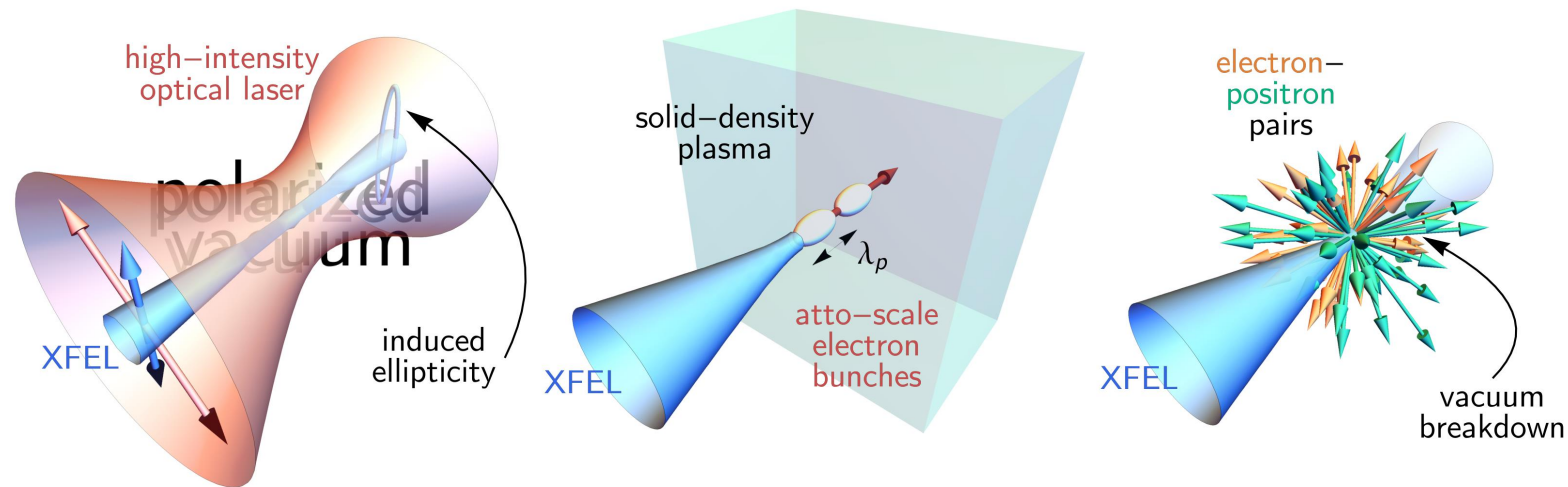


Opportunities in nonlinear classical and quantum physics with intense XFELs



Tom Blackburn

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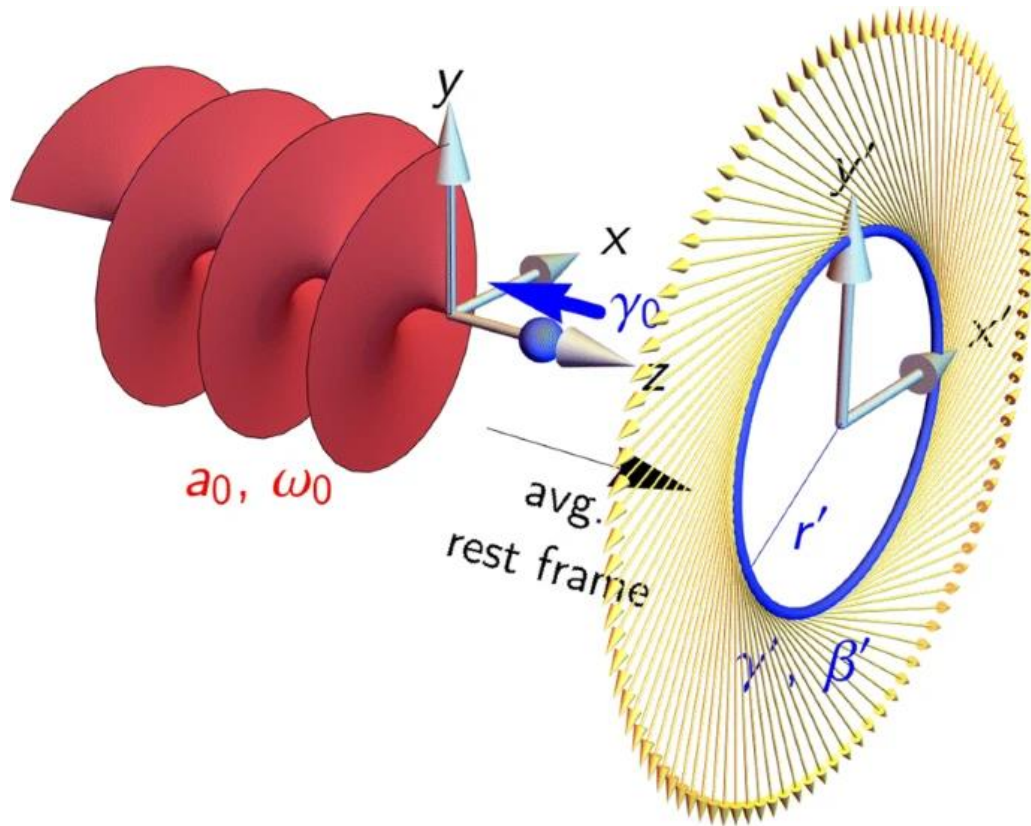
19th January 2023

UK XFEL Townhall – Fundamental Physics,
Quantum Computing and AI

Opportunities in nonlinear classical and quantum physics with intense XFELS

Summary

- Introduction: key parameters, the nonlinear regimes
- XFELs as probes: vacuum birefringence
- XFELs as backgrounds: gamma-ray sources for nuclear photonics
- XFELs as drivers: wakefields in solid-density plasmas...
- ... and Schwinger pair creation in vacuum



- Compare the characteristic frequency of the emitted radiation to the orbital (cyclotron) frequency:

$$\frac{\omega'}{\omega_c} \approx a_0^3$$

- Harmonic order of the emitted radiation, or number of participating photons.

$$a_0 = \frac{eE \frac{\hbar}{mc}}{\hbar\omega}$$

force exerted by electric field

Compton length

photon energy

- Alternatively...
- Electron does not interact with one photon of the radiation field, but many.
- Determines the importance of higher order terms in a perturbative expansion, which only works if a_0 is much smaller than one.

force exerted by electric field

Compton length

$$? = \frac{eE \frac{\hbar}{mc}}{mc^2}$$

electron rest energy

- Alternatively...
- Electron does not interact with one photon of the radiation field, but many.
- Determines the importance of higher order terms in a perturbative expansion, which only works if a_0 is much smaller than one.
- If we swap out the photon energy...

Quantum nonlinearity parameter $\chi_{e,\gamma}$

$$\chi = \frac{eE \frac{\hbar}{mc}}{mc^2}$$

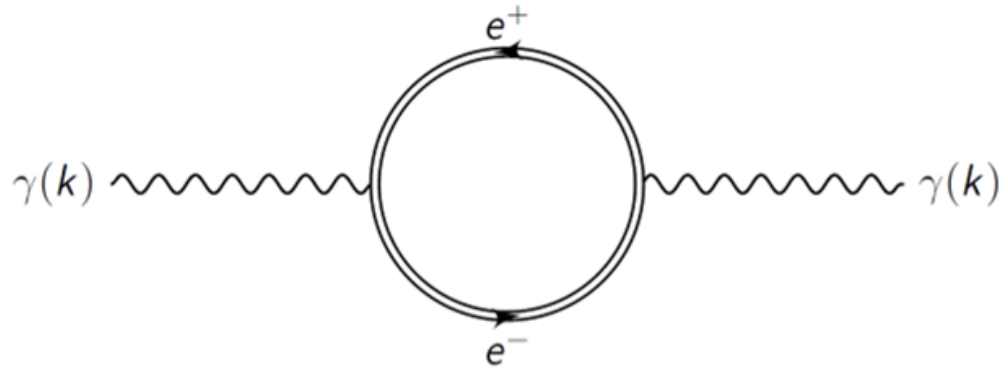
force exerted by electric field

Compton length

electron rest energy

- ... we see if the field does enough work over a Compton length to create an electron,
- or the field strength in the particle rest frame, in units of the critical (Schwinger) field.
- Then quantum effects must be important – even if the field amplitude in the lab frame is far below critical.

photon interacts with virtual electron-positron pair of polarized vacuum



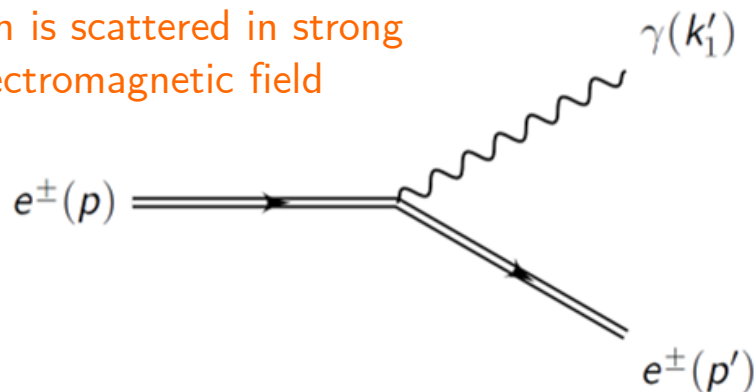
- For a photon, vacuum polarization and electron-positron pair creation are controlled by

$$\chi_\gamma = \frac{e\sqrt{-(F_{\mu\nu}k^\nu)^2}}{m^3} = \frac{\omega|\vec{E}_\perp + \vec{n} \times \vec{B}|}{mE_{\text{crit}}}$$

- For an electron, photon emission (and recoil effects) are controlled by

$$\chi_e = \frac{e\sqrt{-(F_{\mu\nu}p^\nu)^2}}{m^3} \simeq \frac{\gamma|\vec{E}_\perp + \vec{v} \times \vec{B}|}{E_{\text{crit}}}$$

electron is scattered in strong electromagnetic field



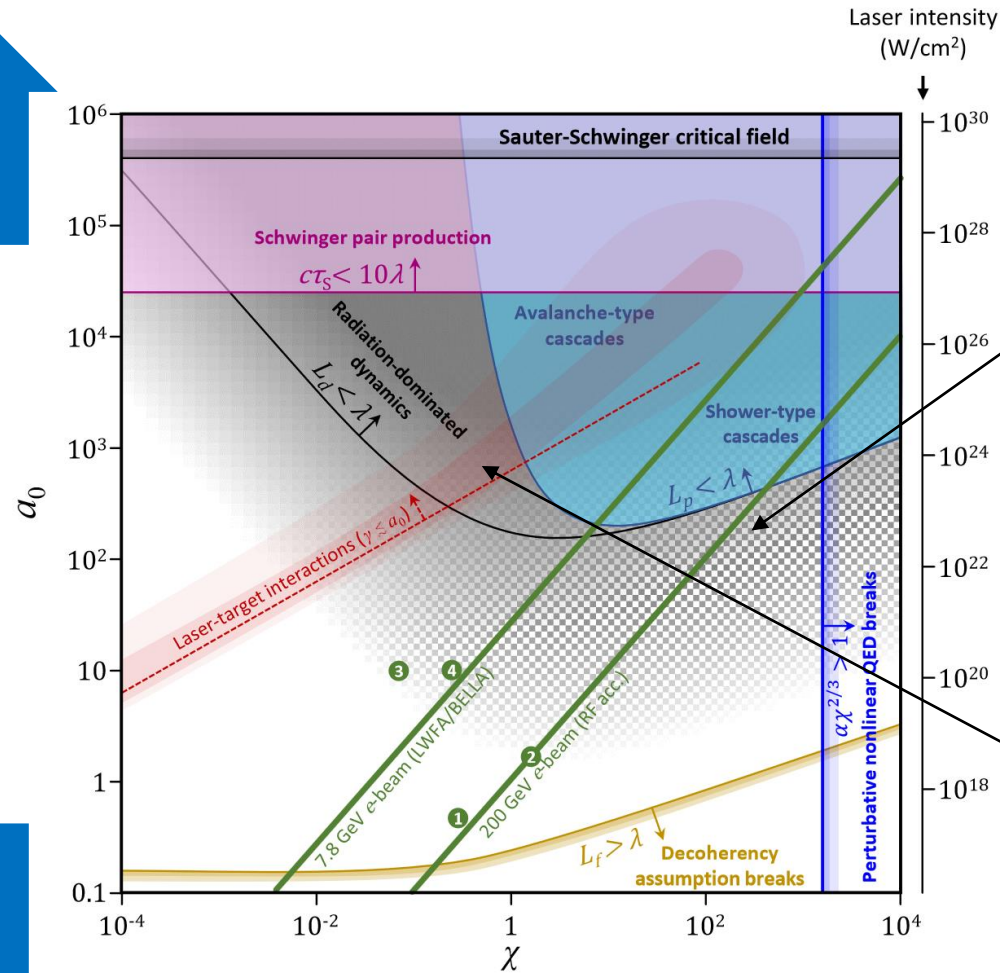
Introduction

Parameter space

high intensity,
quasistatic,
tunnelling,
field can be treated as
being locally constant



low intensity,
weakly multiphoton,
use perturbation
theory



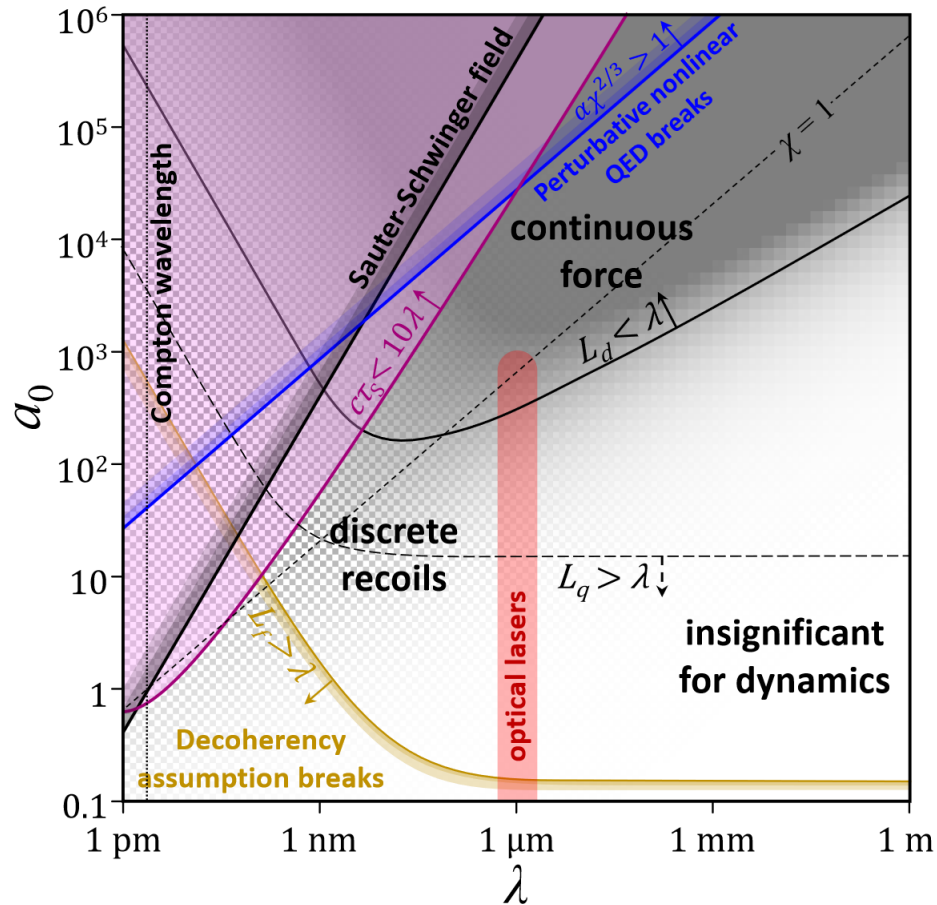
optical laser +
high-energy
electron beam

optical laser +
matter

Gonoskov et al, RMP 2022

Introduction

Wavelength matters



Gonoskov et al, RMP 2022

- Which of nonlinear classical or quantum effects kicks in first depends on the wavelength of the strong field (among other things).
- Most research has focused on high-power optical lasers.
- In this talk I will present some ideas, all of which rely upon the co-location of intense XFELs with other particle and radiation sources.

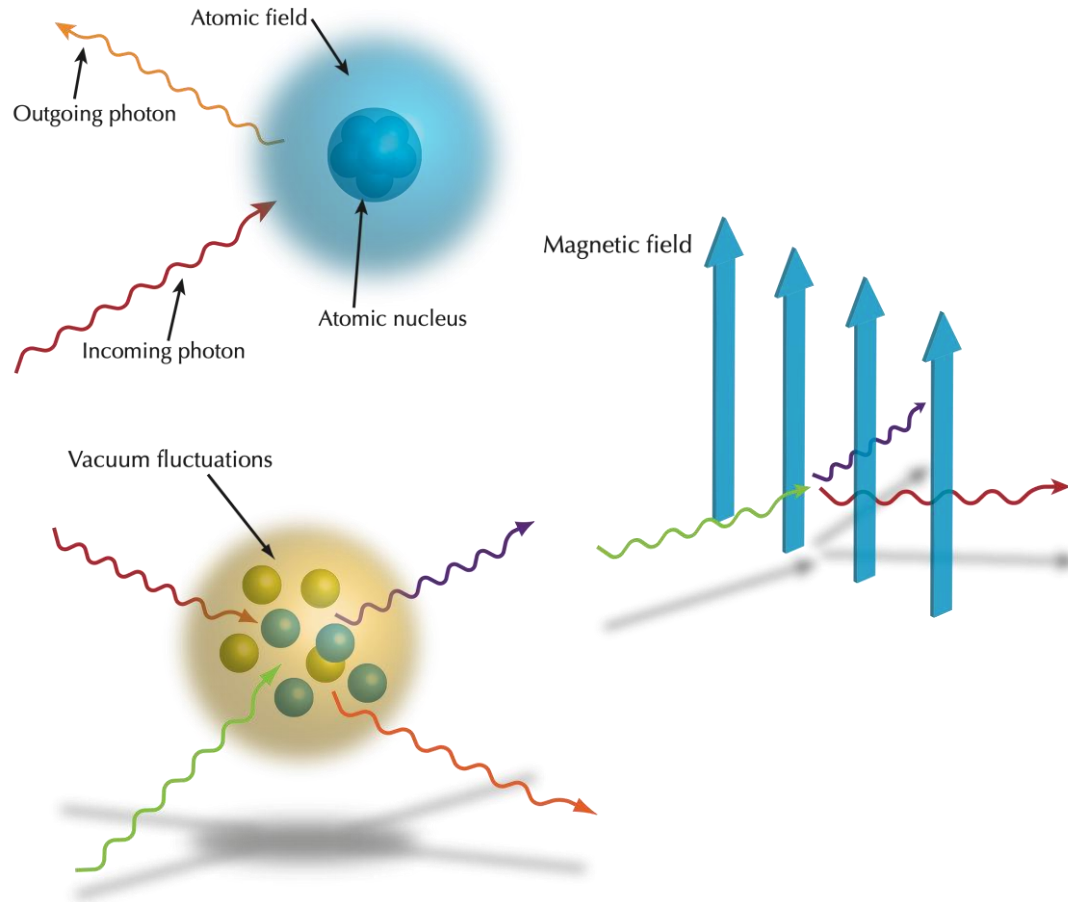
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XFELs as probes

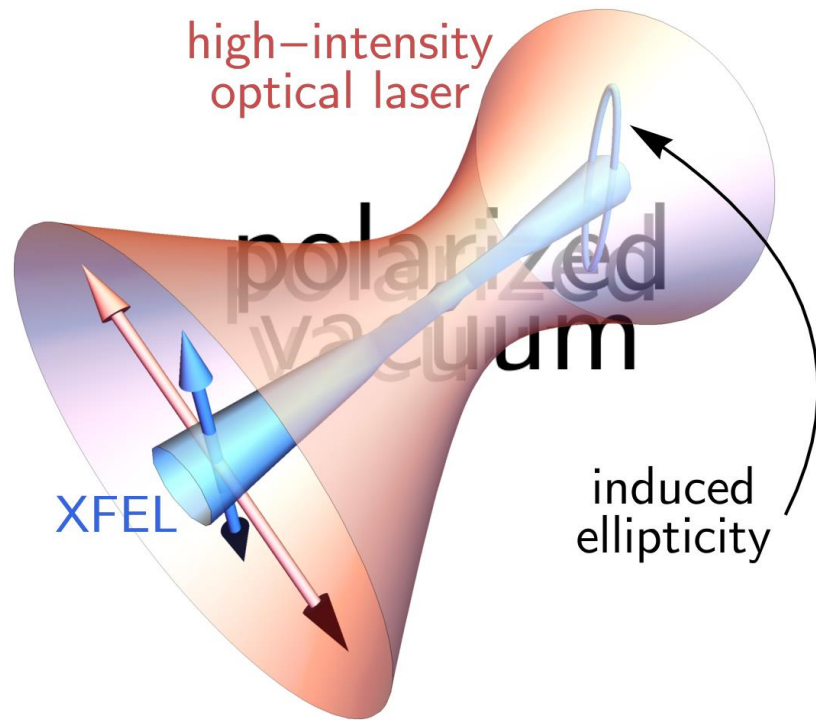
The quantum vacuum



- Photons can interact with each other via fluctuating electron-positron pairs.
- Astrophysical applications; laboratory tests of strong-field QED.
- Many of these cross sections scale positively with increased frequency.
- Top to bottom: Delbruck scattering, photon splitting, four-wave mixing.

XFELs as probes

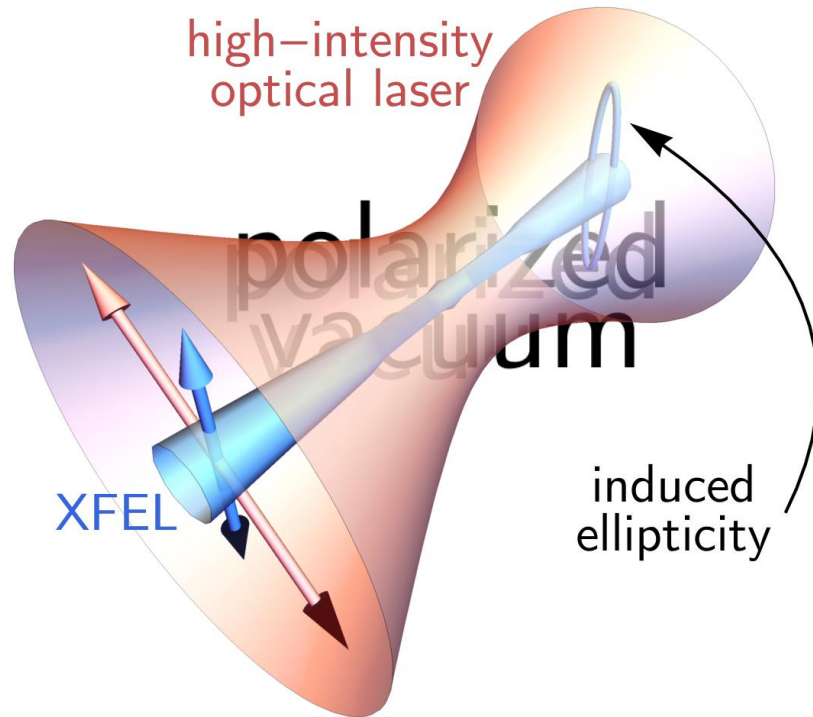
Vacuum birefringence



- Cross section for four-photon scattering $\sigma [\text{cm}^{-2}] = 10^{-31} (\omega/m)^6$ ($\omega =$ photon energy in ZMF, natural units hereon).
- 11 orders of magnitude larger for XFEL-laser than laser-laser (33 for XFEL-XFEL vs laser-laser).
- Use optical laser as “vacuum polarizer” (coherence!) and XFEL as probe [HIBEF, SEL etc].

XFELs as probes

Vacuum birefringence



- Scattering visible in helicity change of the probe X rays, as if the vacuum refractive index were polarization dependent.
- Initially linearly polarized probe acquires a small ellipticity $\delta = \omega_X L \Delta n / 2$, where $\Delta n = 2\alpha I_{\text{laser}} / I_{\text{cr}}$.
- Why X rays? High polarization purity and precision of measurement.

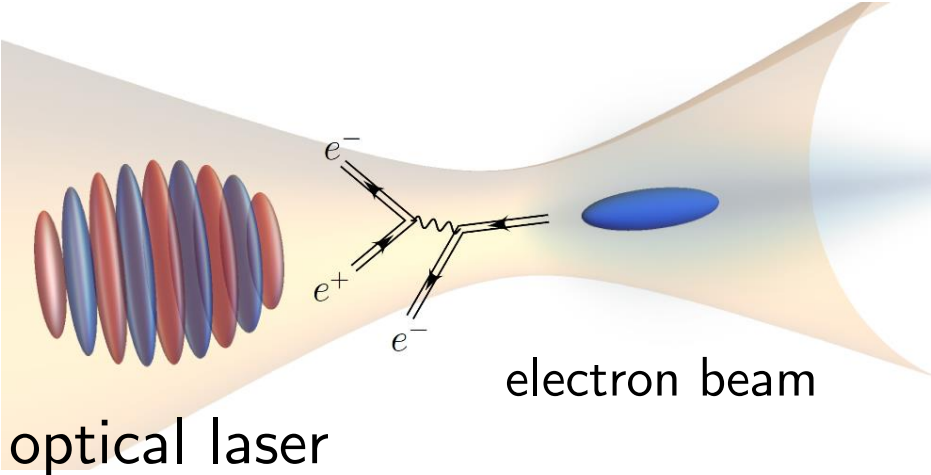
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XFELs as backgrounds

Gamma-ray source



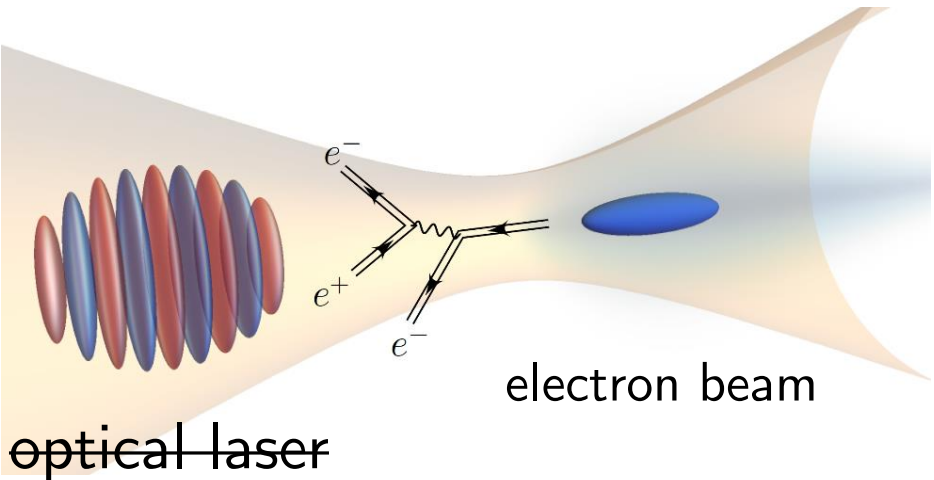
$$R = f \frac{\sigma_T N_L N_e}{\pi(r_L^2 + r_e^2)}$$

photons/s $\leftarrow R$
 BX/s $\leftarrow f$
 no. laser photons $\leftarrow N_L$
 no. electrons $\leftarrow N_e$
 transverse sizes $\leftarrow r_L^2 + r_e^2$

- Compton scattering of electrons is a source of polarized, monochromatic photons.
- The 10 to 100 MeV range useful for nuclear photonics: excitation of resonances, transitions etc
- e.g. 720 MeV electrons + 1 eV optical laser (ELI-NP)
- or 1 GeV + 1 to 5 eV FEL (HIGS)

XFELs as backgrounds

Gamma-ray source



XFEL?

$$R = f \frac{\sigma_T N_L N_e}{\pi(r_L^2 + r_e^2)}$$

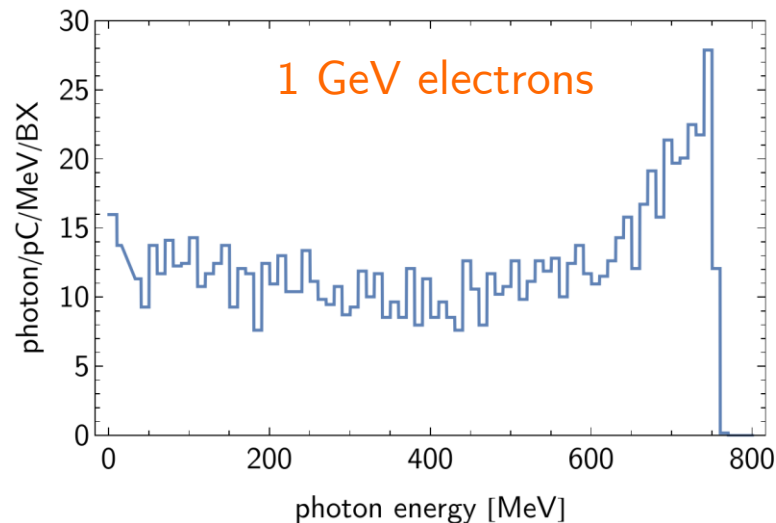
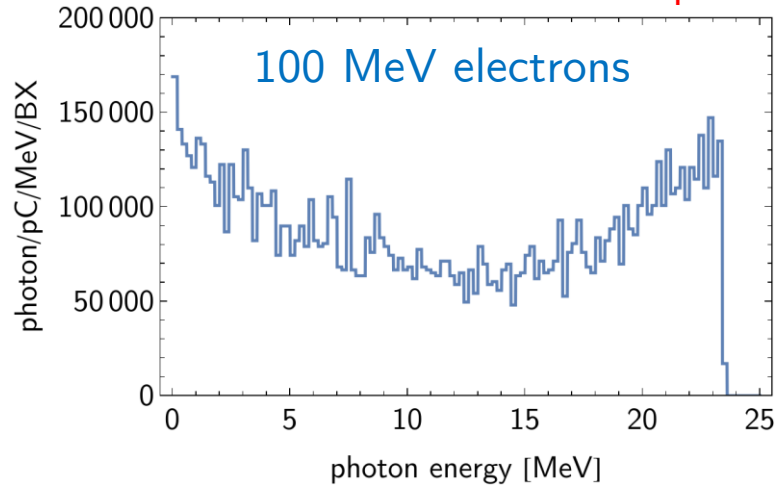
photons/s (pointing to R)
 BX/s (pointing to f)
 no. laser photons (pointing to N_L)
 no. electrons (pointing to N_e)
 transverse sizes (pointing to $r_L^2 + r_e^2$)

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XFELs as backgrounds

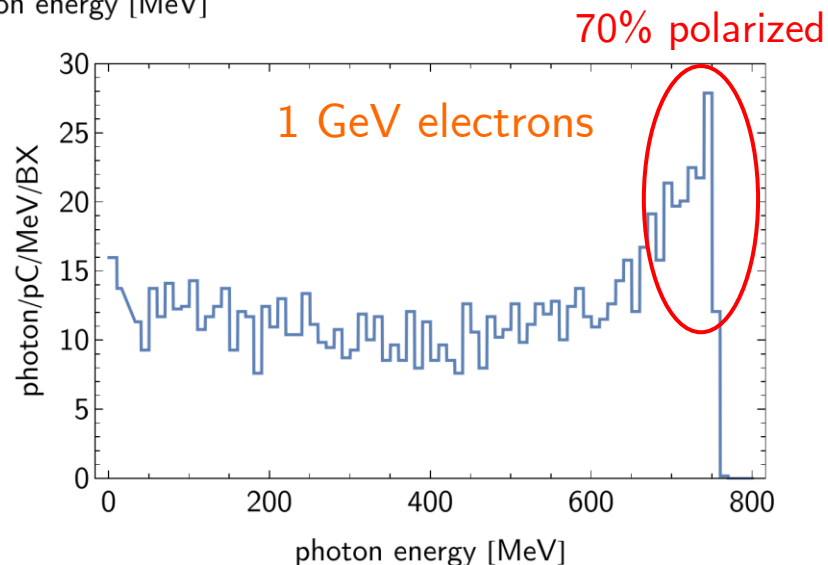
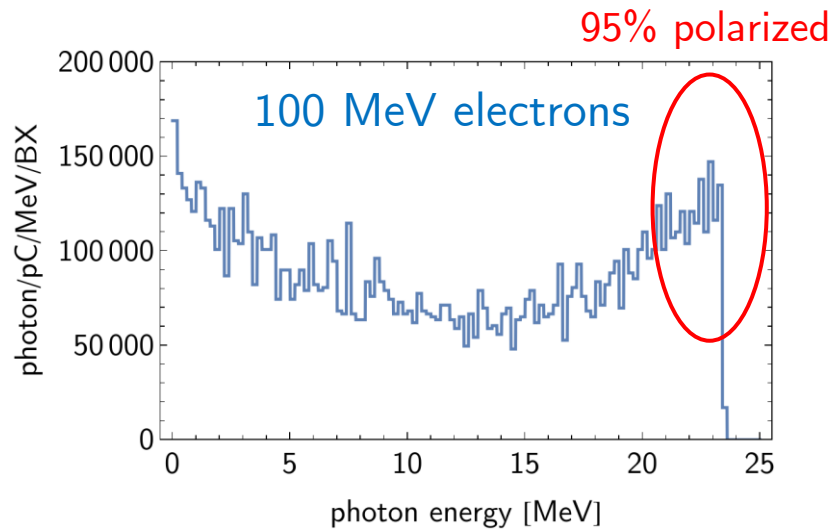
Gamma-ray source

95% polarized



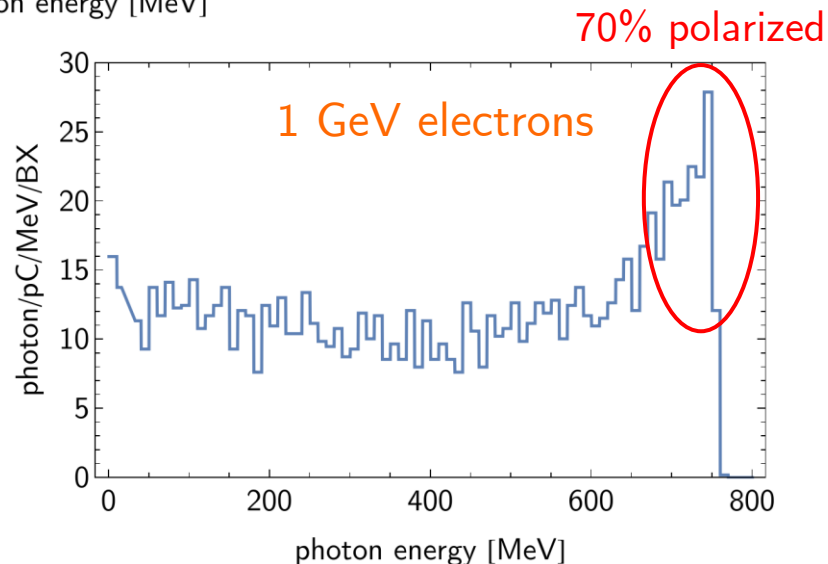
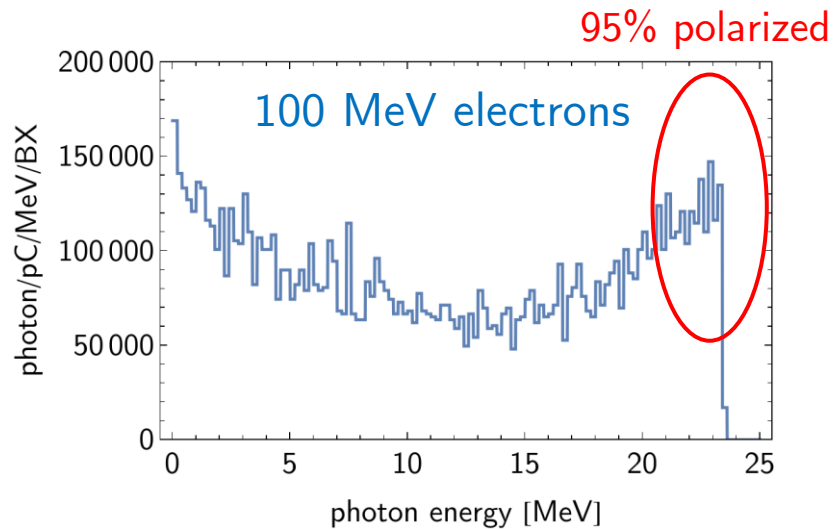
- To achieve the same photon energies by scattering XFEL light requires lower electron energies.
- Example: 100 MeV or 1 GeV electrons + 200 eV light @ 10^{19} W/cm² (300 cycles in duration)

XFELs as backgrounds Gamma-ray source

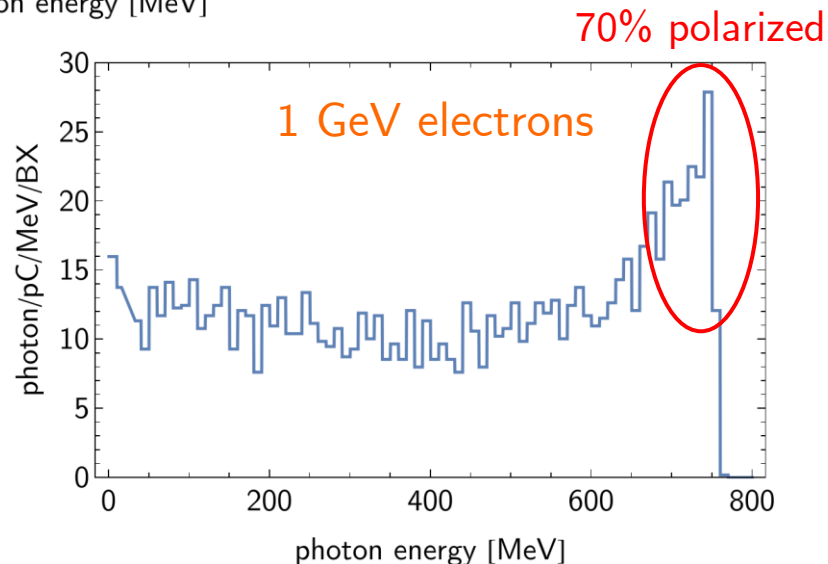
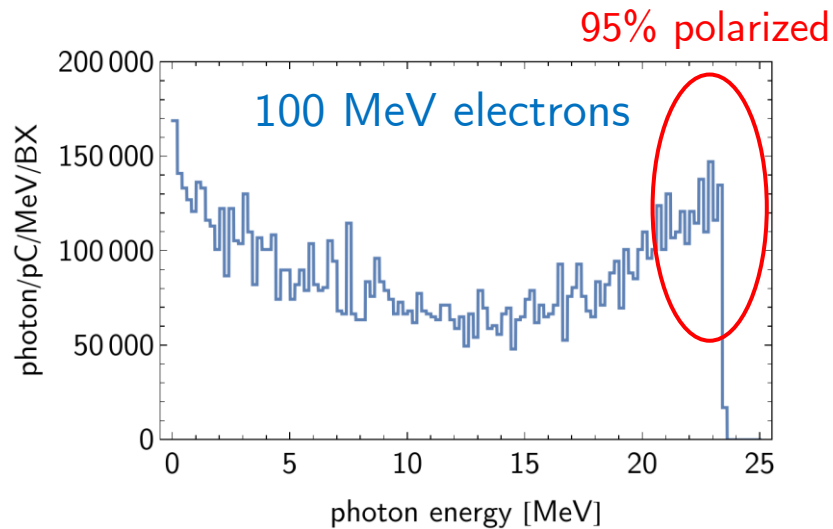


- To achieve the same photon energies by scattering XFEL light requires lower electron energies.
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- Gammas inherit polarization properties of the XFEL light: linear, circular etc
- Monochromatise by angle selection

XFELs as backgrounds Gamma-ray source



- How far can one go?
- Multi-GeV polarized gammas needed for hadron studies in strange sector / photoproduction of QCD exotics
- LEPS2 @ Spring-8 (ICS with 8-GeV electrons + VUV), GlueX (coherent bremsstrahlung)
- At 10 GeV, Compton edge at 9.7 GeV.



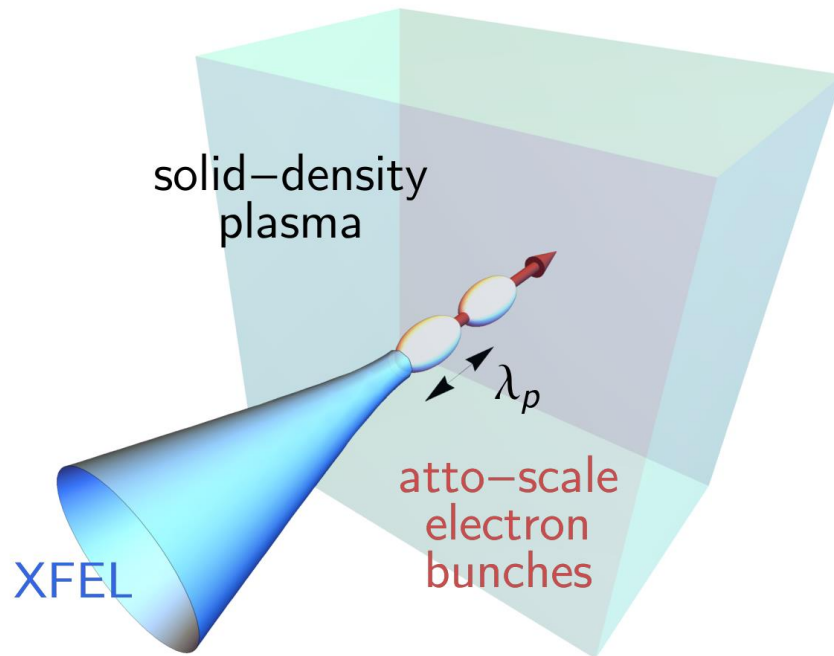
- Recoil effects on the electron beam are significant!
- Investigation of radiation reaction with optical lasers usually in the regime where single-photon recoil is much less important than the accumulated recoil from many emissions.
- Importance of “spin light” increased.

Opportunities in nonlinear classical and quantum physics with intense XFELS

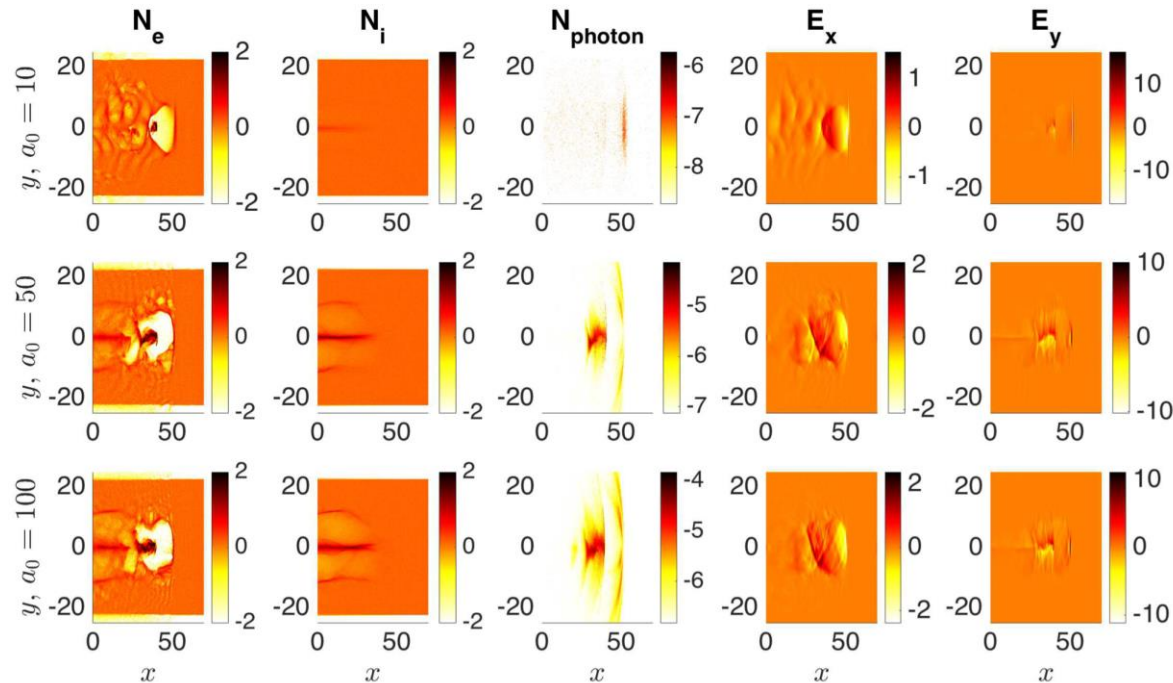
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XFELs as drivers XWFA



- Solid-density plasmas become underdense at XUV/X-ray wavelengths:
 $\lambda_p [\mu\text{m}] = n^{-1/2} [10^{21} \text{ cm}^{-3}]$
- As an example, a relativistically intense XUV pulse could drive wakefield acceleration of electrons [[Tajima EPJST 2014](#)].
- $a_0 > 1 \Rightarrow I > 10^{24} (\lambda_X [\text{nm}])^{-2} \text{ W/cm}^2$



Wettervik, PoP 2018

- Scaling wakefield acceleration to XUV wavelengths relies on similarity theory ($S = n_e/a_0 n_{cr}$).
- But not all physics scales with this parameter, e.g. radiation generation (betatron or nonlinear Compton) [Zhang PRAB 2016].
- PIC simulations show it holds to a wavelength of 1 nm.

Quantity	Scaling
Timescale	λ
Spatial dimensions of cavity	λ
Density	a_0/λ^2
Current density	a_0/λ^2
Current	a_0
Trapped charge	λa_0
Electron energy	a_0
Electromagnetic fields	a_0/λ
Pulse energy	λa_0^2
Ion motion	a_0
Quantum parameter, χ	a_0^2/λ
Radiated energy fraction	$\lambda^{1-\alpha} a_0^{2\alpha-1}$
Photon energy ($\chi \ll 1$)	a_0^3/λ
Photon energy ($\chi \sim 1$)	a_0
Time between emissions ($\chi \ll 1$)	T/a_0
Time between emissions ($\chi \sim 1$)	$T/(\lambda a_0)^{1/3}$

Wettervik, PoP 2018

- Motivation might be **accelerating gradient** (TeV/cm)...
- or perhaps the **current density** or **electron bunch duration**.
- For densities of 10^{23} to 10^{25} cm⁻³, cavity size implies an equivalent duration of 10^{-17} to 10^{-15} seconds.

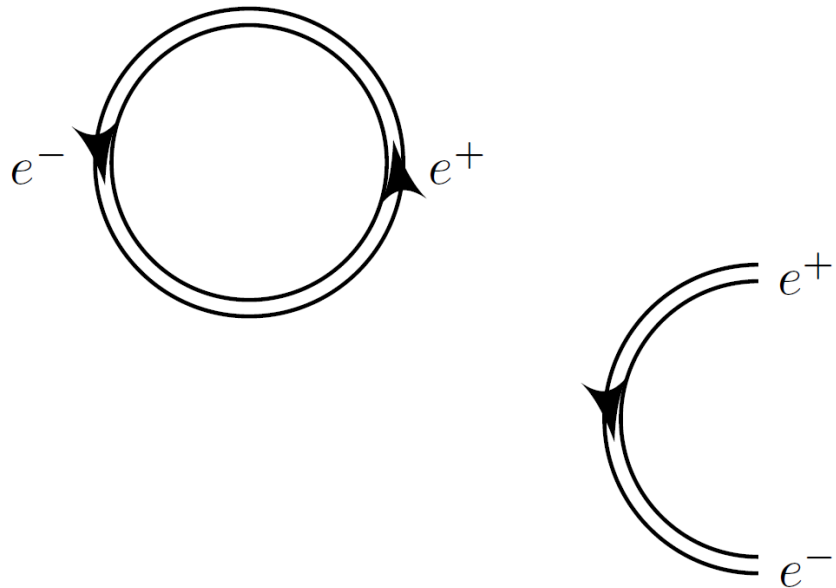
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XFELs as drivers

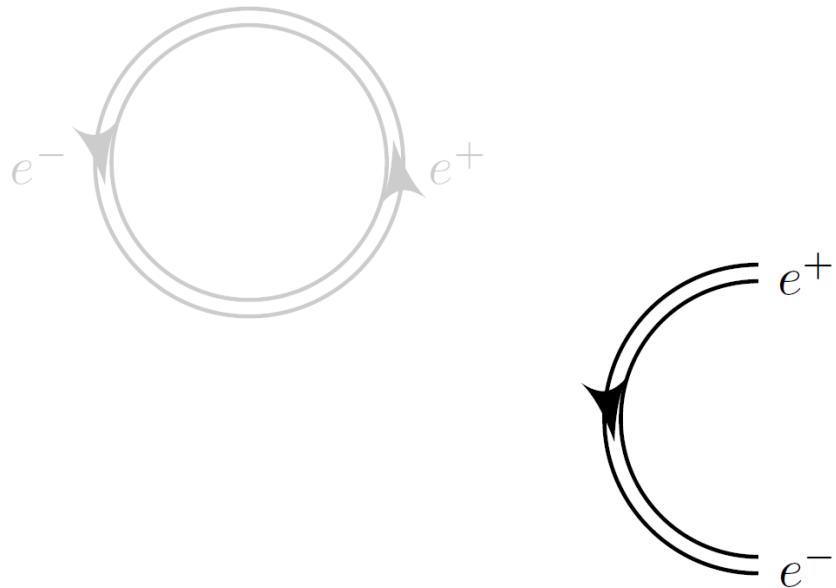
How do fields interact?



- Maxwell's equations are linear, at but ultrahigh field strengths this is not an adequate description anymore.
- Vacuum polarisation (photon-photon scattering, birefringence, dichroism)
- Schwinger pair creation/vacuum pair creation

XFELs as drivers

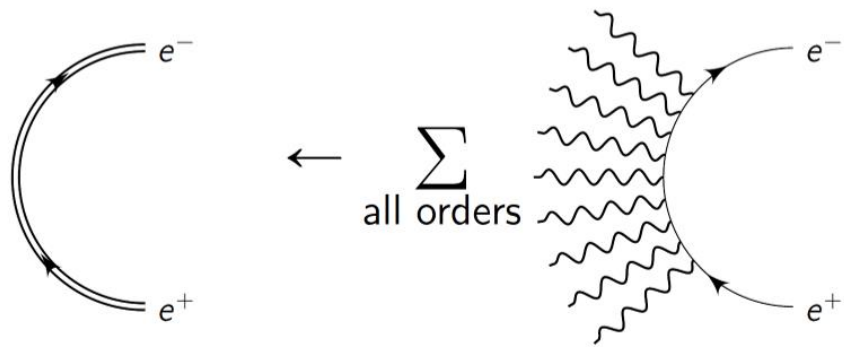
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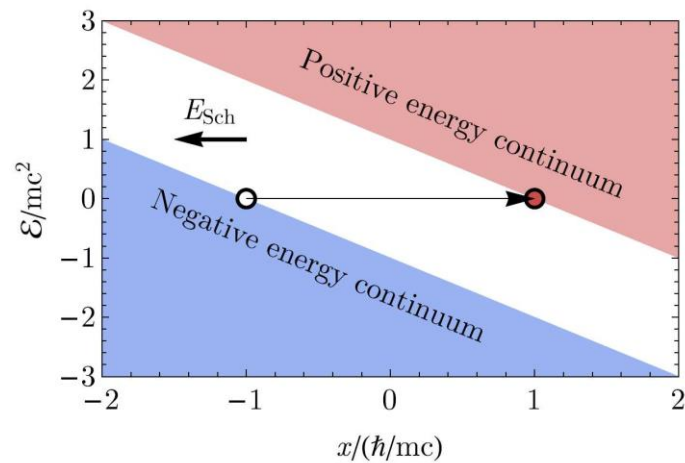
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XFELs as drivers

Schwinger effect



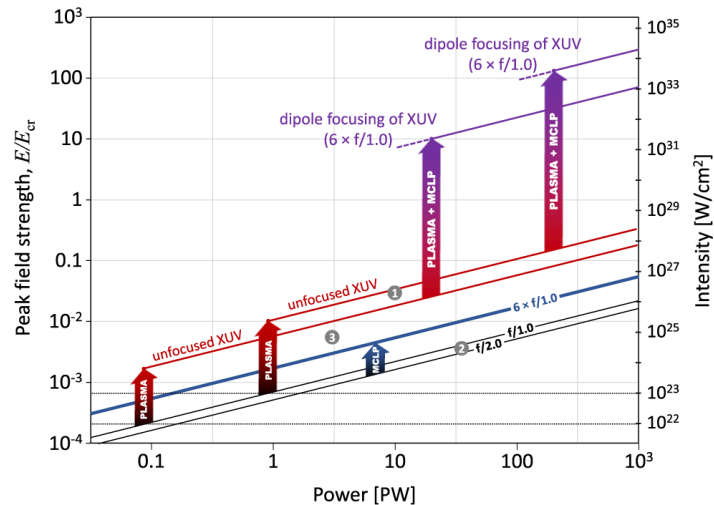
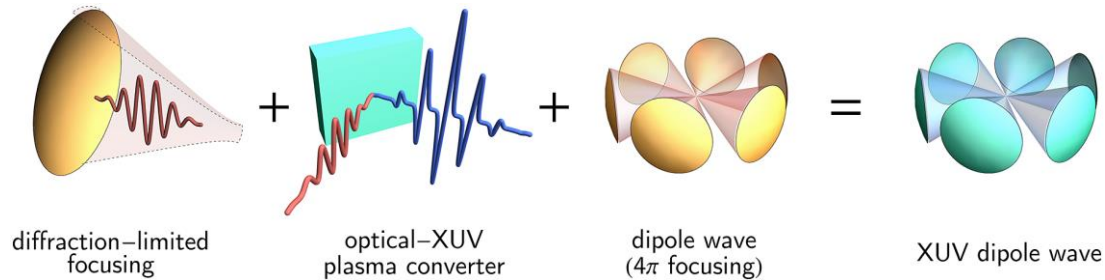
not to be taken too literally...



- The action for a static electric field acquires an imaginary part that is non-perturbative in the electric field strength.
- The characteristic field strength required is $E_{\text{crit}} = 1.3 \times 10^{18}$ V/m...
- ... which corresponds to an intensity of 2×10^{29} W/cm²

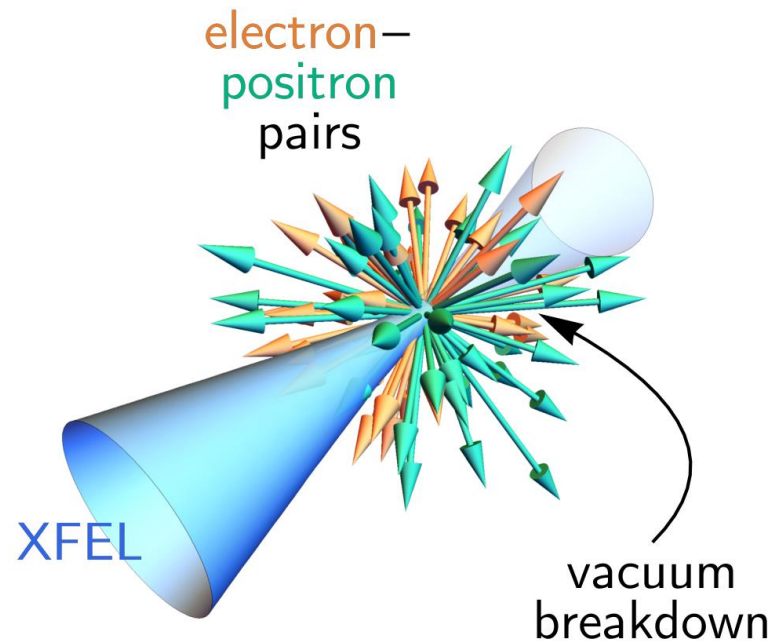
Reaching critical/supercritical field strengths

conversion of optical laser to XUV via plasma interaction [Gonoskov et al, HPLSE 2023]



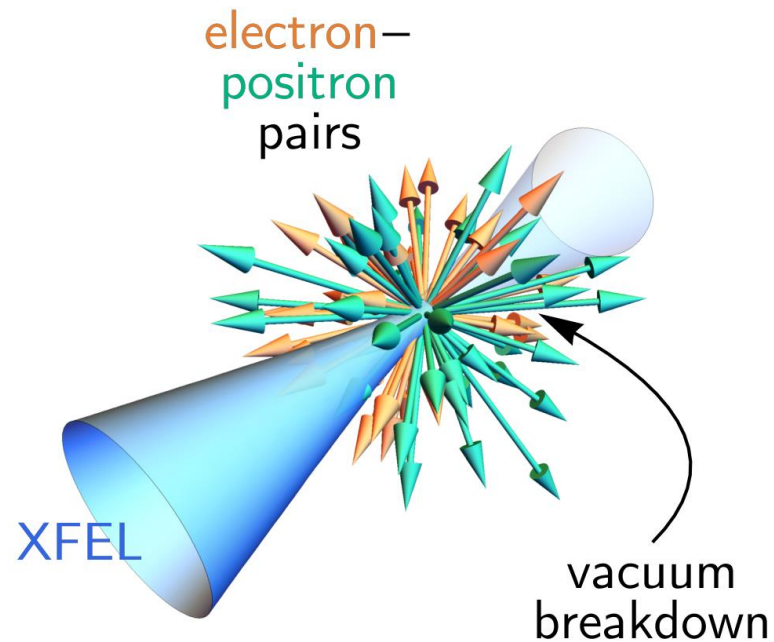
- 2×10^{29} W/cm² ?
- Aside from the technological challenges(!), it's likely impossible to reach an intensity anywhere near this with optical lasers because of pair avalanches [Bulanov et al, PRL 2010; Fedotov et al, PRL 2010].
- Routes toward “extremes” extreme field science rely on some means of conversion to higher frequencies.

Reaching critical/supercritical field strengths



- Takes advantage of smaller focal spot sizes at shorter wavelength
- ... and suppression of pair avalanche growth due to volume factors.
- In principle, Schwinger pair creation observable with a terawatt XFEL... focused to nm focal spot sizes [Ringwald, PLB 2001]

Reaching critical/supercritical field strengths



- Rate depends on invariants $E^2 - B^2$ and $E \cdot B$, not the electric field strength alone.
- Theory question: does a quasistatic model of pair creation work at high frequency/field gradient?
- “Assisted” Schwinger pair creation: reduce field strength required by combining the driver with other high-frequency EM waves.

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