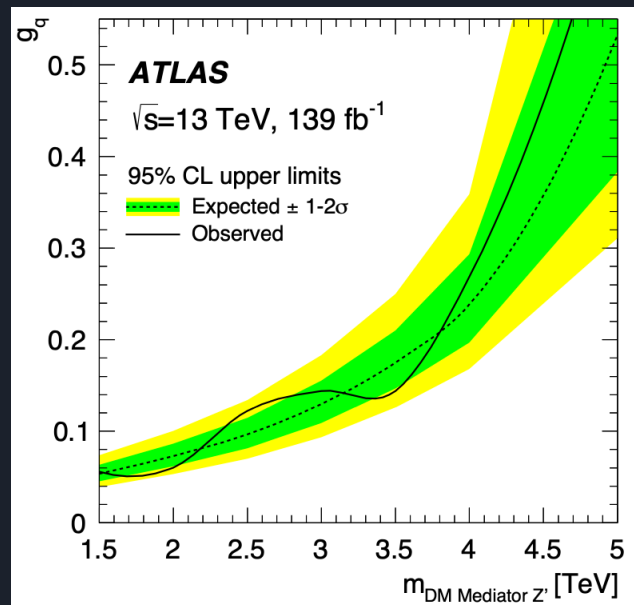
JHEP 11 (2020) 05



$$\frac{d\sigma}{dm_{\ell\ell}} = \frac{d\sigma_{\text{DY}}}{dm_{\ell\ell}} - \eta_{ij} \frac{F_I}{\Lambda^2} + \frac{F_C}{\Lambda^4}$$

Non-resonant dilepton: Interference between the DY and CI processes can be constructive or destructive depending on the sign of η_{ij}

JHEP 03 (2020) 145



Resonant dijet search: The 95% CL upper limits are set on the universal quark coupling g_q as a function of the Z' mass

What are the ingredients for a simplified/collider-friendly DM model?

Basic Ingredients:

- Generic signatures
- Evades constraints
- Manageable no. of parameters
- Promising dark matter candidate
 - ability to satisfy relic density

Spices/garnishes:

- Wide range of possible signatures
- Interesting phenomenology
- Potential synergies
 - decays: prompt vs LLP vs invisible
 - resonant vs non-resonant production
 - complementarity with direct/indirect detection



What did I bring to the table?



Searches for strongly-interacting dark sectors...

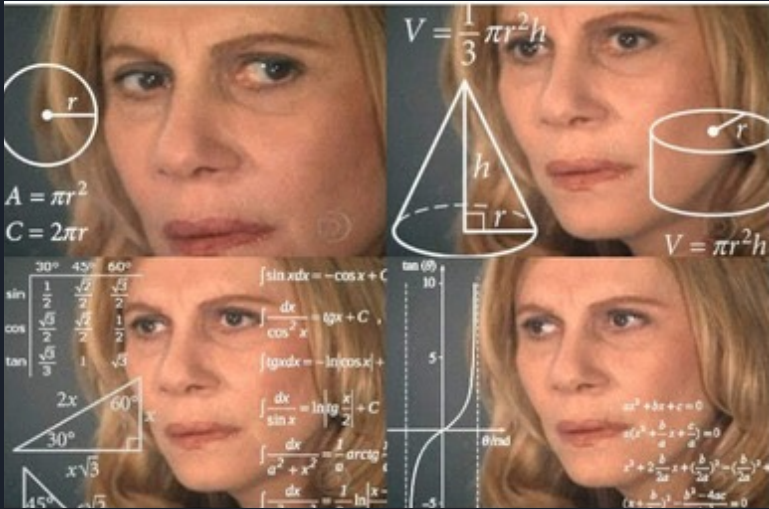


What did I bring to the table?

Searches for strongly-interacting dark sectors...

Confused?

Let's take a step back....



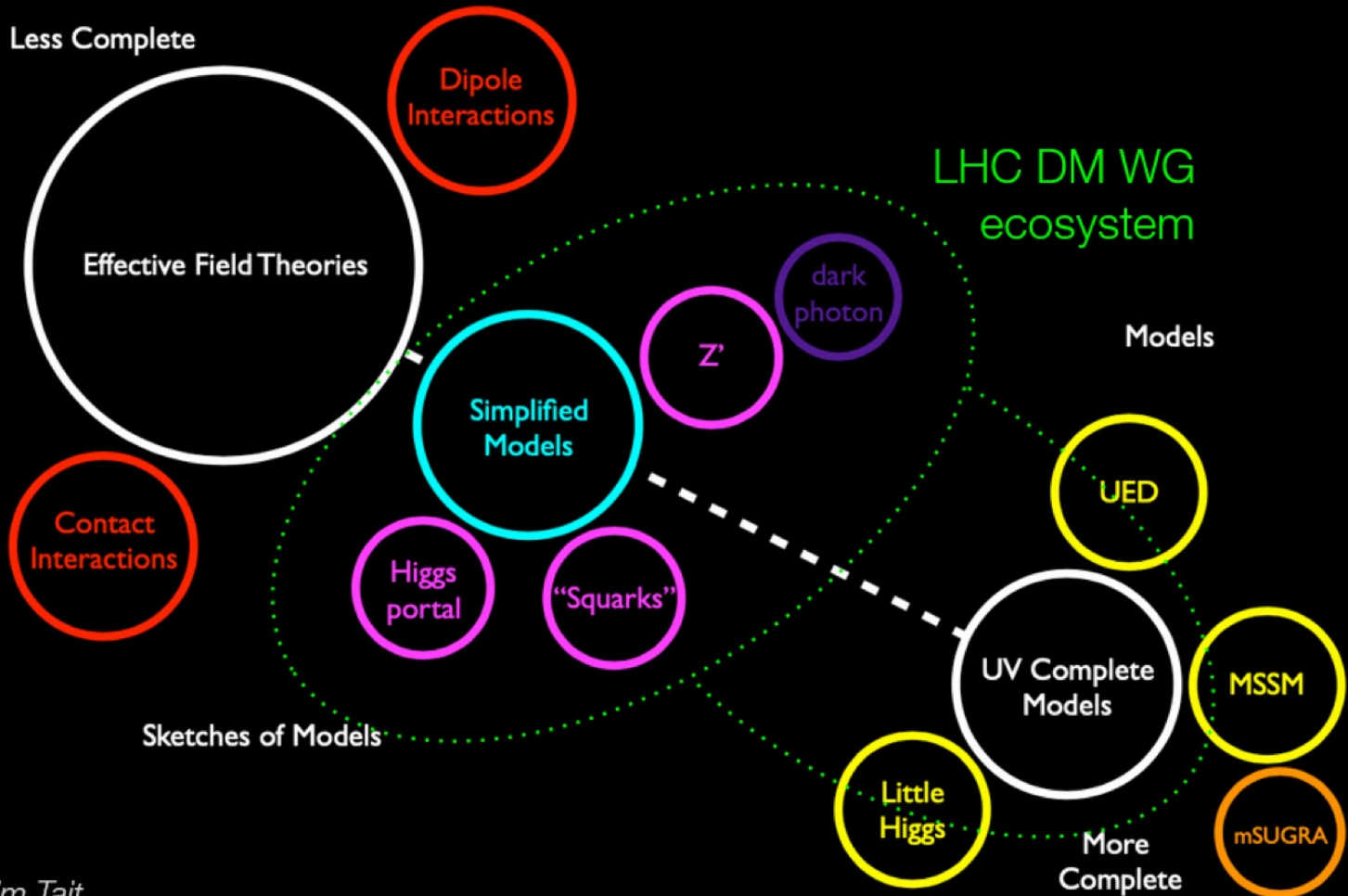


Figure: Tim Tait

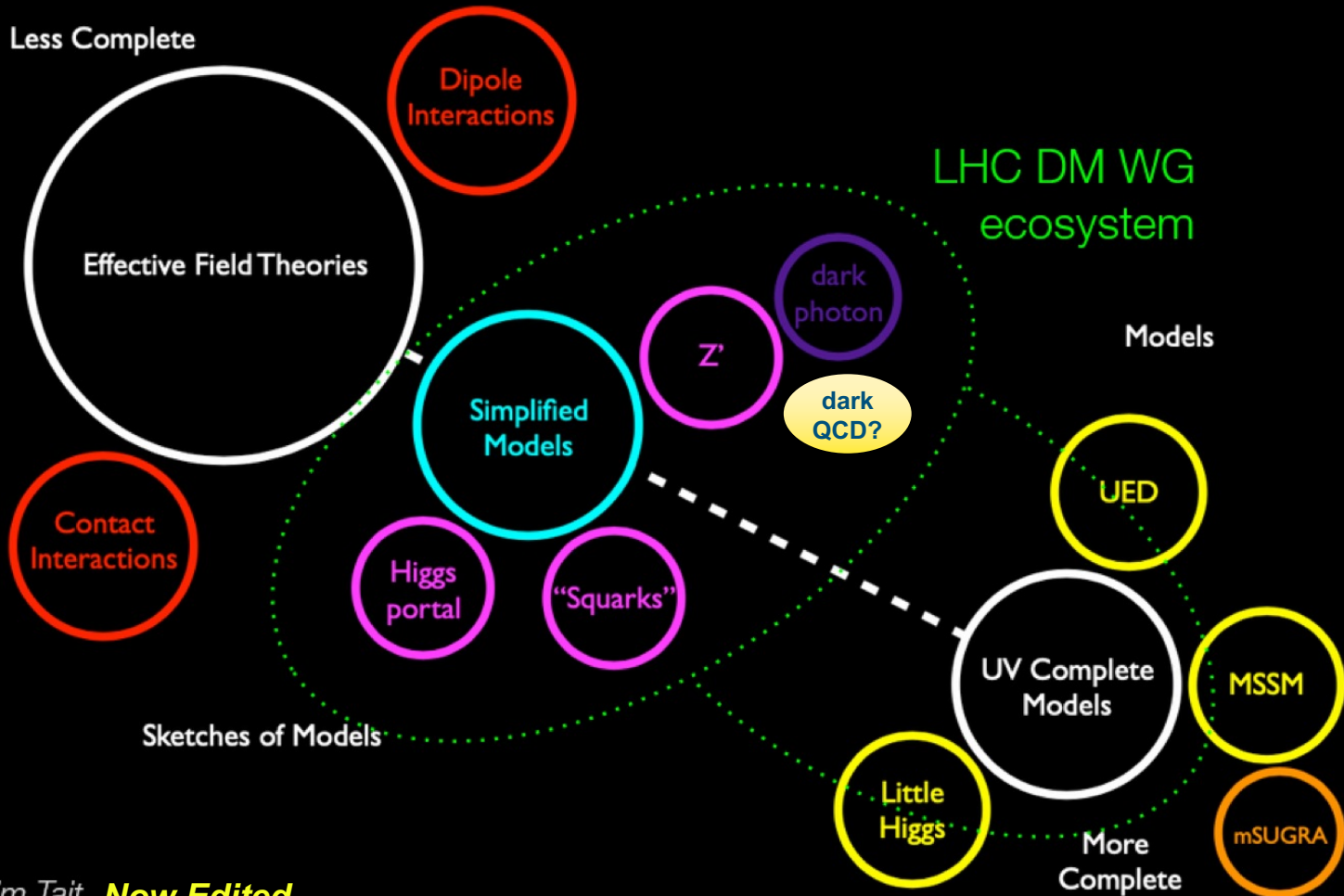


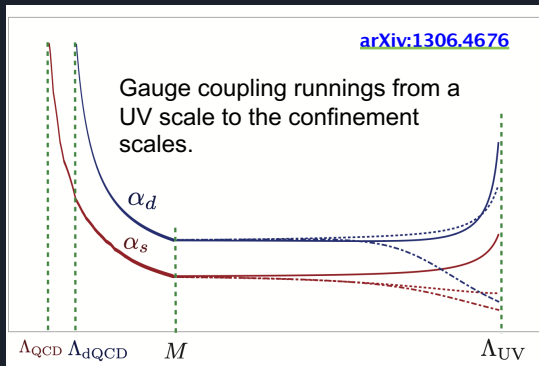
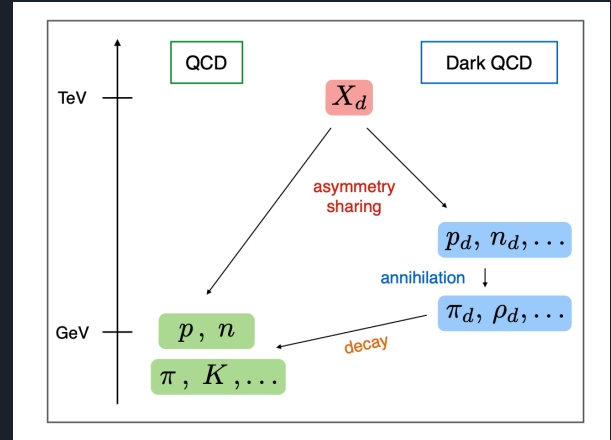
Figure: Tim Tait **Now Edited**

Dark QCD matters!

Hidden Valley may provide cosmologically required DM. No attempt to construct a specific model → set up reasonably generic framework for simulating variety of experimental signatures.

Different dark quark flavours combine to form π_d^+ , π_d^- , π_d^0 , and ρ_d^+ , ρ_d^- , ρ_d^0 (assumed to be produced thrice as much as pions)

- Neutral dark mesons are unstable and (promptly) decays to SM quarks: more likely to decay to b pairs due to need for a mass insertion, to make the angular momentum conservation work out
- Other mesons are (collider-)stable → invisible



Baryon and DM asymmetries shared via a mediator X_d → asymmetry in stable dark baryons.

The symmetric relic density annihilated into dark pions → decay into SM particles.

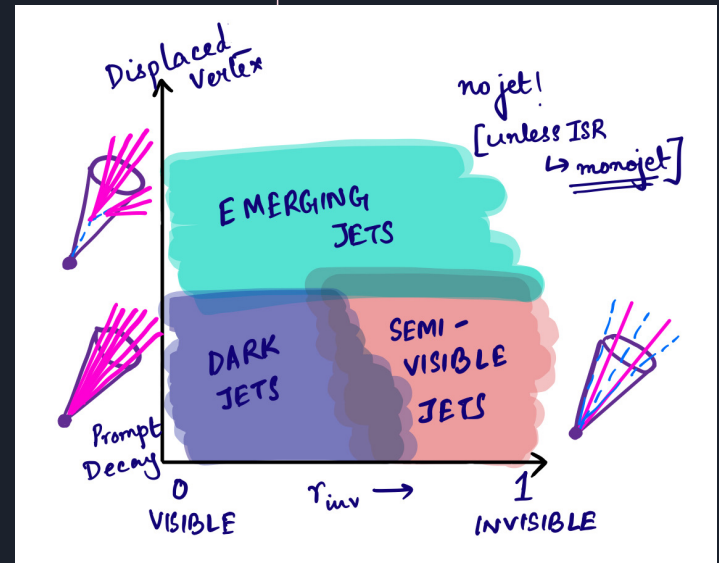
Correct DM relic density obtained when dark baryon masses are in the 10 GeV range.

Different UV boundary gauge couplings can lead to the same perturbative IR fixed points.

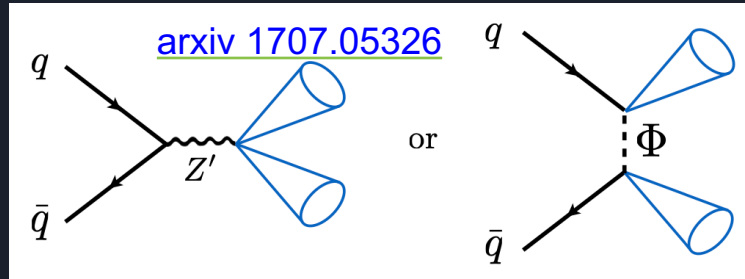
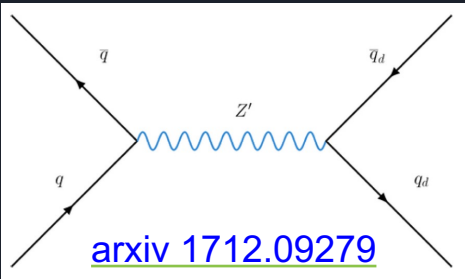
The final state signatures...

Dark hadrons decaying promptly in a QCD-like fashion,
 - fully, or
 - partially back to the visible sector
 Dark hadrons undergoing displaced decays in a QCD-like fashion

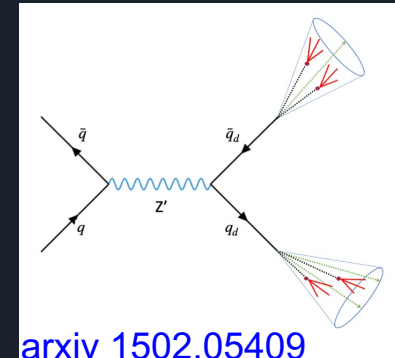
Each signature has a distinctive feature:
 -> Dark jets: unique substructure
 -> Emerging jets: displaced objects
 -> Semi-visible jets: substantial missing energy



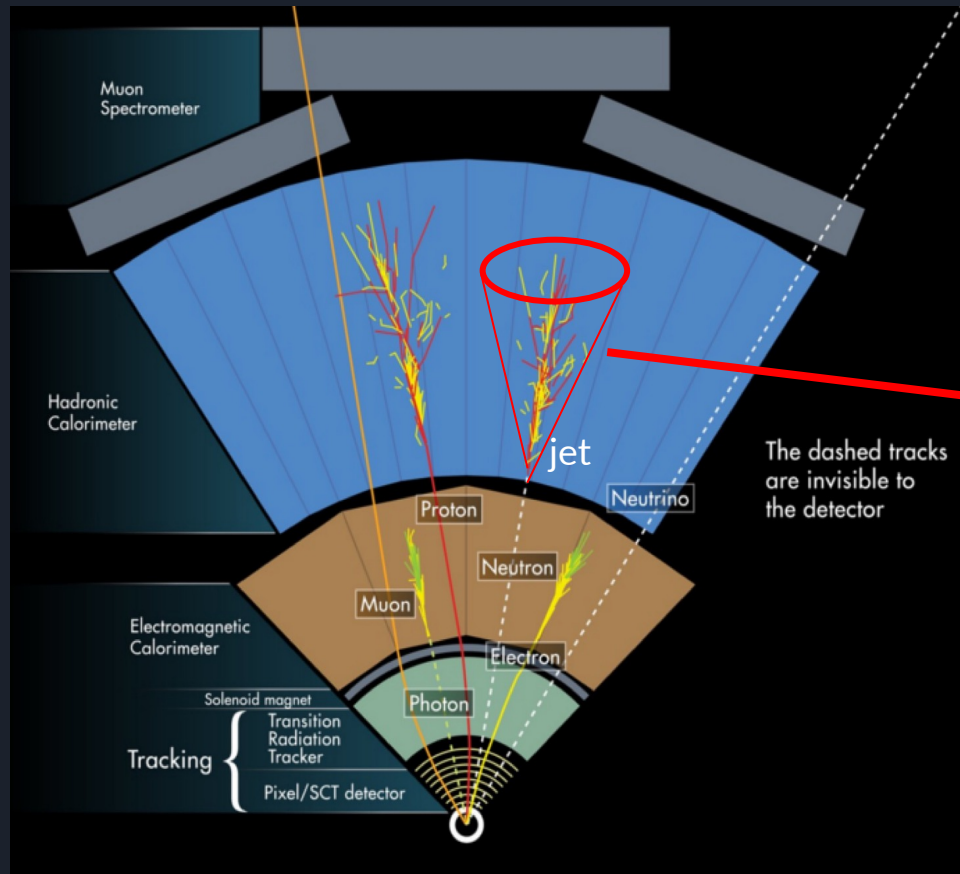
Showering using Pythia hidden valley module: at best a guesstimate!



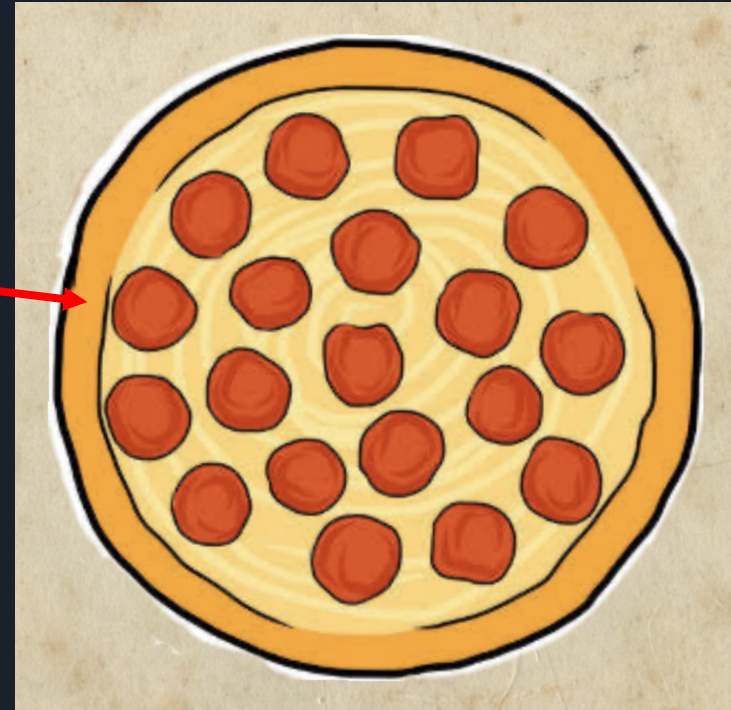
$$R_{inv} = \frac{\text{\#stable dark hadrons}}{\text{\#all dark hadrons}}$$



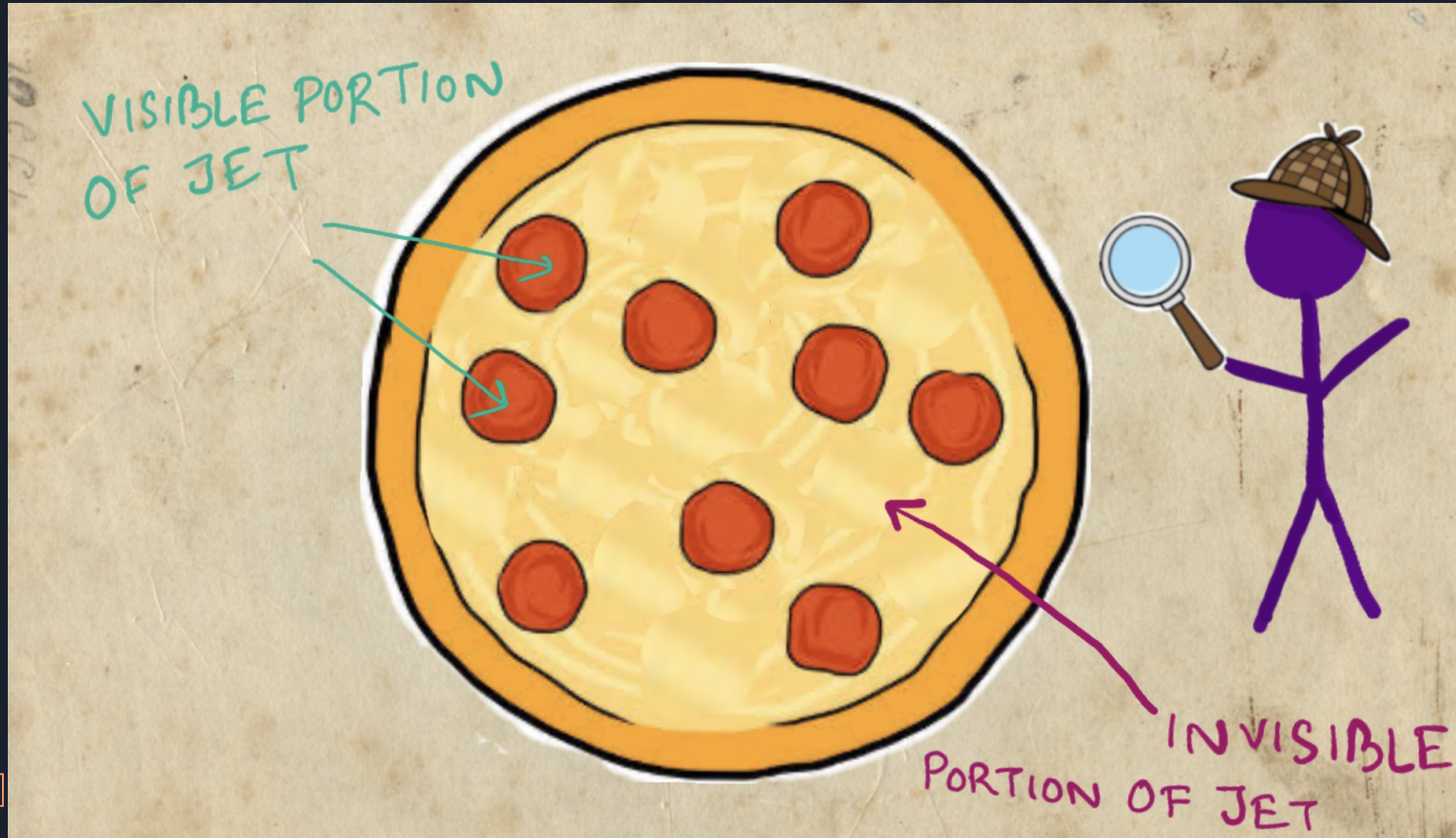
The final state signatures...



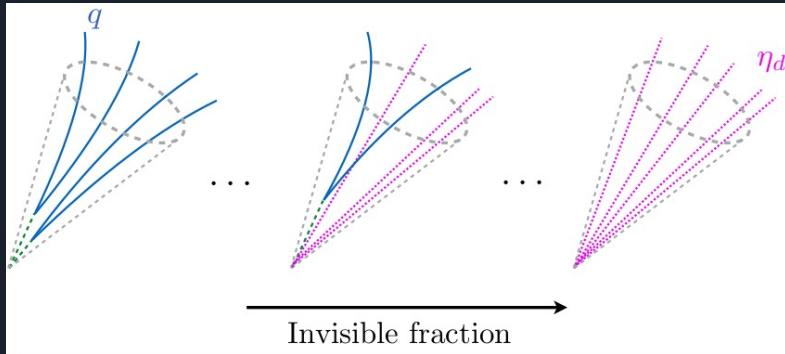
The dashed tracks are invisible to the detector



Semi-visible jets



ATLAS semi-visible jets t-channel production



Lagrangian containing the interaction and kinetic terms:

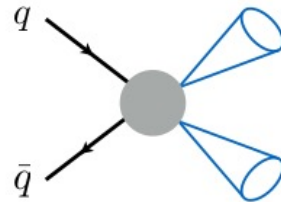
$$\mathcal{L}_{\text{dark}} \supset -\frac{1}{2} \text{tr} G_{\mu\nu}^d G^{d\mu\nu} - \bar{\chi}_a (i\not{D} - M_{d,a}) \chi_a$$

Model Parameters:

1. M_ϕ = Mass of Scalar Bi-fundamental mediator
2. $R_{\text{inv}} = \text{\#stable dark hadrons} / \text{\#all dark hadrons}$
3. M_d = Mass of dark hadrons

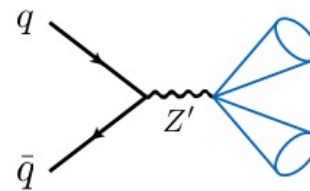
Dark sector \rightarrow SU(2) gauge theory with coupling $\alpha_d = \frac{g_d^2}{4\pi}$, containing two fermionic states.
Assuming minimal flavour-violation, light-flavour production channels dominate.

Contact Operator



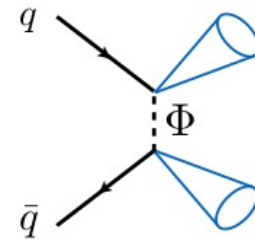
UV completions

s-channel



or

t-channel

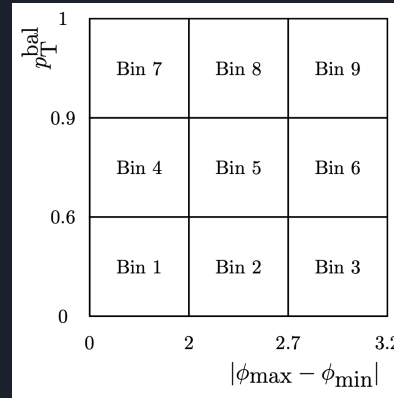
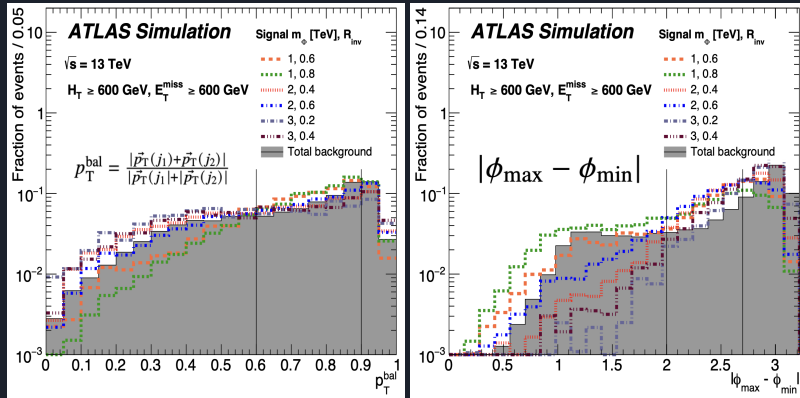


Link to the paper:
[1707.05326](https://arxiv.org/abs/1707.05326)

ATLAS semi-visible jets t-channel

Looking at two central $R=0.4$ jets, MET trigger, $\Delta\phi$ (closest jet, MET) < 2.0 , leading jet $p_T > 250$ GeV, $HT > 600$ GeV, $MET > 600$ GeV

Two key observables used to design a 9 bin grid: Yield in each bin treated as an observable



The signal events typically have high MET — better sensitivity for signals with higher mediator masses and R_{inv} fraction if search is performed at a high MET range.

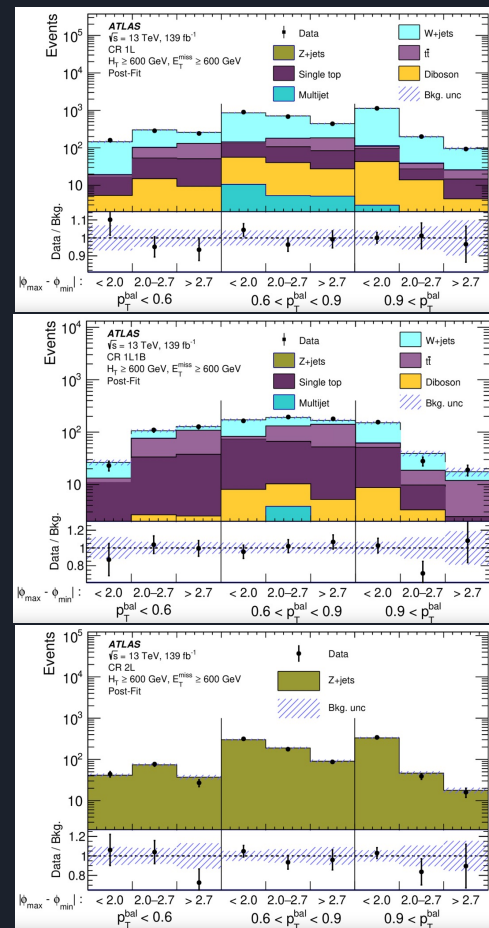
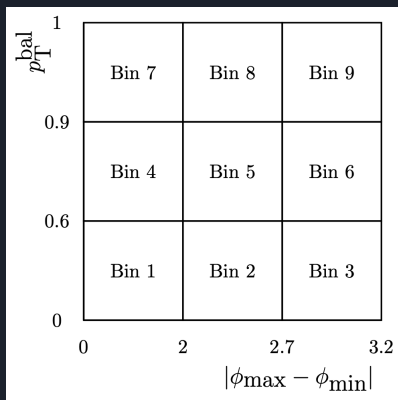
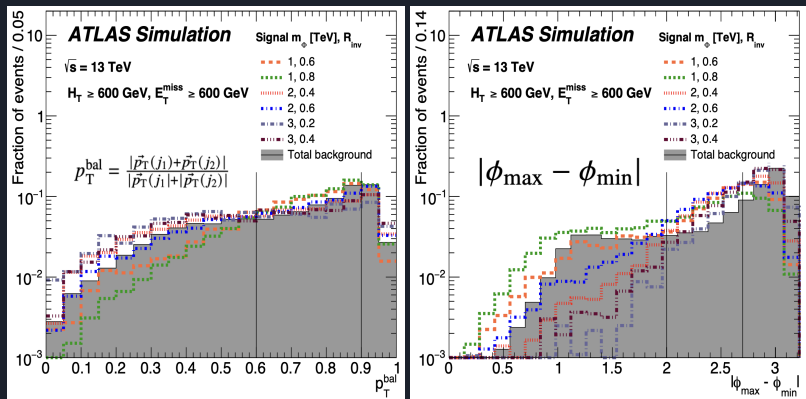
Background samples:

W/Z+jets, ttbar,
 singletop, multi-jet,
 diboson

ATLAS semi-visible jets t-channel

Looking at two central R=0.4 jets, MET trigger, $\Delta\phi$ (closest jet, MET) < 2.0, leading jet $p_T > 250$ GeV, HT > 600 GeV, MET > 600 GeV

Two key observables used to design a 9 bin grid: Yield in each bin treated as an observable



The signal events typically have high MET — better sensitivity for signals with higher mediator masses and R_{inv} fraction if search is performed at a high MET range.

$MET > 600$ GeV and $HT > 600$ GeV – define SR & CR (1L, 1L1B & 2L)

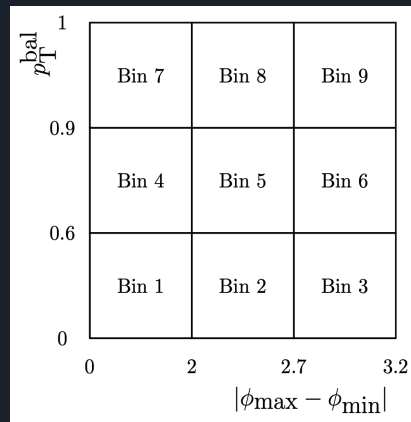
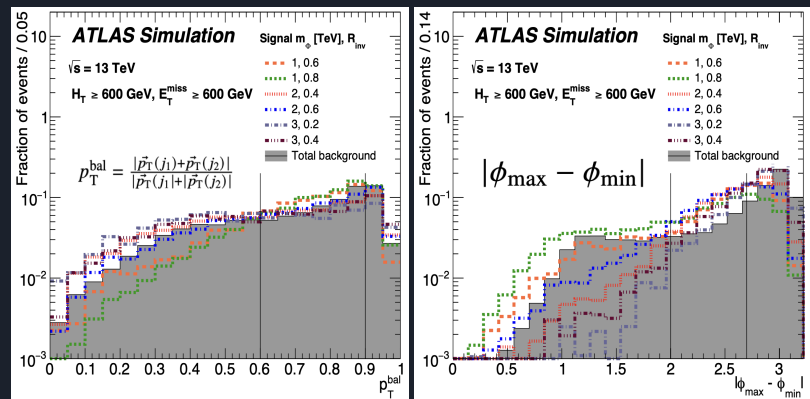
Partially data-driven method, simultaneously fit SR and three CRs to obtain scale factors for each bg process

Background samples:
 W/Z+jets, ttbar,
 singletop, multi-jet,
 diboson

ATLAS semi-visible jets t-channel

ATLAS-EXOT-2022-37

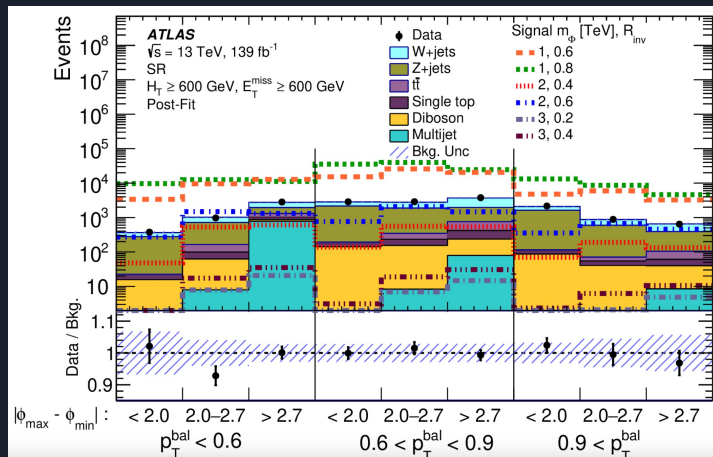
Two key observables used to design a 9 bin grid: Yield in each bin treated as an observable



The signal events typically have high MET — better sensitivity for signals with higher mediator masses and R_{inv} fraction if search is performed at a high MET range.

$MET > 600 \text{ GeV}$ and $H_T > 600 \text{ GeV}$ — define SR & CR (1L, 1L1B & 2L)

Partially data-driven method, simultaneously fit SR and three CRs to obtain scale factors for each bg process — absence of signal, good postfit agreement!

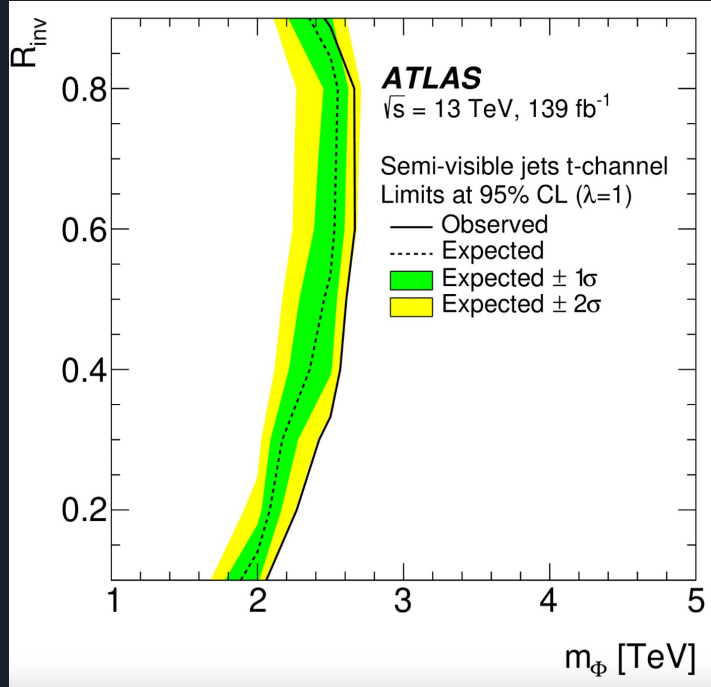
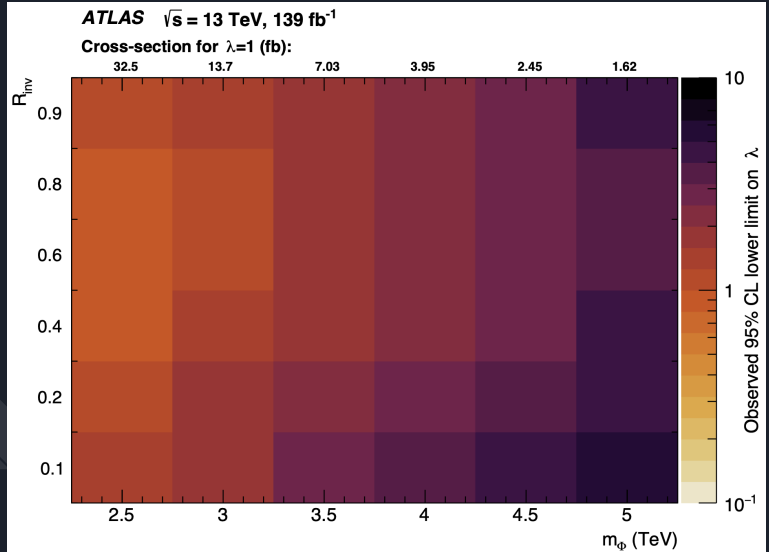


Results

First semi-visible jets result from ATLAS... but not the last!

The largest post-fit effects: signal modelling uncertainties ~8%, Z+jets modelling uncertainties ~7%, top process modelling uncertainties ~4%. The rest of the contributions are less than 2%.

- Excellent agreement between data and background prediction
- Assuming a coupling strength of unity between the mediator, a Standard Model quark and a dark quark, mediator masses up to 2.7 TeV can be excluded.
- For mediator mass of 2.5 TeV or higher can also express the limits in terms of the q-q_d-φ vertex coupling strength λ, as XS ~ λ⁴

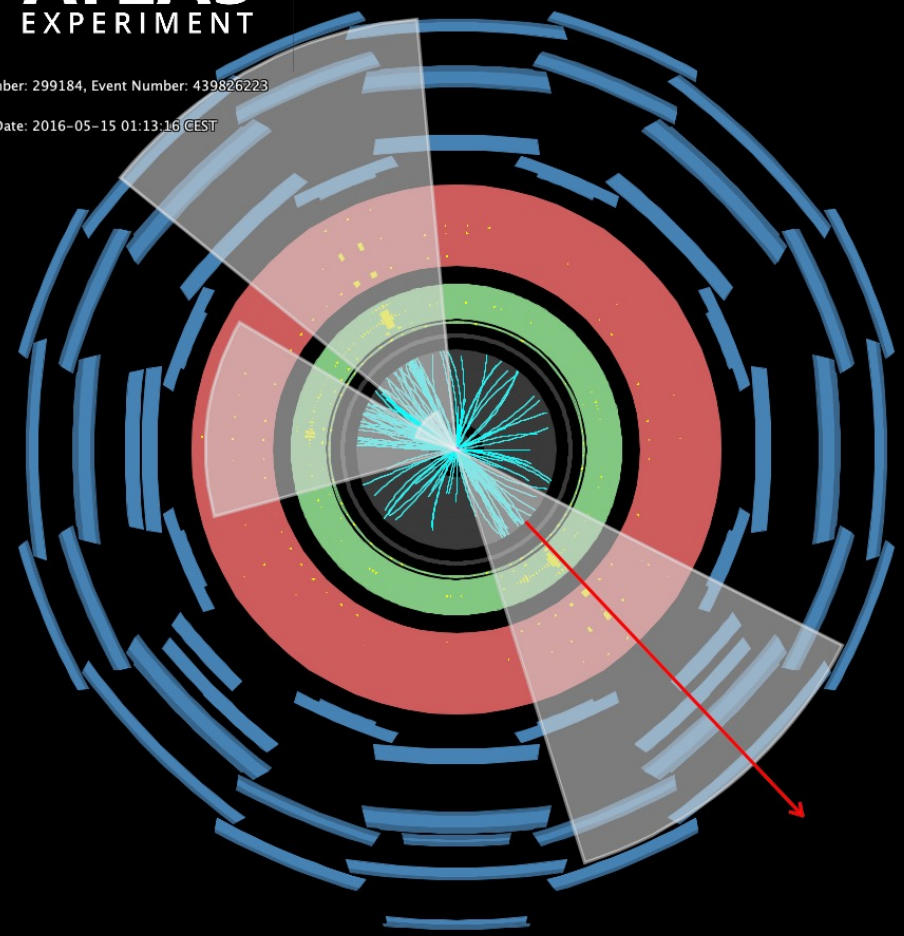


Results

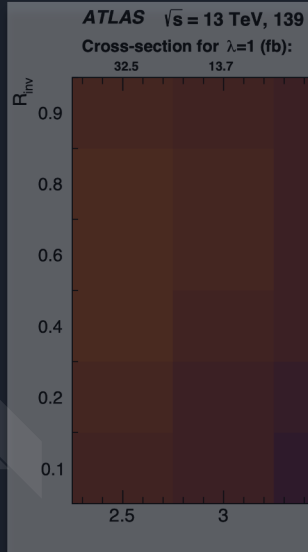
First
last



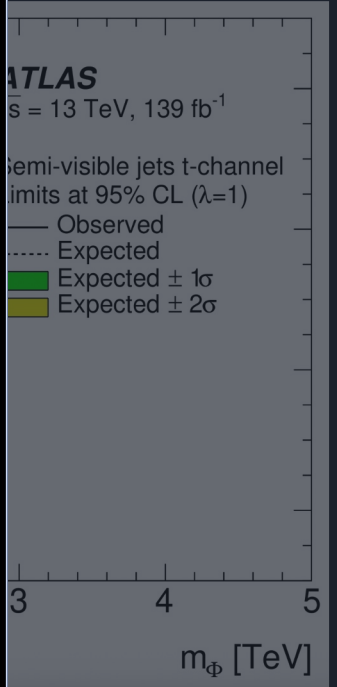
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- Excellent agreement
- Assuming a coupling to the Standard Model quarks, a mediator mass of 2.7 TeV can be excluded
- For mediator mass of 2.7 TeV, the exclusion limits in terms of the q - q_d - ϕ coupling are



effects: signal modelling
Z+jets modelling uncertainties
modelling uncertainties ~4%.
contributions are less than 2%.



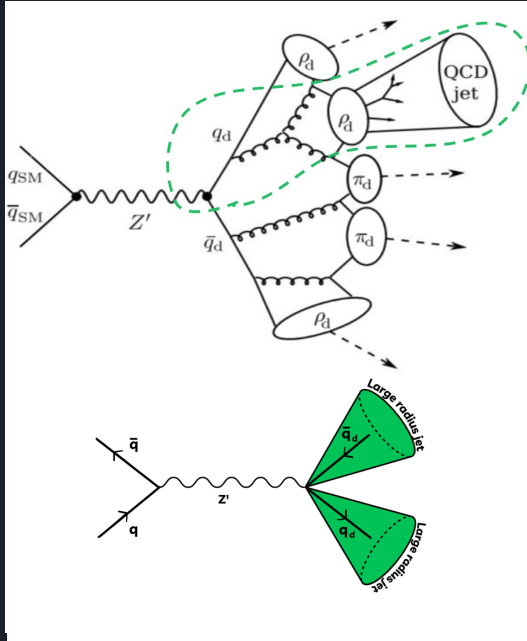
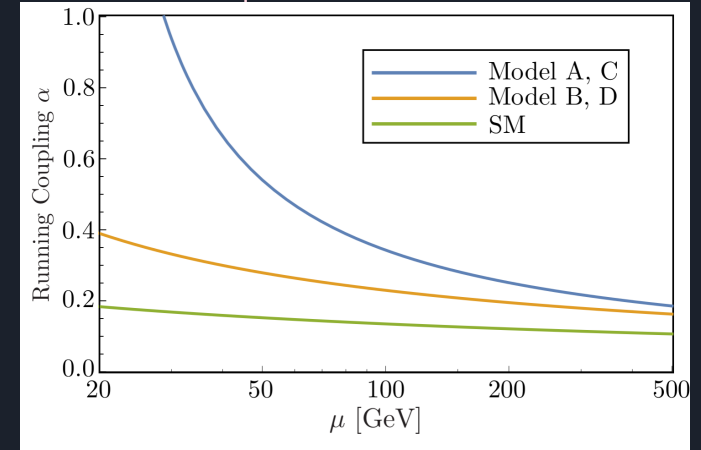
ATLAS dark jet resonances

4 different benchmark models (A, B, C, D)

Higher running coupling for dark sector models.

Main differences with respect to SM jets:

- Higher number of soft particles (and higher number of tracks).
- Larger jets, due to **double hadronization**



Signal masses generated using the Hidden Valley module of Pythia 8.235

- Signal masses per model from $m_{Z'}=1.5$ TeV to $m_{Z'}=3.5$ TeV in bins of 250 GeV for low masses and then 500 GeV
- Signal width not very narrow ($\sim 15\%$)

Model	n_f	Λ_d (GeV)	$\tilde{m}_{q'}$ (GeV)	m_{π_d} (GeV)	m_{ρ_d} (GeV)	π_d decay mode
A	2	15	20	10	50	$\pi_d \rightarrow c\bar{c}$
B	6	2	2	2	4.67	$\pi_d \rightarrow s\bar{s}$
C	2	15	20	10	50	$\pi_d \rightarrow \gamma'\gamma'$ with $m_{\gamma'} = 4.0$ GeV
D	6	2	2	2	4.67	$\pi_d \rightarrow \gamma'\gamma'$ with $m_{\gamma'} = 0.7$ GeV

ATLAS dark jet resonances

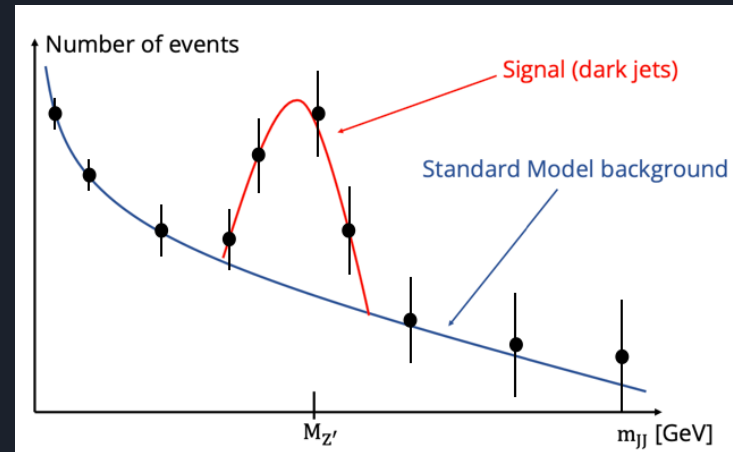
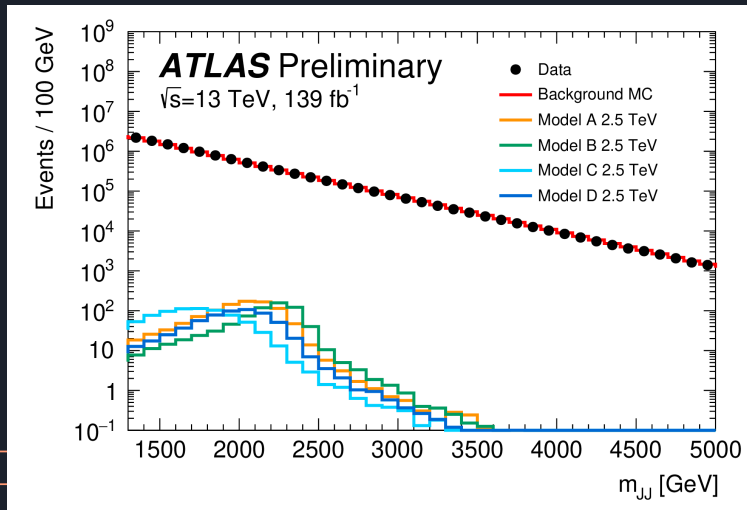
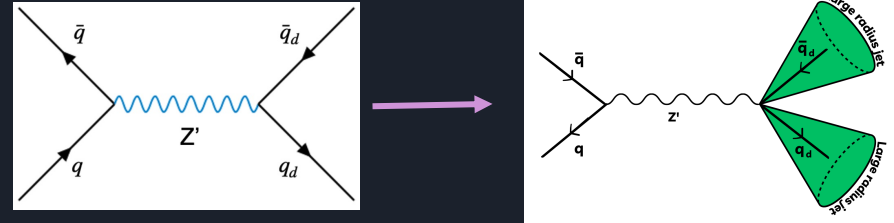
- **Reconstructed Objects:** Two Large radius jets (Trimmed LCTopo jets)
- **Trigger on Large-R jets**
- **Preselection cuts:** $|\eta(j_{1,2})| < 2.0$

$$p_T(j_1) > 500 \text{ GeV}, p_T(j_2) > 400 \text{ GeV}$$

$$m(j_{1,2}) > 50 \text{ GeV}$$

Signal scenario: Two dark jets

- Tag two dark jets using substructure information
- Look for a resonance over dijet background



ATLAS dark jet resonances

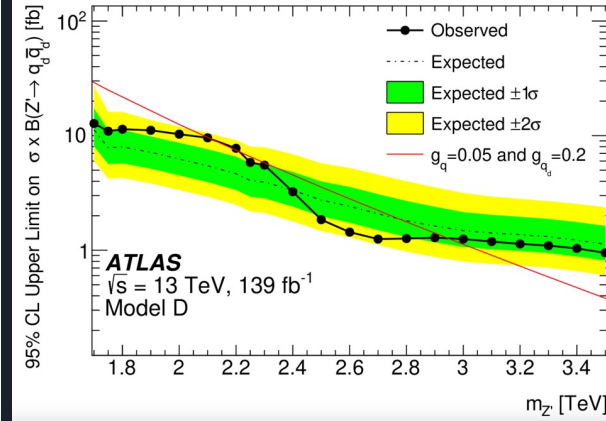
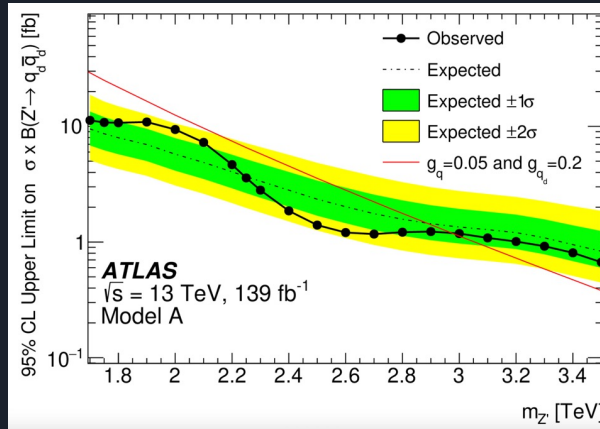
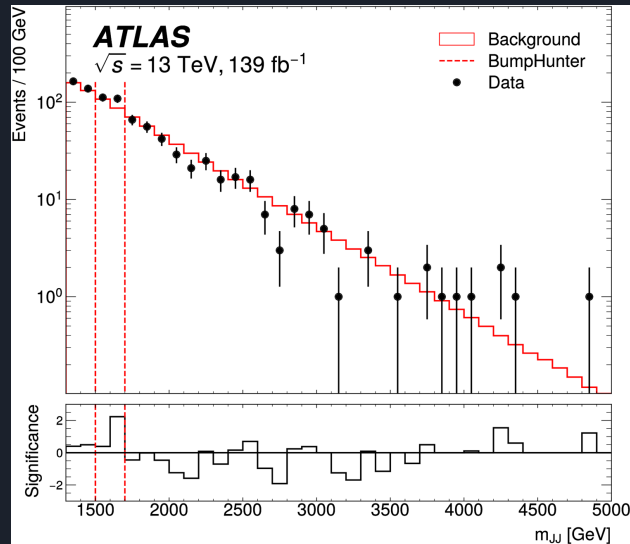
Two large R jets with a high track multiplicity
 \rightarrow Tag dark jets using substructure information \rightarrow study ungroomed no. of tracks associated to the jets, n_{ETK} , which is decorrelated from dijet mass \rightarrow obtain bkg. template from dedicated CRs \rightarrow look for a m_{JJ} resonance over the dijet background [BumpHunter]

Search for Resonant Production of Dark Quarks in the Di-jet Final State with the ATLAS Detector

This paper presents a search for a new Z' resonance decaying into a pair of dark quarks which hadronise into dark hadrons before promptly decaying back to Standard Model particles. This analysis is based on proton-proton collision data recorded at $\sqrt{s} = 13$ TeV with the ATLAS detector at the Large Hadron Collider between 2015 and 2018, corresponding to an integrated luminosity of 139 fb^{-1} . After selecting events containing large-radius jets with high track multiplicity, the invariant mass distribution of the two leading jets is scanned to look for an excess above a data-driven estimate of the Standard Model multi-jet background. No significant excess of events is observed and the results are thus used to set 95% confidence-level upper limits on the production cross-section times branching ratio of the Z' to dark quarks as a function of the Z' mass for various dark-quark scenarios.

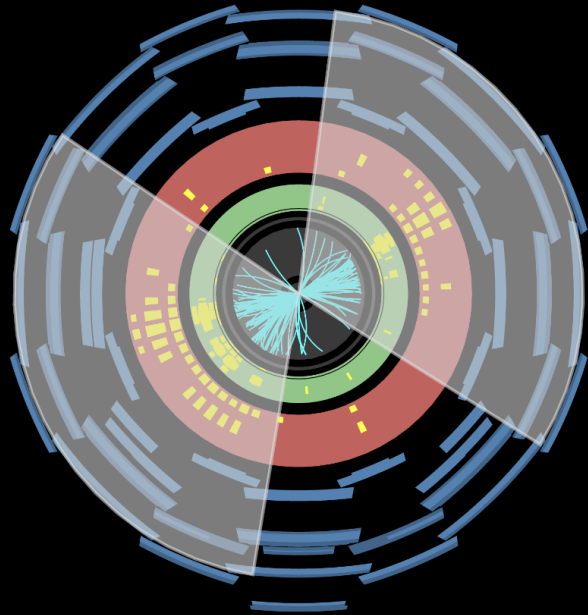
Highest excess: between $1500 < m_{JJ} < 1700$ GeV, with a p-value of 0.63

[JHEP 02 \(2024\) 128](#)



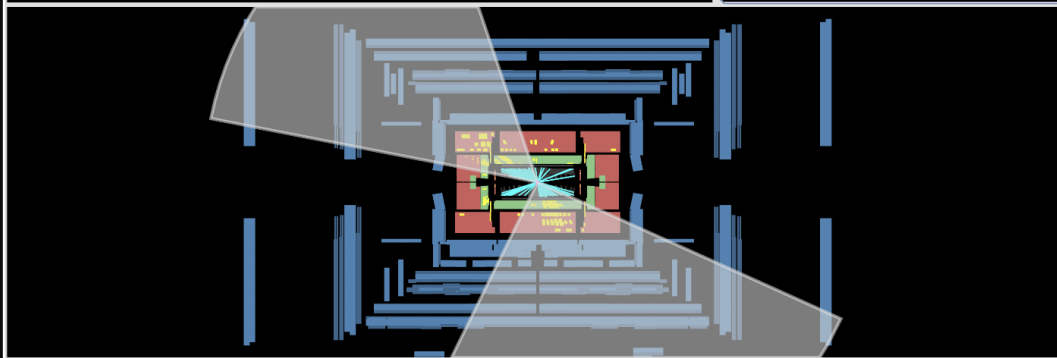
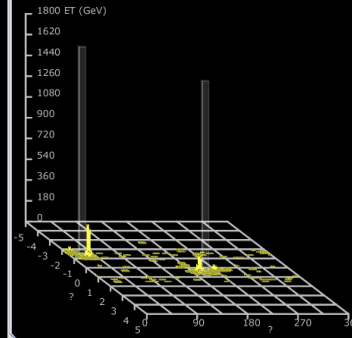
ATLAS dark

Two largeR jets with a high
 → Tag dark jets using substructure
 ungroomed no. of tracks
 decorrelated from dijet
 dedicated CRs → look for
 background [Bumphunter]



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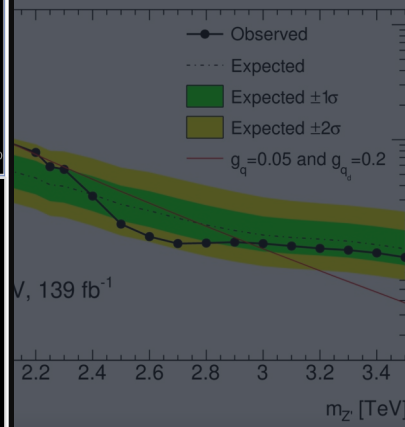
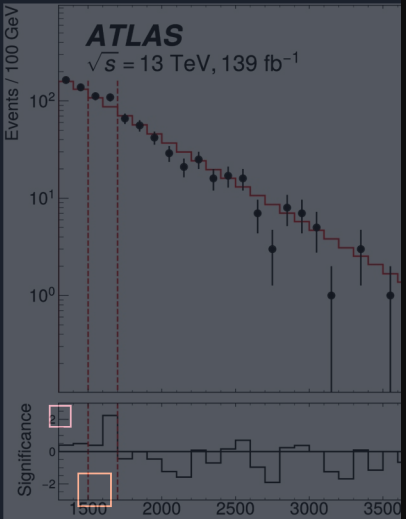
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Search for Dark Quarks in the Di-jet Final State at the ATLAS Detector

decaying into a pair of dark quarks which hadronise into Standard Model particles. This analysis is based on 139 fb⁻¹ of data collected at 13 TeV with the ATLAS detector at the Large Hadron Collider. After selecting the signal region, the invariant mass distribution of the di-jet system is compared to a data-driven estimate of the Standard Model background. If a significant excess is observed and the results are thus used to set upper limits on the cross-section times branching ratio of the Z' to dark-quark scenarios.

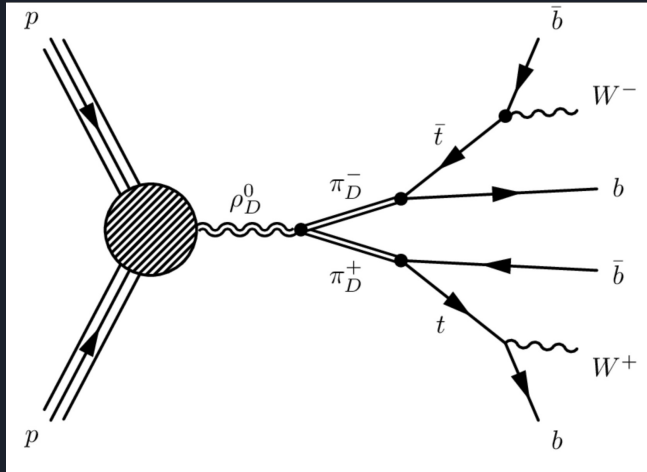
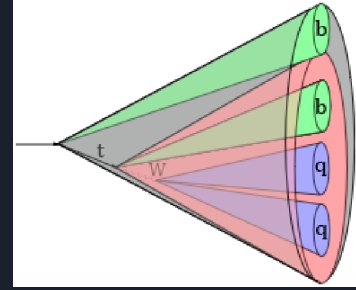
Highest excess: between 1500 and 2000 GeV



ATLAS dark mesons search

For $\eta (= \frac{m_{\pi_d}}{m_{\rho_d}}) < 0.5$, ρ_d decay to π_d pairs, resulting in $t\bar{t}b\bar{b}$ and $t\bar{t}b\bar{b}$ signatures (in all hadronic decay mode)

Dark pion identified from jet mass and the flavour tagging of its constituents, reconstruction using reclustering of anti- k_R radius 0.4 PFlow jets with radius 1.2

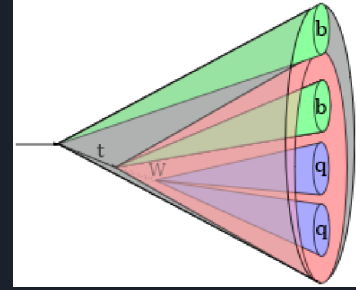


Excellent postfit agreement, limit on dark pion masses, large gain in coverage

ATLAS dark mesons search

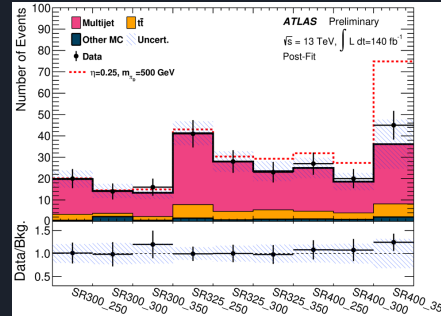
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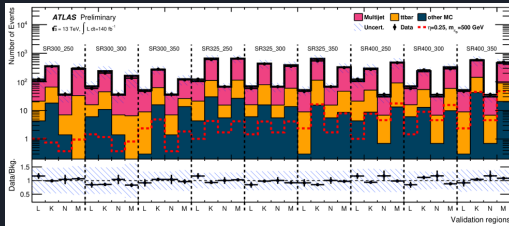
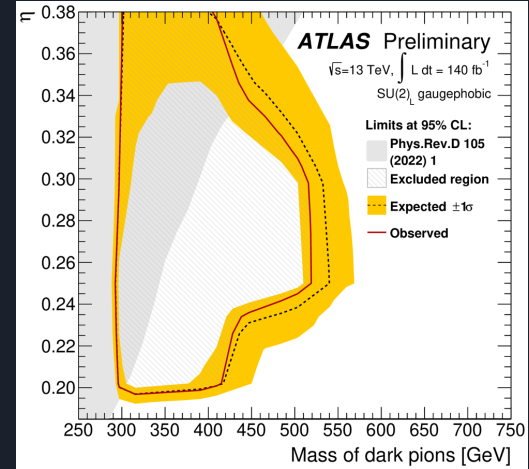


Discriminating variables:

- $\Delta R(j, b_2)$: Distance between a RC jet and the second closest b-jet
- $m_{bb} / p_{T,bb}$: for the closest b-jet pair to an RC jet to suppress QCD
- $m_{jet, R=1.2}$: Mass of the RC jet, main discriminant



Agreement in VRs



4-variable ABCD method to estimate multijet background based on b-tag and RC jet mass

Excellent postfit agreement, limit on dark pion masses, large gain in coverage

Advertisement break

MITP Youngst@rs dark showers workshop

(Oct 2023): [Link](#)

Scope of the workshop: Aim to build collaboration and motivate cross-talk between the experimental and theory community dedicated towards developing and understanding the strongly interacting dark sector.

- **All discussions summarised in report, now on arXiv [[arXiv:2311.16330](#)]**



arXiv:2311.16330v1 [hep-ph] 27 Nov 2023

MITP Colours in Darkness workshop summary report

Jon Butterworth^a, Cesare Cazzaniga^b, Aran Garcia-Bellido^c, Deepak Kar^d, Suchita Kulkarni^e, Pedro Schwaller^f, Sukanya Sinha^g, Danielle Wilson-Edwards^g and Jose Zurita^h

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^cDepartment of Physics and Astronomy, University of Rochester, Rochester NY, USA

^dSchool of Physics, University of Witwatersrand, Johannesburg, South Africa

^eInstitute of Physics, NAWI Graz, University of Graz, Graz, Austria

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^hInstituto de Física Corpuscular, CSIC-Universitat de València, Paterna, Spain

* **Corresponding author(s): Sukanya Sinha** *

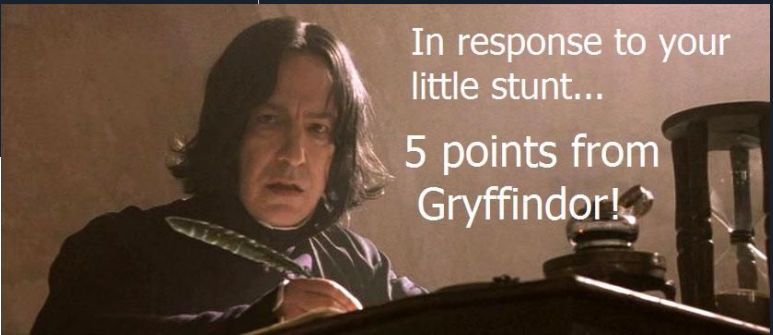
ABSTRACT

This report summarises the talks and discussions that took place over the course of the MITP Youngst@rs Colours in Darkness workshop 2023. All talks can be found at this URL: <https://indico.mitp.uni-mainz.de/event/377/>.

1 Introduction

In recent years, there has been an increase in the number of search programmes exploring the possibility of a “dark sector” beyond the Standard Model (BSM) using LHC data. To date, dark matter (DM) searches at the Large Hadron Collider (LHC) have usually focused on WIMPs (Weakly Interacting Massive Particles), but since the standard signatures have found no compelling evidence, several recent phenomenology papers have explored the possibility of accessing the dark sector with unique collider topologies. If dark mesons exist, their evolution and hadronization procedure are currently little constrained. They could

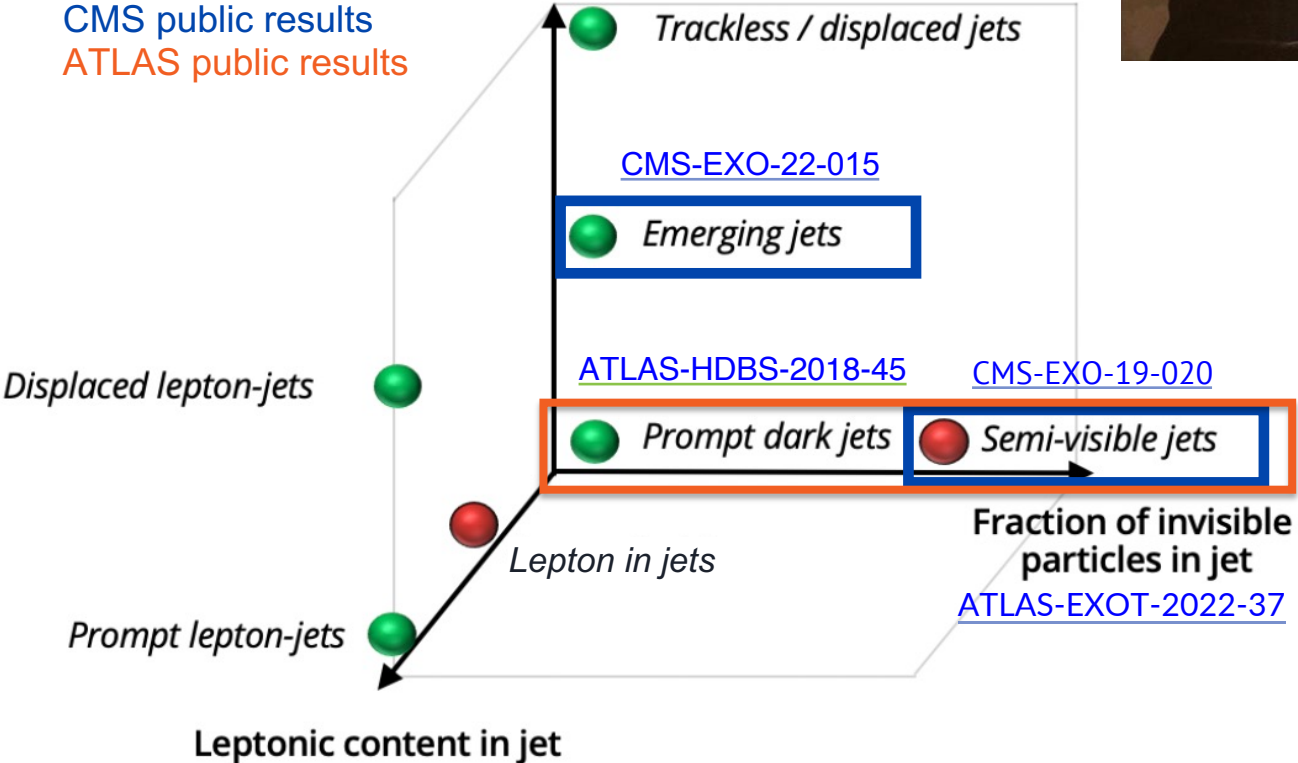
*sukanya.sinha@cern.ch



In response to your little stunt...
5 points from Gryffindor!

Lifetime of dark particles in jet

CMS public results
ATLAS public results



New topologies of interest

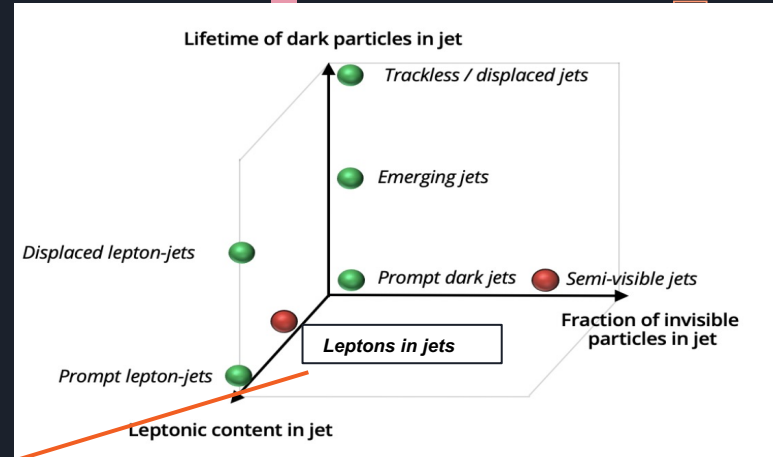
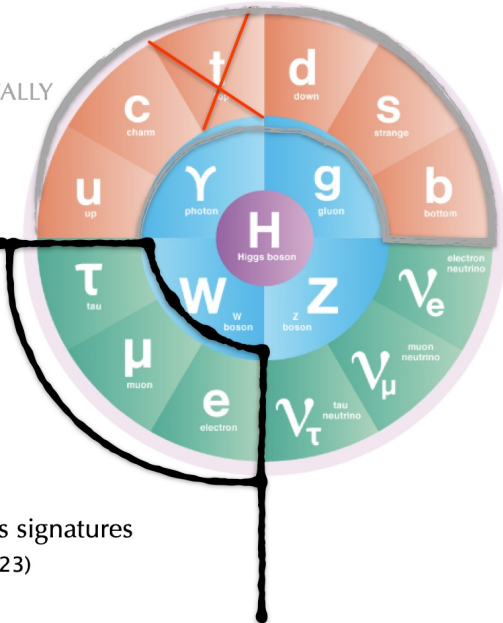
EXPLORED THEORETICALLY
& EXPERIMENTALLY

OUR FOCUS

Leptons-enriched
signatures for SVJs

Eur. Phys. J. C 82, 793 (2022)

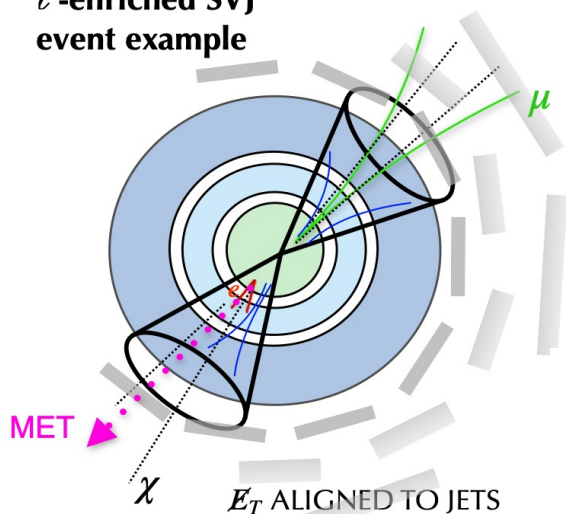
→ τ leptons-enriched SVJs signatures
Eur. Phys. J. C 83, 599 (2023)



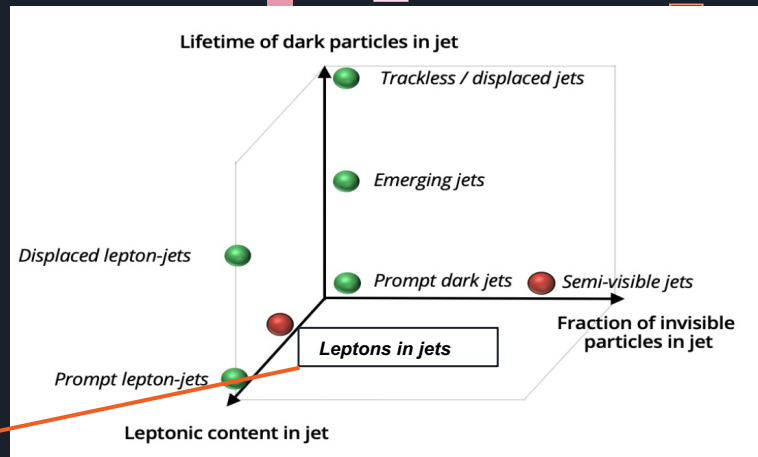
Semi-visible jets with μ/e leptons

[arXiv:2206.03909v3](https://arxiv.org/abs/2206.03909v3)

ℓ -enriched SVJ
event example



From C. Cazzaniga's DS workshop [talk](#)



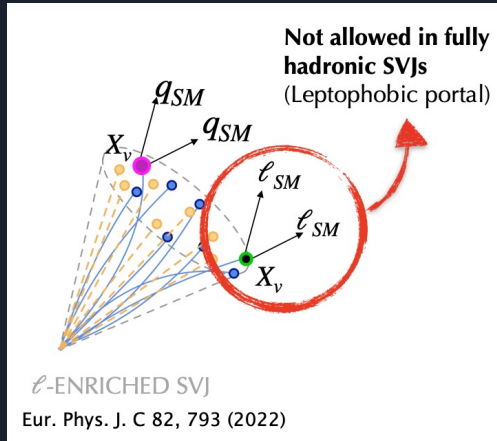
Leptophobic Z' mediator to evade constraints from high-mass di-lepton searches

Phenomenological aspects of semi-visible jets with μ/e

Simple modification to vanilla svj signature

- Favour dark bound state decays to both leptons and quarks
- Add an additional dark photon as a mediator, satisfying $m_{A'} \ll m_{Z'}$
- prompt decays of A' are allowed.

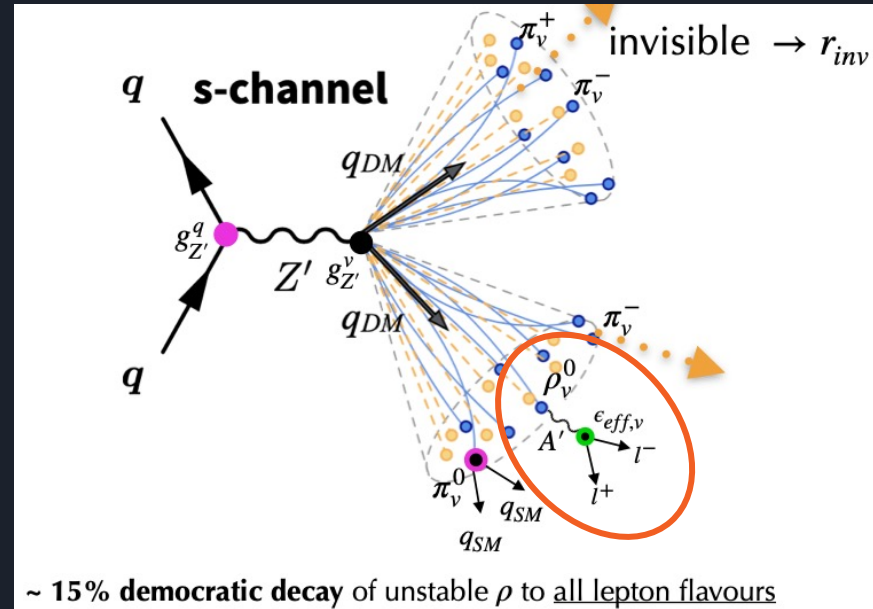
Decay width of DM bound states scales as $\sim 1/M^4 \rightarrow$ suppressing the off-shell Z' mediated decays



New decay chain:
 $\rho_v^0 \rightarrow A' A'$; $A' \rightarrow e^+ e^-$

Lagrangian \mathcal{L}'' for the A' messenger sector is:

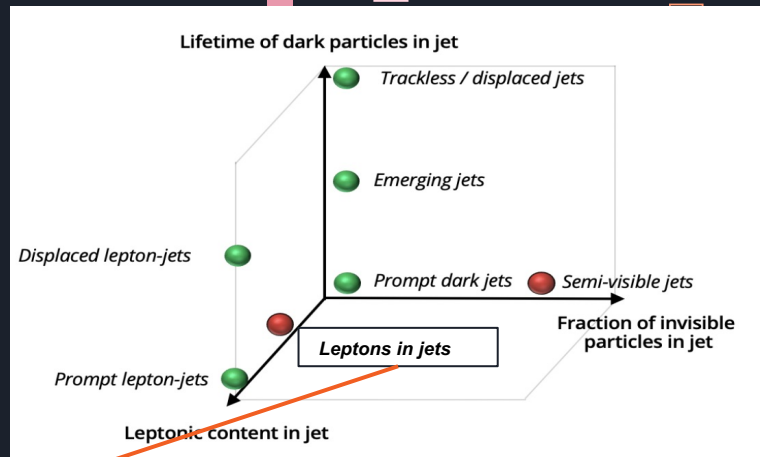
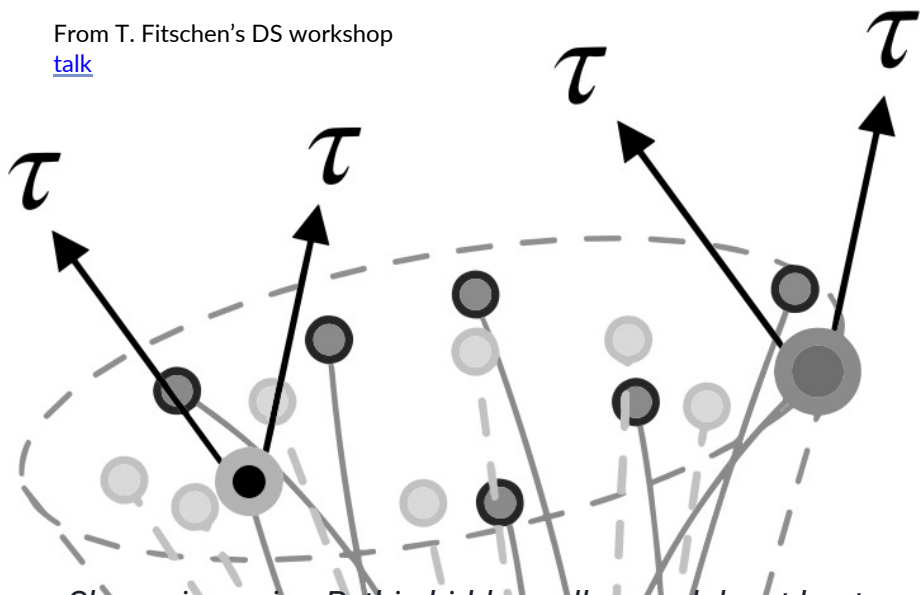
$$\mathcal{L}'' \supset -\frac{1}{4}F_{\mu\nu}[A']F^{\mu\nu}[A'] + \frac{1}{2}M_{A'}^2 A'_\mu A'^\mu - \epsilon e Q A'_\mu J_{q,SM}^\mu - \epsilon e Q A'_\mu J_{l,SM}^\mu - g_{A'}^v A'_\mu J_v^\mu. \quad (1)$$



Semi-visible jets with tau leptons

[arXiv:2212.11523v2](https://arxiv.org/abs/2212.11523v2)

From T. Fitschen's DS workshop
[talk](#)



Leptophobic Z' mediator to evade constraints from high-mass di-lepton searches

Phenomenological aspects of semi-visible jets with tau leptons

Simple modification to vanilla svj signature

→ Standard Z' coupling to dark quarks

→ Add an additional Z' coupling to SM quarks and leptons

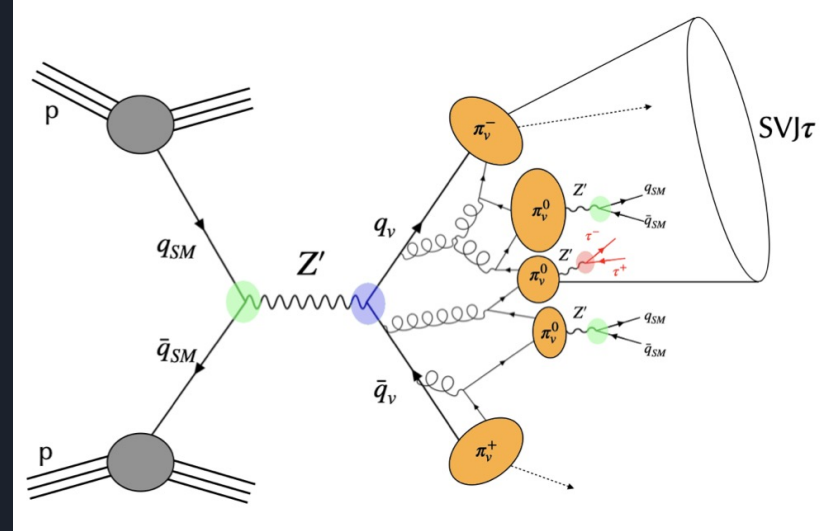
- Z' coupling to e/mu suppressed

Z' coupling to SM q and ℓ :

$$\begin{aligned} \mathcal{L}_{SM} \supset & -Z'_\mu \bar{u}_i \gamma^\mu (g_{ij}^{uR} P_R + g_{ij}^{uL} P_L) u_j \\ & -Z'_\mu \bar{d}_i \gamma^\mu (g_{ij}^{dR} P_R + g_{ij}^{dL} P_L) d_j \\ & -Z'_\mu \bar{e}_i \gamma^\mu (g_{ij}^{eR} P_R + g_{ij}^{eL} P_L) e_j \end{aligned}$$

Z' coupling to dark q_v :

$$\mathcal{L}_{q_v} \supset -Z'_\mu \bar{q}_{vi} \gamma^\mu (g_{ij}^{q_{vR}} P_R + g_{ij}^{q_{vL}} P_L) q_j$$



⇒ Dark mesons decaying into SM:

$$\Gamma_{\pi_a \rightarrow f_i \bar{f}_i} = \frac{N_c F_\pi^2}{32\pi} (\Delta_{ii}^f \Delta_a^{q_v})^2 \frac{m_{\pi_a} m_{f_i}^2}{M_{Z'}^4} \sqrt{1 - \frac{4m_{f_i}^2}{m_{\pi_a}^2}}$$

B-phillic SVJ

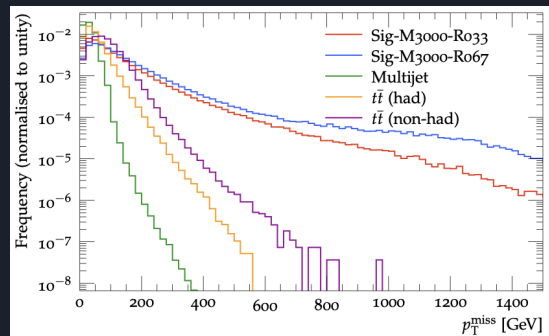
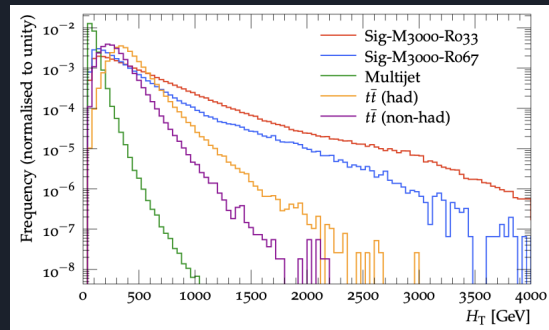
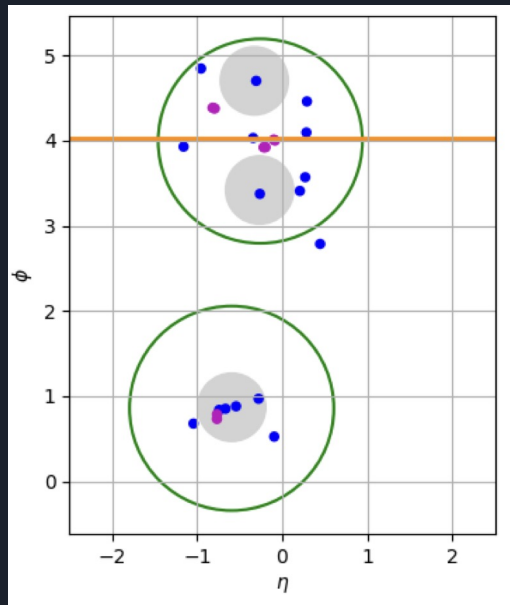
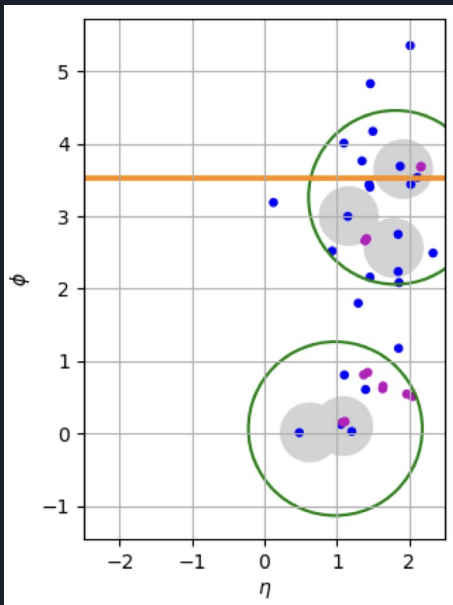
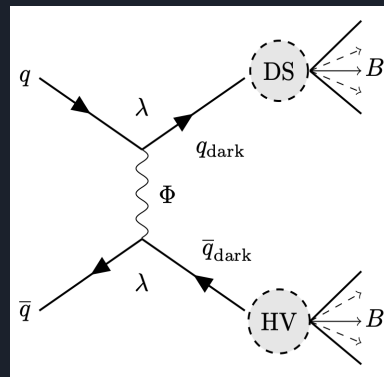
DK, WN, SS: [arXiv:2207.01885](https://arxiv.org/abs/2207.01885)

Extra handle on SVJs in collider search, no significant collider constraints

If mass of π_d^0 is lighter than other hadrons:

Helicity flipping suppression forces $\pi_d \rightarrow b \bar{b}$ (if $2m_b < m_{\pi_d^0} < 2m_t$)

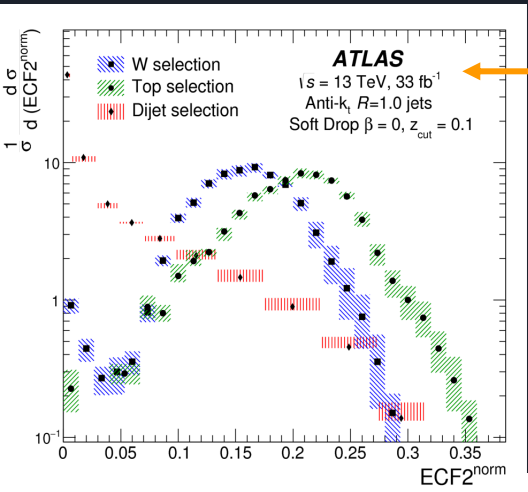
i.e., b-philic mode of hidden valley pion decay will be preferred unless there is a hierarchical κ matrix that opposes mass enhancement.



Requiring high H_T and p_T^{miss} can result in having a good S/\sqrt{B} significance.

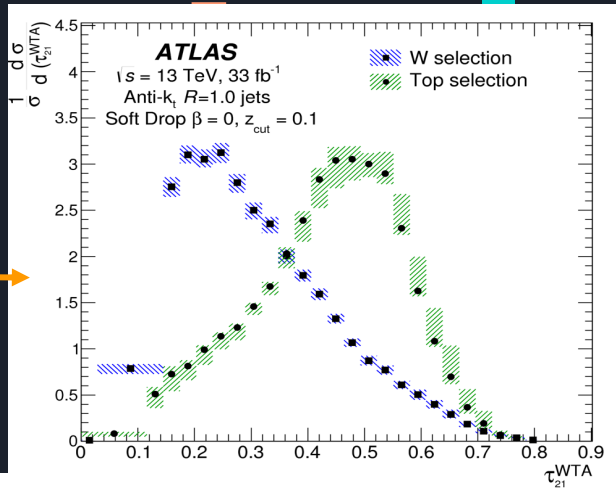
signal reconstruction improved by using [variable-radius jets](#)

Using jet substructure observables: Example plots from ATLAS

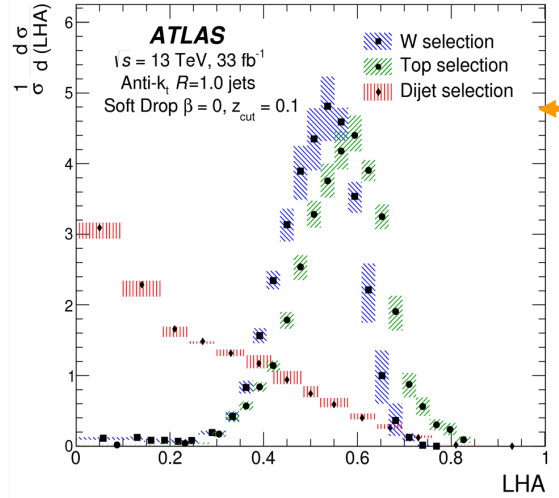
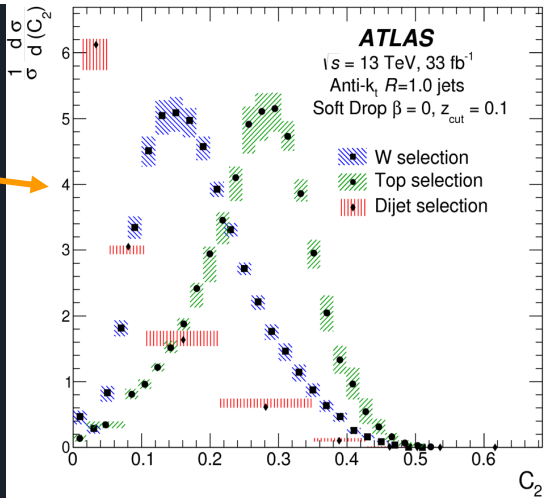


Energy correlation functions:
 ECF2:
 multi-prong has higher values

N-Subjettiness:
 τ_{21} : Lower values indicate more 2 subjet like behaviour

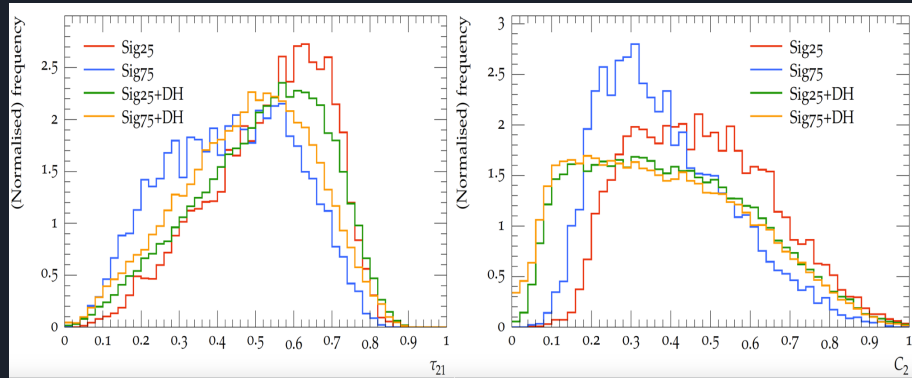


Energy correlation double ratios:
 C2: higher value has more subjects



Les Houches Angularity:
 higher value means
 hard radiations are more
 separated

 [Link to paper](#)



Comparing jet substructure variables to see if SVJ substructure is different from light quark/gluon jets (BG). Do they behave more multi-pronged as opposed to mostly single prong?

Comparison can be done in p_T bins or in m/p_T bins, picked the former, as there is no resonance in t-channel.

Conclusions:

1. The substructure becomes less two-pronged with visible and dark hadrons in them, and the absence of the dark hadrons create the two-pronged structure \rightarrow The substructure is created by the interspersing of visible hadrons with dark hadrons.
2. Specific hidden valley parameter configurations can reduce the dark shower model dependent features of the signal jets.

Cards utilised by Snowmass White paper... [link](#)

Summary

- Several avenues of strongly interacting dark sector open for exploration
- General idea evolving around the need of more signature based searches
- Can probe unusual collider phase-space corners by exploiting existing wealth of jet substructure observables
- **First bounds set on these kind of signatures in the t -channel production mode from ATLAS & CMS (many more to come)**



BACKUP

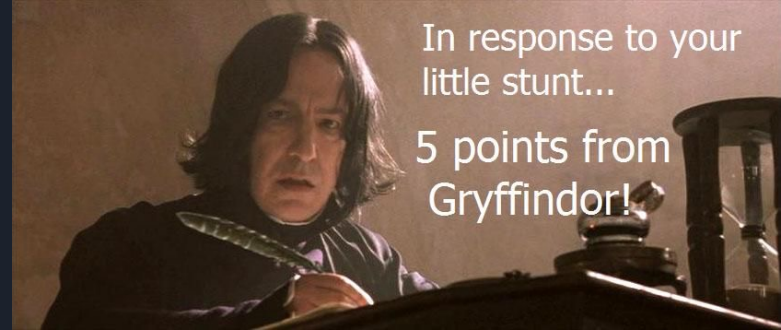


SVJ t-channel systematic uncertainties

- Largest contribution from theoretical components ($\sim 25\%$ on signal cross-sections mostly from scale variations).
 - Apart from usual scale and PDF variations, also included $t\bar{t}$ and single top I/FSR variation, ME and PS variation by using alternate generators, DR/DS subtraction scheme difference for tW .
 - W +jets split into heavy and light flavour, and an extra 30% normalisation uncertainty was used for heavy flavour, since Sherpa 2.2 has been found to underestimate V +heavy-flavour by about a factor of 1.3
 - There is known mismodelling in multijet processes, so a data-otherMC vs multijet reweighting is done in $250 < MET < 300$ GeV in 9bin distribution \rightarrow the reweighting factors are obtained in bin 3,6,9, and applied to 1-3, 4-6, 7-9 respectively.
- Standard experimental uncertainties: JES/JER, MET soft term, luminosity, PU reweighting, flavour tagging, reconstruction/identification/isolation/trigger efficiencies on muon and tau leptons.

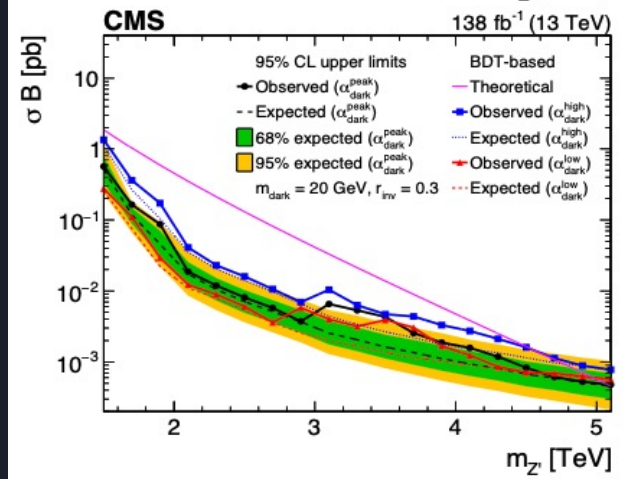
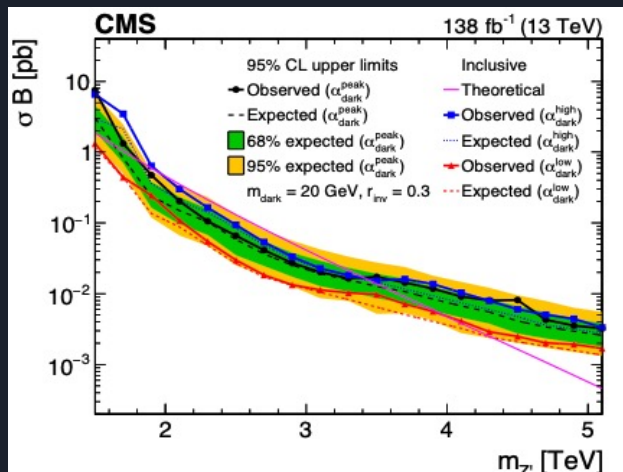


SVJ t-channel fit Strategy



- Maximum likelihood function fit performed using the product of all relevant Gaussian and Poisson PDFs and 9-bin yields, using MC templates.
- The signal region (0L SR) 9-binned histograms are fitted simultaneously with 1LCR, 1L1BCR and 2LCR.
- Dedicated systematic uncertainties are applied to the 0L SR, 1L CR, 1L1BCR & 2L CR 9-binned histogram.

The largest post-fit effects on the shape are signal modelling uncertainties up to 8%, Z+jets modelling uncertainties up to 7%, and top process modelling uncertainties up to 4%. The rest of the contributions are less than 2%.

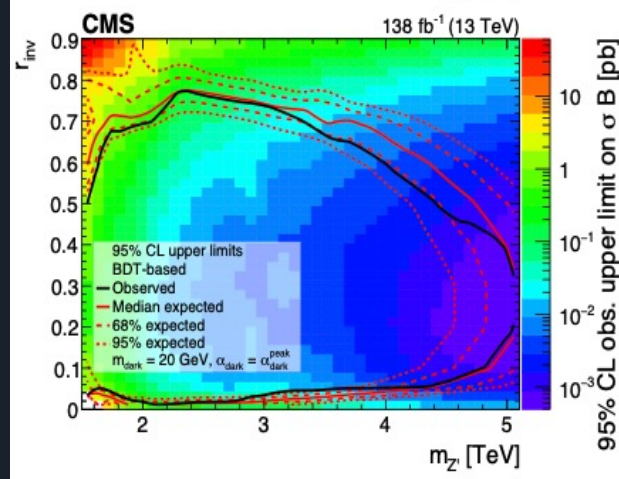
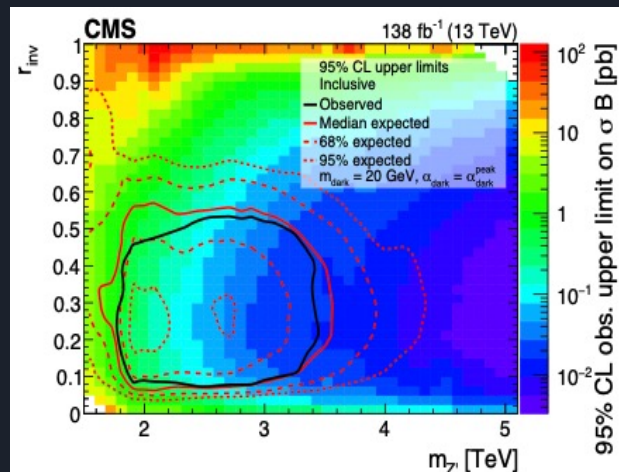


Inclusive signal regions

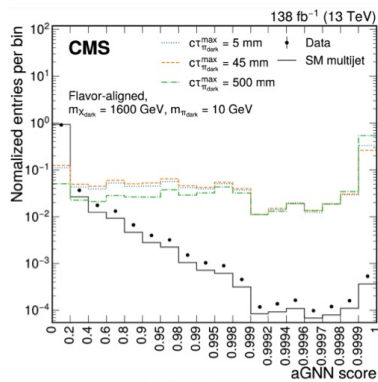
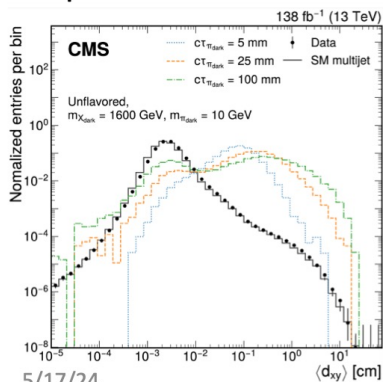
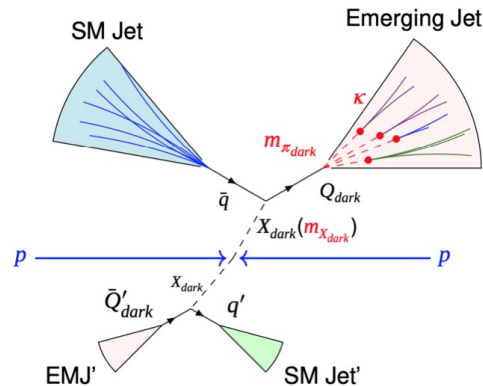
Excludes models with
 $1.5 < M_Z < 4.0 \text{ TeV}$ and
 $0.07 < R_{\text{inv}} < 0.53$

BDT - based regions

Excludes models with
 $1.5 < M_Z < 5.1 \text{ TeV}$ and
 $0.01 < R_{\text{inv}} < 0.77$
 For $M_{\text{dark}} = 20 \text{ GeV}$



- Signal signatures on detector:
 - Semi-long-lived dark mesons decay back to SM particles, forming **SM showers emerging from vertices** finite distances away from collision points;
 - Tree level two SM jets + 2 emerging jets;
- Cut & Count analysis
- Background estimated using control samples in data based on jet misidentification probabilities.

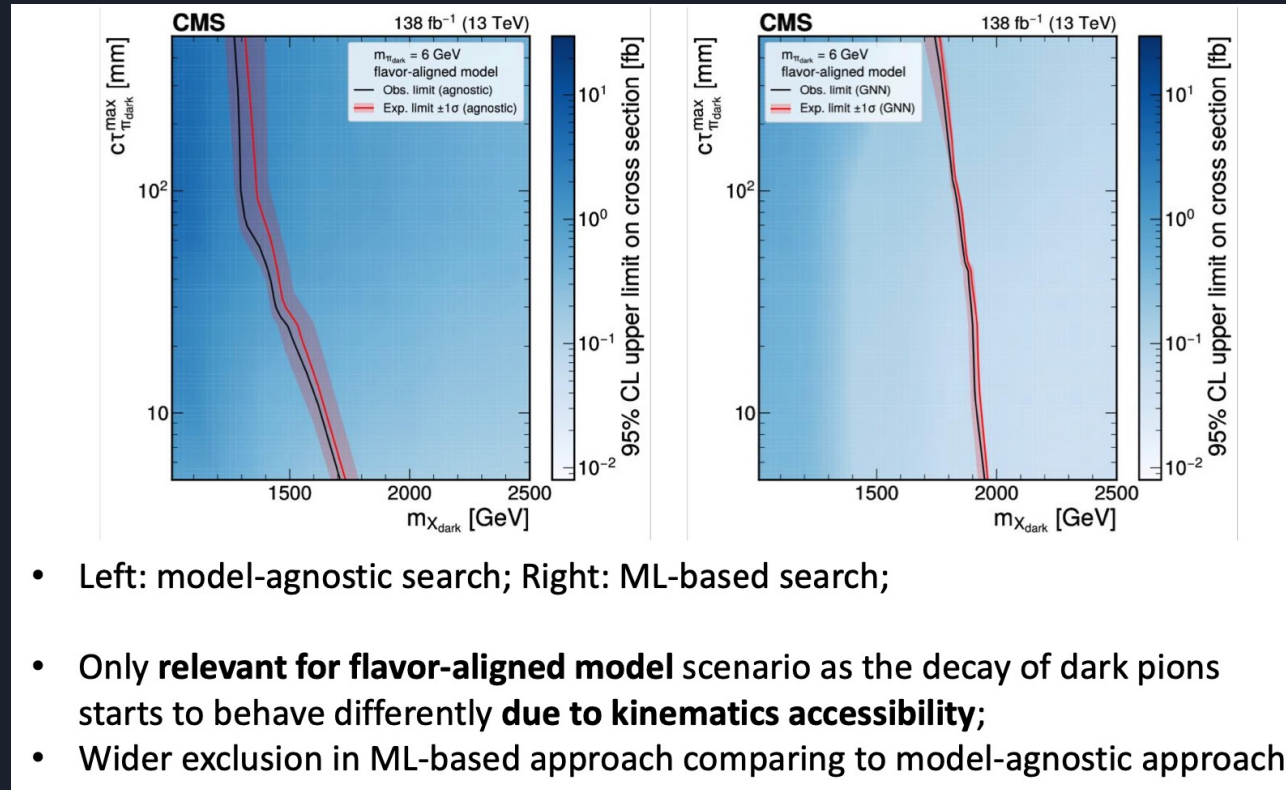


Two approaches:

- **Model-agnostic approach** for reinterpretability;
- **ML-based approach** to maximize sensitivity for the chosen models under search.

CMS EJ SEARCH RESULTS

From L. Wang's [talk](#)



2HDM+a

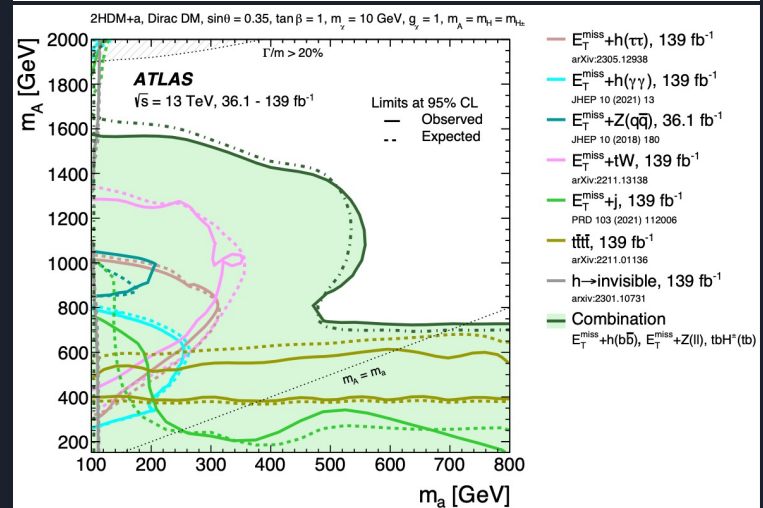
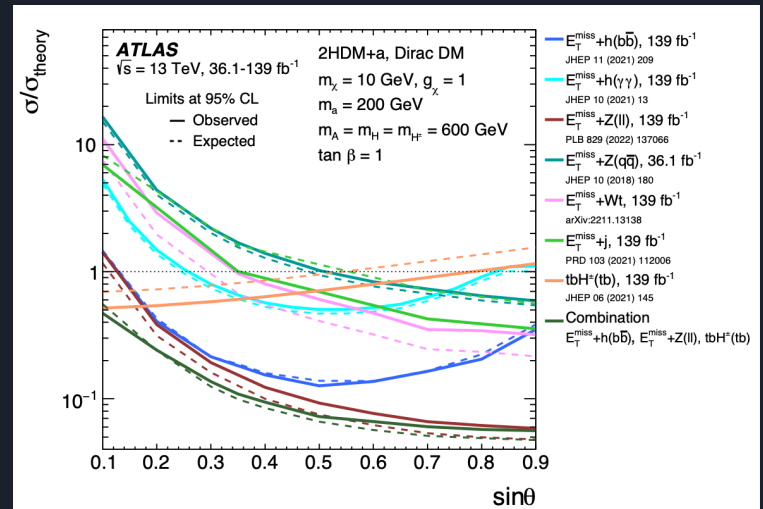
LHC DMWG report: [ArXiv 1810.09420](https://arxiv.org/abs/1810.09420)

2HDM containing an additional pseudoscalar boson which mediates the interactions between the visible and the dark sector

- gauge invariant & renormalisable extension of simplified pseudoscalar model
- DM candidate: singlet under SM gauge group, usually a Dirac fermion
- CP-odd mediator [pseudo scalar to bypass constraints from DD]

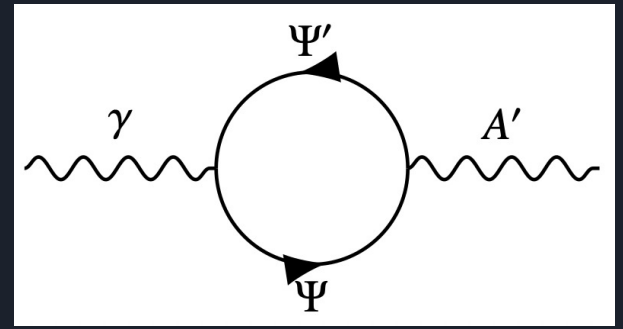
FREE PARAMETERS OF THE THEORY:

- * masses of the heavy Higgs ($m_A = m_{H^\pm} = m_{H^\pm}$)
- * mass of pseudo-scalar mediator, m_a
- * mass of DM particle, m_χ
- * sine of mixing angle b/w CP-odd states a & A , $\sin\theta$
- * \sqrt{EV} ratio, $\tan\beta$



ATLAS 2HDM+a summary paper
<https://arxiv.org/abs/2306.00641>

Dark-photons

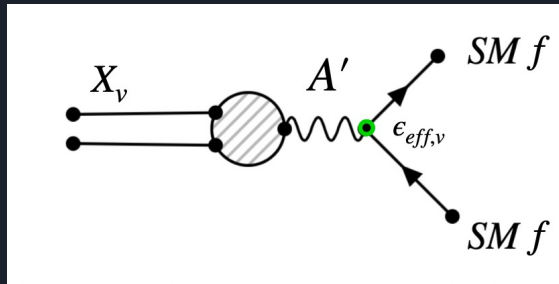


Vector Portal: Add a $U(1)'$ whose massive "dark" gauge boson mixes kinetically with SM photon

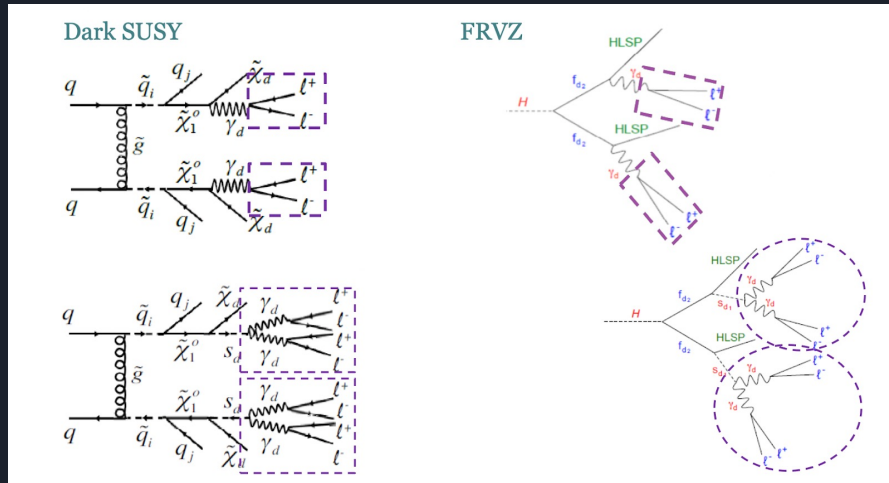
Higgs Portal: Add dark scalar singlet that spontaneously breaks $U(1)'$ and mixes with SM Higgs

Hidden Valley: sector of dark particles, interacting amongst themselves

- Lowest particle in Valley forced to decay to SM due to mass gap or symmetry
- "Portal" coupling both to SM and HV operators, can be A'



Benchmark models for limit setting



Falkowsky-Ruderman-Volansky-Zupan model:
 Pair of dark fermions produced in the Higgs boson decay

dark fermion decays in turn to a dark photon + a lighter dark fermion assumed to be the Hidden Lightest Stable Particle (HLSP).

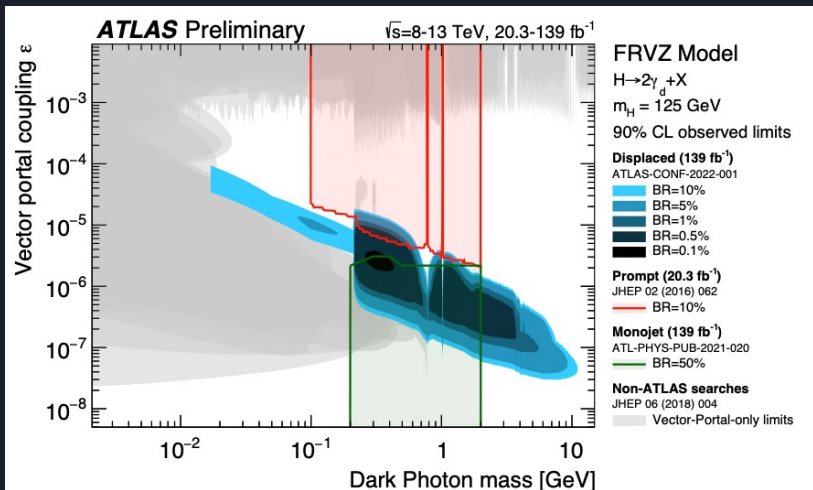
dark photon (vector mediator) mixes kinetically with the SM photon and decays to leptons or light hadrons.

Dark SUSY:

Neutralino \rightarrow dark photon and susy DM, and dark photon decaying to pair of Leptons

Neutralino \rightarrow susy DM, and pair of dark photons decaying to pair of Leptons

Summary of constraints from colliders



[ATL-PHYS-PUB-2022-007](#)

Exclusions also available for lepton+jets in high mass regions from ATLAS, CMS and LHCb dedicated searches.

