

The QTNM collaboration: a project for neutrino mass measurement

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(on behalf of the QTNM collaboration)

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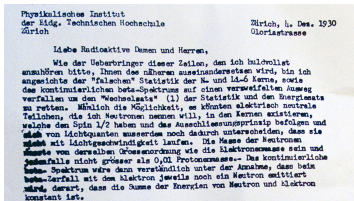
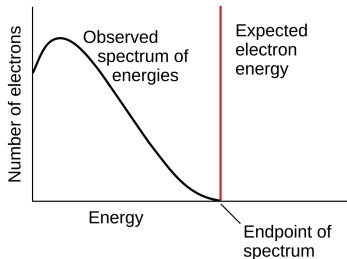


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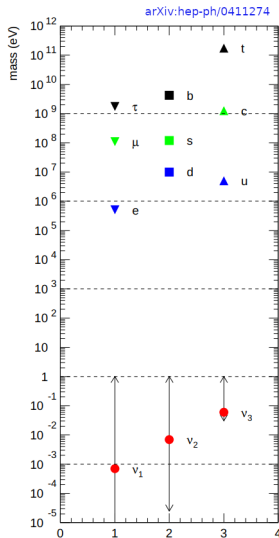
The neutrino

- Existence first postulated by Pauli in 1930 to explain shape of β decay spectrum
- Directly detected by Cowan & Reines in 1956
- Three flavours discovered: ν_e , ν_μ , ν_τ . All appeared to be massless.



Neutrino oscillations

- Evidence from atmospheric, solar, reactor and accelerator neutrinos all confirms the existence of **neutrino oscillations**
- 2015 Nobel Prize awarded to Takaaki Kajita & Arthur B. Macdonald *“for the discovery of neutrino oscillations, which shows that neutrinos have mass”*
- Oscillations arise from mixing between flavour and mass eigenstates of neutrinos
- **Neutrino mass scale very different** from other fermions



Neutrino oscillations and neutrino mixing

- Mixing between flavour and mass eigenstates given by

$$|\nu_i\rangle = \sum_{\alpha} U_{\alpha i} |\nu_{\alpha}\rangle$$

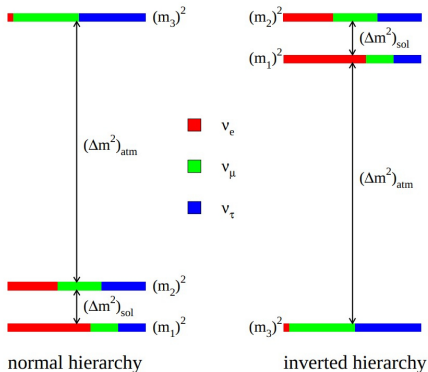
where

$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix}$$

is a unitary matrix

- Oscillations controlled by the **matrix U** and the squared differences between the mass eigenstates, $\Delta m_{ij}^2 = m_i^2 - m_j^2$
- These Δm_{ij}^2 control the **length/energy scale** at which oscillations occur

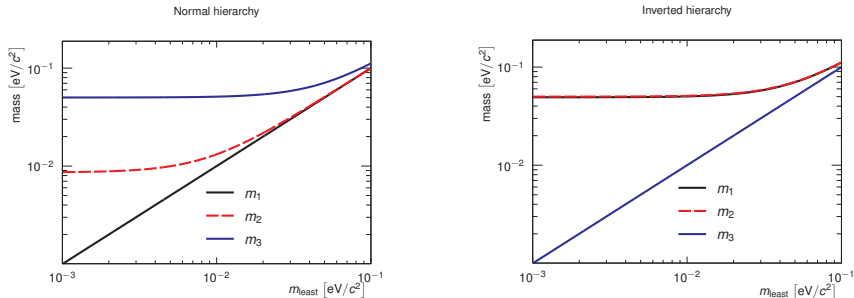
Neutrino mass hierarchy



[arXiv:1310.4340](https://arxiv.org/abs/1310.4340)

- Differences between m_i^2 known from oscillations
- **Ordering** of mass eigenstates currently **unknown**
- Lightest mass eigenstate is either m_1 (normal hierarchy) or m_3 (inverted hierarchy)

Possible neutrino masses



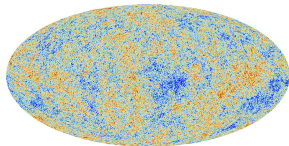
- It is possible the lightest mass eigenstate (either m_1 or m_3) may in fact be massless
- Masses of the other eigenstates are then constrained by the mass splittings

Why measure the neutrino mass?

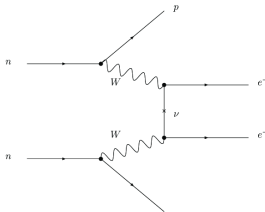
- One of the most abundant particles in the universe – and we don't know its mass!
- Very different mass scale suggests different mass generation mechanism (compared to just Higgs)
- Connected to various other areas of physics:
 - Lepton number violation
 - Cosmology
 - Sterile neutrinos

Measuring the neutrino mass

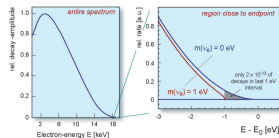
Cosmological measurements



Neutrinoless double β -decay

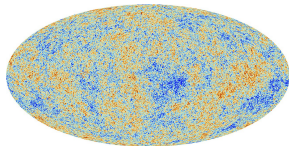


Direct measurement of β -decay



Measuring the neutrino mass

Cosmological measurements

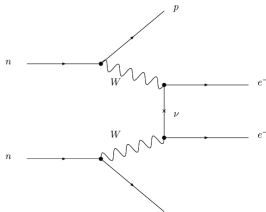


$$\Sigma = \sum_i m_i$$

$$\Sigma < 0.111 \text{ eV}c^{-2}$$

arXiv:2007.08991 [astro-ph.CO] (2021)

Neutrinoless double β -decay

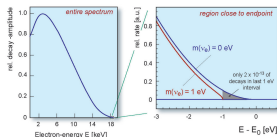


$$m_{\beta\beta} = \sum_i (U_{ei})^2 m_i$$

$$|m_{\beta\beta}| < 0.036 - 0.156 \text{ eV}c^{-2}$$

arXiv:2203.02139 [hep-ex] (2022)

Direct measurement of β -decay

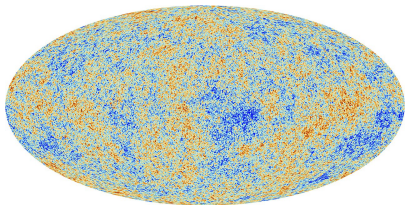


$$m_{\beta} = \sqrt{\sum_i |U_{ei}|^2 m_i^2}$$

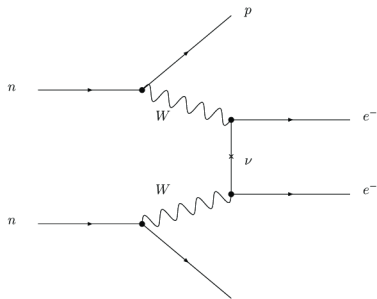
$$m_{\beta} < 0.8 \text{ eV}c^{-2}$$

Nat. Phys. 18, 160-166 (2022)

The first two have their issues



Relies on cosmological models



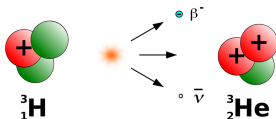
Only works if neutrinos are Majorana particles

- Neither of these are **model-independent** measurements in the same way that direct measurement is

Measurements of β -decay

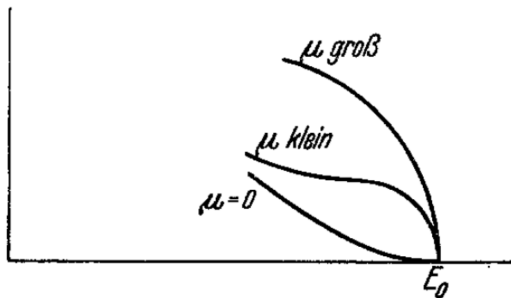


- For β -decay the total energy of the initial state is well known and the kinematics of the **final state** can be **precisely measured**
- Can use energy and momentum conservation to constrain the neutrino mass
- Processes such as this often referred to as **'direct measurement'**
- Isotope commonly used is **tritium**



Direct measurement

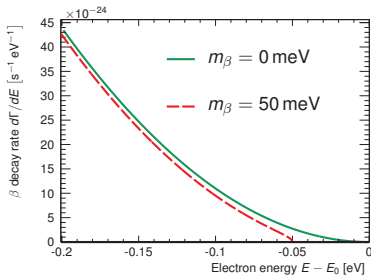
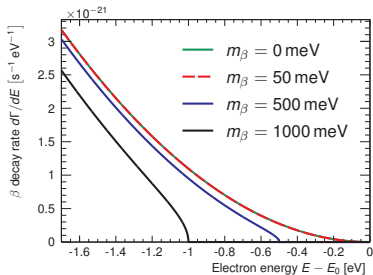
Z. Phys. **88**, 161 (1934)



- An old idea – Fermi suggested the shape of the β -ray spectrum could be used to determine the neutrino mass in 1934, as did Perrin separately in 1933

Tritium β -decay spectrum

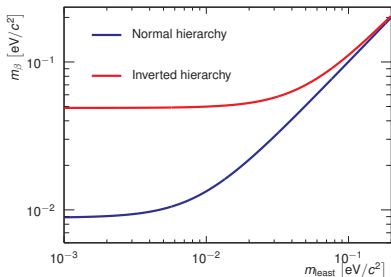
$$\frac{d\Gamma}{dE} \approx 3r_0 (E_0 - E) \left[(E_0 - E)^2 - m_\beta^2 \right]^{1/2} \Theta(E_0 - E - m_\beta)$$



- Information about m_β is encoded in last eV of energy spectrum (endpoint energy, $E_0 \approx 18.6$ keV)

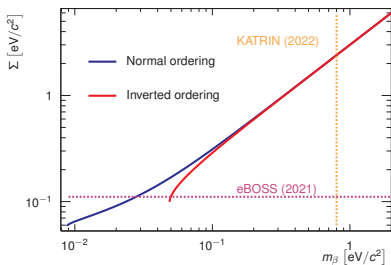
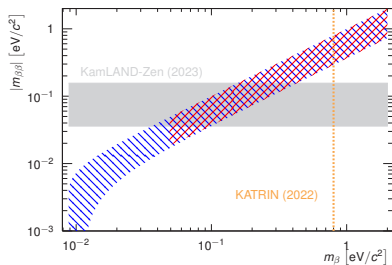
Limits on m_β

$$m_\beta = \sqrt{\sum_i |U_{ei}|^2 m_i^2}$$



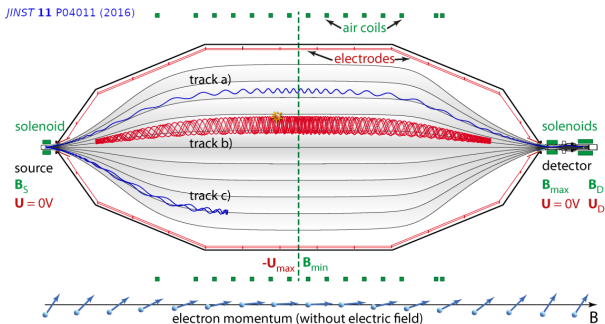
- Mass splittings from oscillation experiments provide a lower limit on m_β
 - For normal hierarchy $m_\beta \gtrsim 9 \text{ meV}$
 - For inverted hierarchy $m_\beta \gtrsim 50 \text{ meV}$
- An experiment with a sensitivity below $m_\beta \approx 50 \text{ meV}$ will determine the **mass hierarchy** (if still unknown)
- A sensitivity of 9 meV gives us a **guaranteed discovery**

What can measuring m_β tell us?



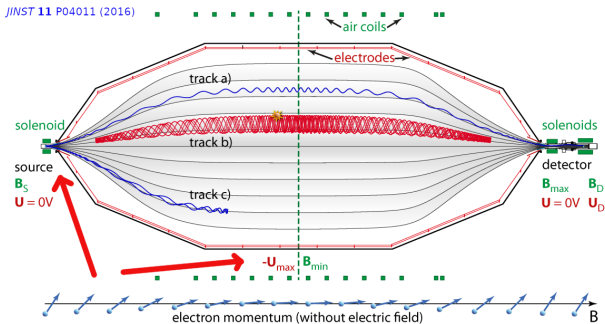
- Results can augment those from $0\nu\beta\beta$ experiments and cosmology
- Red here is IH, blue is NH

MAC-E filter – Current state of the art



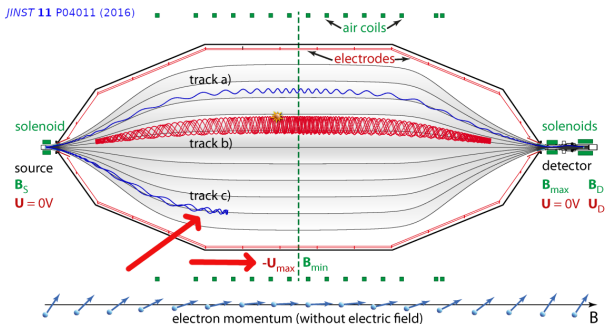
■ Magnetic Adiabatic Collimation – Electrostatic

MAC-E filter – Current state of the art



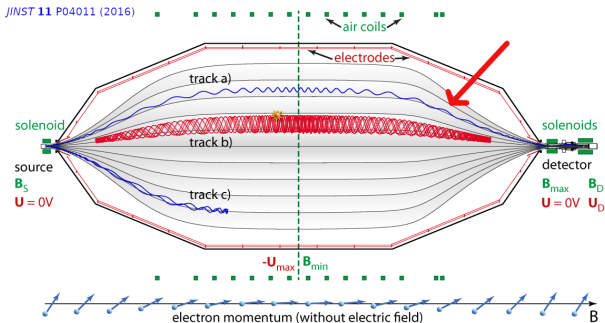
- **Magnetic Adiabatic Collimation – Electrostatic**
- Electrons emitted in source region with high magnetic field, B_S , and travel adiabatically along field lines to analysing region with much lower field, B_{min}

MAC-E filter – Current state of the art



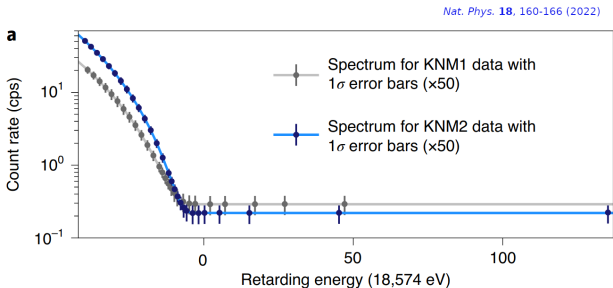
- **Magnetic Adiabatic Collimation – Electrostatic**
- Retarding potential at central analysing plane prevents electrons without sufficient energy from passing

MAC-E filter – Current state of the art



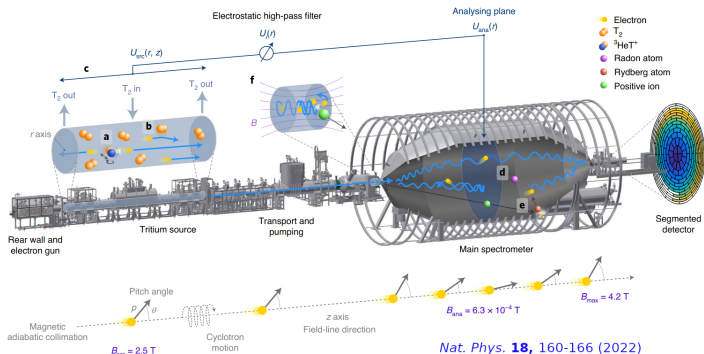
- **Magnetic Adiabatic Collimation – Electrostatic**
- Those electrons with sufficient energy to pass the potential barrier are re-accelerated and detected

MAC-E filter – Current state of the art



- Repeat this for different retarding potentials in order to generate spectrum

KATRIN experiment



- Current best limits on m_β are produced by the KATRIN experiment
 – $m_\beta < 0.8 \text{ eV}/c^2$
- Expected final sensitivity of $0.2 \text{ eV}/c^2$
- 70 m long beamline, spectrometer is 9.8 m in diameter and 23.3 m in length held at pressure of 10^{-11} mbar

Limitations of MAC-E filters

- To increase statistical power, can increase source size
- However, source thickness is limited by $\sigma n \leq 1$ to avoid collisional losses
- For a MAC-E filter:

$$\begin{aligned}\frac{\Delta E}{E} &= \frac{B_{\text{ana}}}{B_{\text{src}}} \\ &= \left(\frac{R_{\text{src}}}{R_{\text{ana}}} \right)^2\end{aligned}$$

- Therefore, increasing R_{src} requires a corresponding **increase** in the **spectrometer size**

Limitations of MAC-E filters

- Impractical to scale KATRIN up

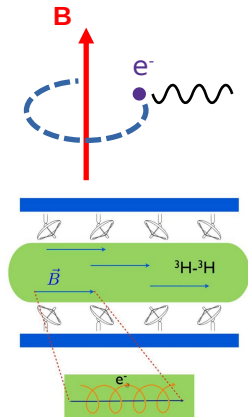


- We require a different technique for $m_\beta < 0.2 \text{ eV}/c^2$

CRES overview

- **C**yclotron **R**adiation **E**mission Spectroscopy
- Concept pioneered by Project 8 collaboration^a
- β -decay electrons immersed in B-field emit **EM radiation** – frequency depends only on **electron energy** and **B-field strength**
- $E_{\text{kin}} = Q_{\beta} = 18.6 \text{ keV}$, $B = 1 \text{ T}$
- $f = 27 \text{ GHz}$, $\lambda \sim 1 \text{ cm}$, MW radiation
- Radiation collected with antenna, waveguide or resonant cavity

$$f = \frac{1}{2\pi} \frac{eB}{m_e + E_{\text{kin}}/c^2}$$



^aMonreal, B.; Formaggio, J. A. *Phys. Rev. D* **2009**, *80*.

CRES advantages

- Frequency measurements can reach precision of $\Delta f/f \sim 10^{-6}$



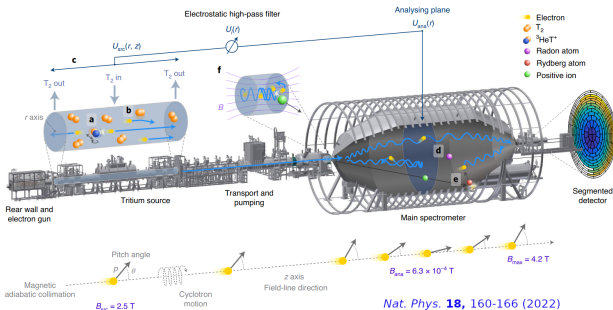
A. L. Schawlow

“Never measure anything but
frequency”

CRES advantages

- Frequency measurements can reach precision of $\Delta f/f \sim 10^{-6}$
- **No losses** while transporting e^- from the source to the detector

Unlike MAC-E filters, no need to transport electrons from source to detector – fewer electrons lost to scattering



CRES advantages

- Frequency measurements can reach precision of $\Delta f/f \sim 10^{-6}$
- No losses while transporting e^- from the source to the detector
- The source (tritium gas) is **transparent** to MW radiation

CRES advantages

- Frequency measurements can reach precision of $\Delta f/f \sim 10^{-6}$
- No losses while transporting e^- from the source to the detector
- The source (tritium gas) is transparent to MW radiation
- **Differential spectrum** measurements

- MAC-E filters are ‘integral filters’ – intensity above a point in the spectrum is counted
- Extra time required for measuring background, spectrum intensity, endpoint energy
- Differential spectrometers do not have this issue

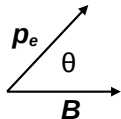
CRES challenges

- Radiated **powers** are very **small**

$$\text{Radiated power} \approx \frac{2\pi e^2 f_0^2}{3\epsilon_0 c} \frac{\beta^2 \sin^2 \theta}{1 - \beta^2}$$

where $f_0 = \frac{1}{2\pi} \frac{eB}{m_e} = 27.9925 \text{ GHz}$ for $B = 1 \text{ T}$

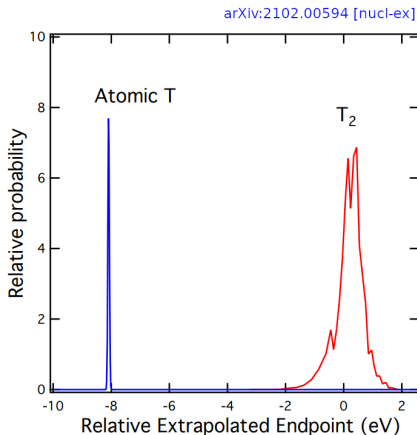
$P = 1.17 \text{ fW}$ for $B = 1 \text{ T}$ and $\theta = \frac{\pi}{2}$



CRES challenges

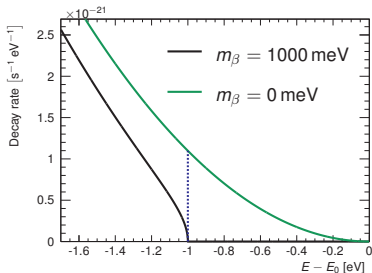
- Radiated powers are very small
- **Atomic tritium** source required

- Molecular tritium has rotational and vibrational excitations that broaden the endpoint peak
- Production and preservation of atomic tritium is a key challenge for any future experiment

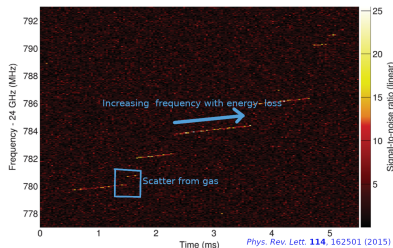


CRES challenges

- Radiated powers are very small
 - Atomic tritium source required
 - Need to trap and observe $\sim 10^{20}$ tritium atoms for \sim year
-
- Last eV of the spectrum contains 2.9×10^{-13} of the events
 - Necessitates an intense source

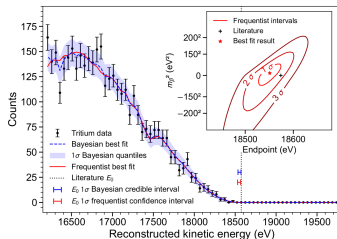


Project 8



- Above: CRES signal from 30 keV ^{83m}Kr decay electrons in Phase I

- Phase II with molecular tritium in a ~ 1 T field
- Detected 3742 events over 82 days



Ashtari Esfahani et al. arXiv:2303.12055 [nucl-ex]

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Quantum Technologies for Neutrino Mass Collaboration

- **Proposal goal: build a demonstrator apparatus for determining neutrino mass via CRES from tritium β -decay – CRESDA**
- This entails:
 - Demonstration of confinement of atoms with **densities of $10^{12} - 10^{13} \text{cm}^{-3}$** , scalable to ultimate exposure of **$\sim 10^{20}$ T atoms**
 - Magnetic field mapping with **$< 1 \mu\text{T}$ absolute precision** and **$\sim 1 \text{mm}$ spatial resolution**
 - Observation of CRES electrons (non-tritium source), potentially with **quantum noise limited** MW detection systems
 - Experiment to be built at **University College London**

Potential project pathway



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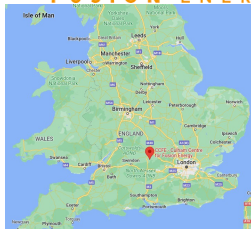


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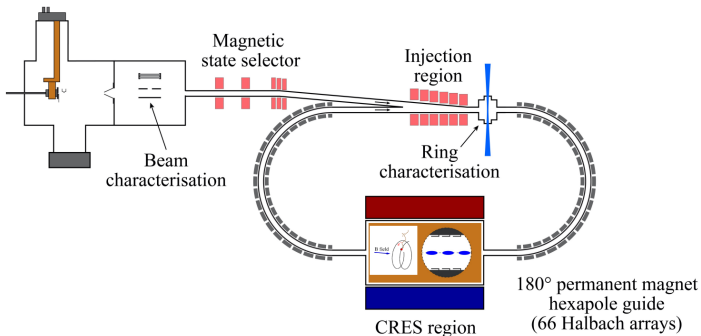
Quantum Technologies for Neutrino Mass Collaboration

- Basic technology demonstration at UCL (2021–2025)
- Tritium demonstrations at Culham (2025–2030+)
- Final neutrino mass experiment with $\sim 10 - 50\text{meV}$ sensitivity at Culham or similar facility (2030–2040)



CRESDA outline

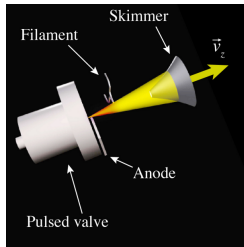
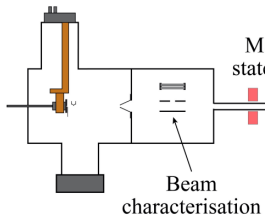
H/D/T atom supersonic beam discharge source (30 K)



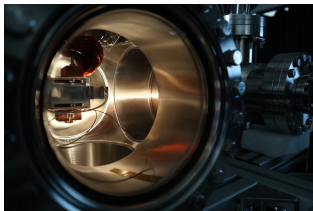
- Consists of source, atomic trapping storage ring and instrumented CRES region

Atomic source

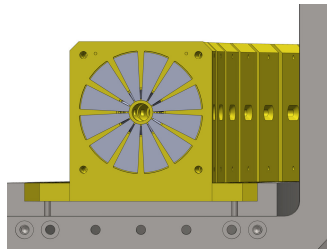
- When measuring the endpoint of atomic tritium any **molecular tritium** contamination acts as a **background**



- Cryogenic (30 K) pulsed supersonic source
- $\text{H}_2/\text{D}_2/\text{T}_2$ dissociation using DC discharge seeded with electron from tungsten filament

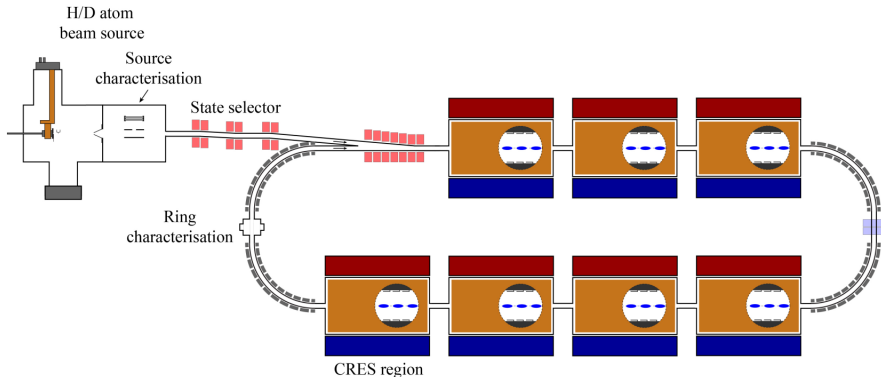


Storage ring



- Need to confine neutral atoms and prevent them forming molecules on surfaces
- Advantages of storage ring design include:
 - **Lower atomic losses** than occur when loading a trap (no deceleration and cooling required)
 - **Separates magnetic field requirements** for optimal high-frequency resolution CRES from magnetic trapping
 - **Scalability** (see next slide...)

Scalability of storage ring concept



- One option: multiple CRES modules connected by single storage ring

Electron trapping

$$f = \frac{1}{2\pi} \frac{eB}{m_e + E_k/c^2} = \frac{eB}{2\pi} \left(\frac{E_{\text{tot}}}{c^2} \right)^{-1}$$

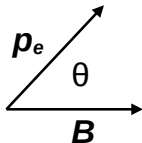
$$\begin{aligned} \frac{df}{dE_{\text{tot}}} &= -\frac{eBc^2}{2\pi} (E_{\text{tot}})^{-2} \\ &= 51 \text{ kHz eV}^{-1} \end{aligned}$$

for $B = 1 \text{ T}$, $E_k = 18.6 \text{ keV}$

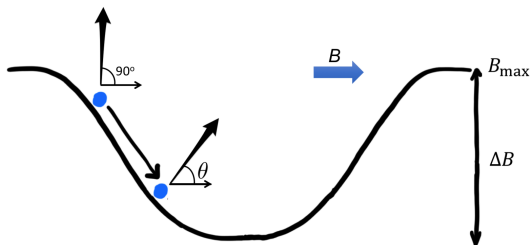
\therefore for $\Delta E = 1 \text{ eV}$, we require $\Delta f \approx 50 \text{ kHz}$

$$t_{\text{obs}} \sim \frac{1}{\Delta f} \quad \therefore t_{\text{obs}} \geq 20 \mu\text{s}$$

An endpoint electron travelling at 89° to B -field will travel 275 m parallel to B -field in this time



Electron trapping

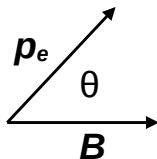
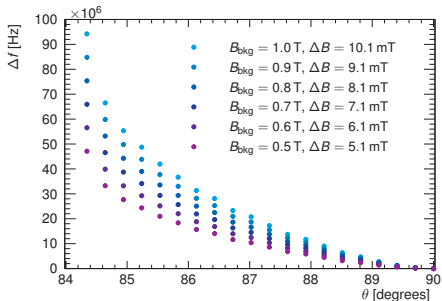


$$\theta \geq \sin^{-1} \left(\sqrt{1 - \frac{\Delta B}{B_{\max}}} \right)$$

Phys. Rev. C **99**, 055501 (2019)

- **Solution:** Trap β -decay electrons in a 'no-work' trap where they can be continuously observed for 10s or 100s of μ s
- **Local minimum** in magnitude of background B -field
- Require trapping field of order **1 mT** against **1 T background**
- Trap design has large effect on range of observed frequencies – key to understand this

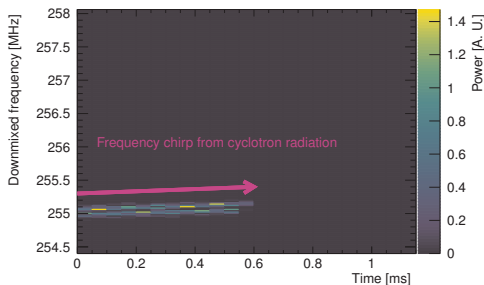
Importance of electron trap design



- Electrons with pitch angle, $\theta < 90^\circ$ experience different average B -field as a result of motion in trap
- Deeper magnetic traps contain more electrons but these will radiate at different frequencies (for the same energy)
- Other features of the signal must be used to get best energy resolution

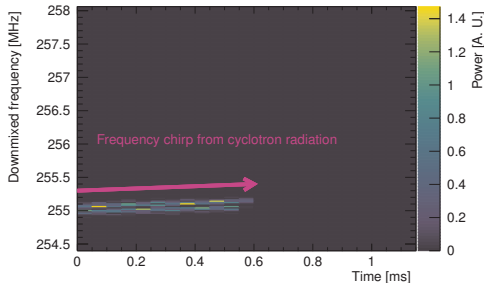
Simulation & analysis

- Work ongoing to optimise trap and RF collection design
- Can simulate an (idealised) decay electron signal in a variety of traps using **custom software**
- Simulation includes electron propagation, RF collection and basic model of signal processing



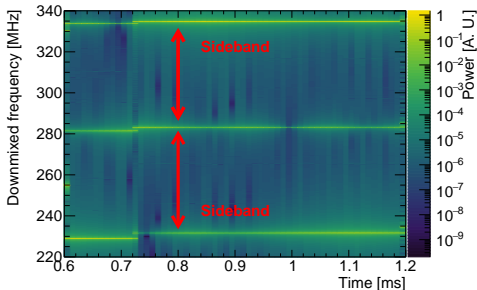
Simulation & analysis

- Many interesting CRES signal features that can be used to constrain **electron energies**
- Chirp rate of our signal tells us about the rate of energy loss



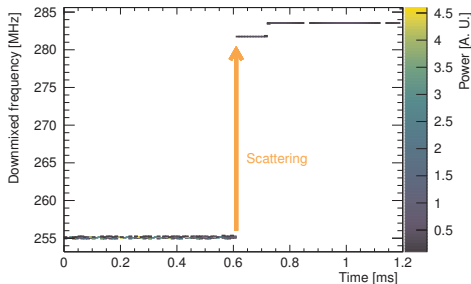
Simulation & analysis

- Many interesting CRES signal features that can be used to constrain **electron energies**
- Chirp rate of our signal tells us about the rate of energy loss
- Sidebands caused by AM and FM of signal – tell us about axial motion of electron

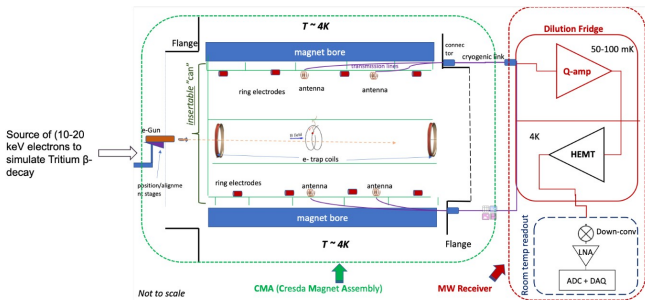


Simulation & analysis

- Many interesting CRES signal features that can be used to constrain **electron energies**
- Chirp rate of our signal tells us about the rate of energy loss
- Sidebands caused by AM and FM of signal – tell us about axial motion of electron
- Discrete jumps in frequency caused by scattering from residual gas – need to fully understand frequency of these scatters in order to reconstruct original frequency



CRES Magnet Assembly (CMA)

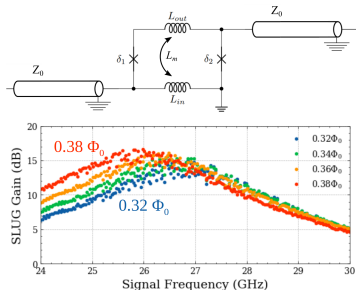
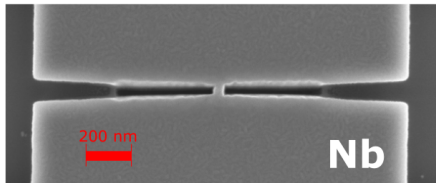


- Aim to detect CRES at **cryogenic temperatures** (minimise noise)
- Exploring use of **quantum-limited amplifiers** to overcome small signal power
- Possible **electron source** to calibrate energy

Quantum-limited microwave amplifiers

- CRES signal is very weak (~ 1 fW)
- State-of-the-art HEMT amplifiers have noise temperatures of about 7 K
- We require **quantum-limited** amplifiers:
 - Superconducting **L**ow-Inductance **U**ndulatory **G**alvanometer (SLUG)
 - Superconducting parametric amplifiers

SLUG amplifiers

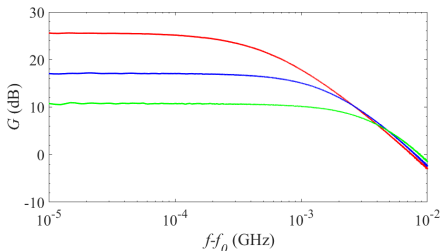


- Superconducting **L**ow-inductance **U**ndulatory **G**alvanometer
- Non-parametric low-noise cryogenic amplifier under development at NPL
- Nb nanobridge junctions fabricated (above left) to allow operation at **high frequencies**
- Numerical simulations and experiments ongoing to **characterise** and **tune** SLUGs for our use

Superconducting parametric amplifiers



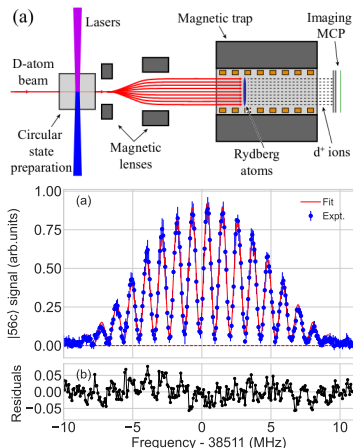
~ 1 cm



- Members at Cambridge Quantum Sensors group have designed, fabricated and tested **high-gain parametric amplifiers** based on **superconducting resonators**
- **Superconducting NbN** optimised as **amplifier material**
- Right: Power gain vs frequency – high gains with a bandwidth of several MHz. Distribution is two-sided about central RF frequency
- Noise measurement of amplifiers ongoing and making good progress

Magnetometry

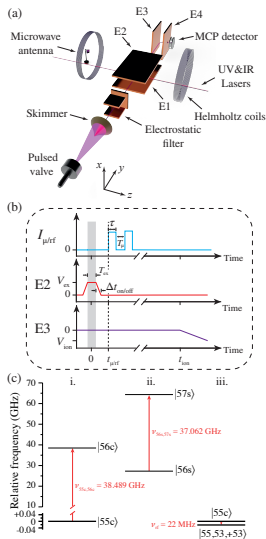
- Measuring electron energy with resolution of 10^{-6} requires that B -field be known to **similar level**
- Deuterium or tritium atoms can be used as **quantum sensors** for B -field mapping to better than $1 \mu\text{T}$
- **Circular Rydberg states** prepared in beam and passed through CRES volume
- Pulses of MW radiation drive Rydberg-Rydberg transitions – sensitive to B -field
- Potential **spatial resolution** of 1 mm



Magnetometry

- Current results:
 - Absolute field precision of $\pm 2 \mu\text{T}$, relative $\pm 900 \text{ nT}$
 - Spatial resolution of $\pm 0.87 \text{ mm}$
 - Electrometry precision of $85 \mu\text{V cm}^{-1}$
- Paper detailing method and results published this year^a

^aZou, J.; Hogan, S. D. *Phys. Rev. A* **2023**, *107*.



Summary

- The neutrino mass scale remains unknown but the answer has the potential to provide key constraints in several areas
- Current measurement techniques (MAC-E filters) are at their limits and cannot take us to an experiment with guaranteed discovery potential
- CRES is a recent technique that allows the measurement of electron energy at unprecedented precision
- The QTNM collaboration is building on unique quantum techniques to demonstrate the viability of a CRES experiment
- Provides the exciting possibility of having the ultimate neutrino mass experiment in the UK



Thanks for listening!



UNIVERSITY OF CAMBRIDGE



Swansea University
Prifysgol Abertawe



Backup

Are calorimetric techniques the solution?

- Can embed isotopes in microcalorimeters – decay via **electron capture**
- Calorimeter measures the atomic de-excitation energy (minus the neutrino's energy)
- ^{163}Ho used by ECHo and HOLMES collabs.
 - In order to get required energy resolution for a high enough activity may **require thousands or 100s of thousands of microcalorimeters** operated at mK level

