

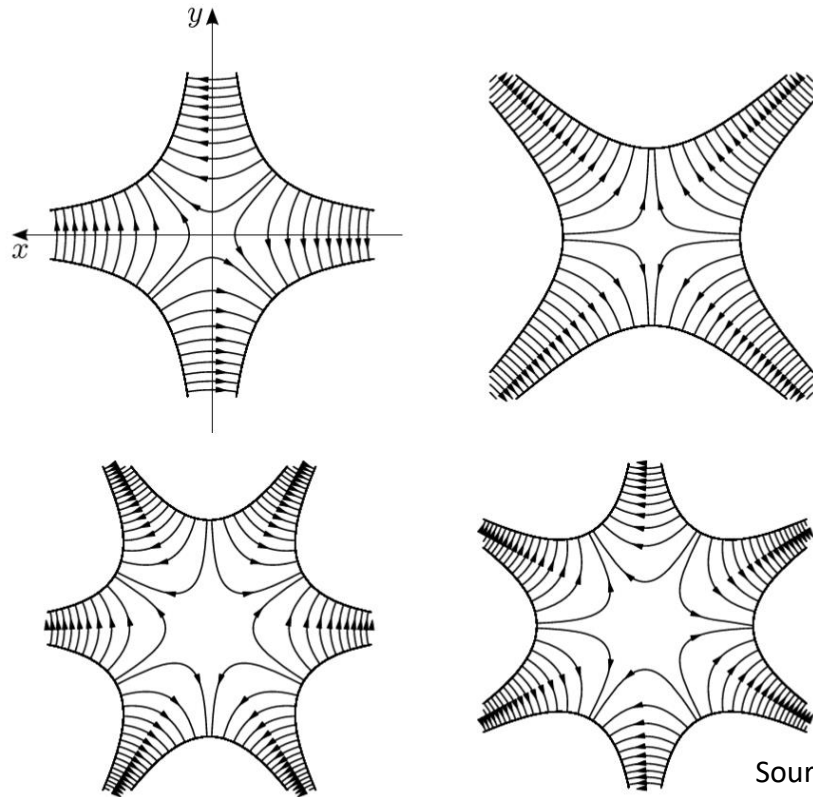
# Improved Beam Dynamics Through Control of Magnetic Fringe Fields

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# Magnets in particle accelerators

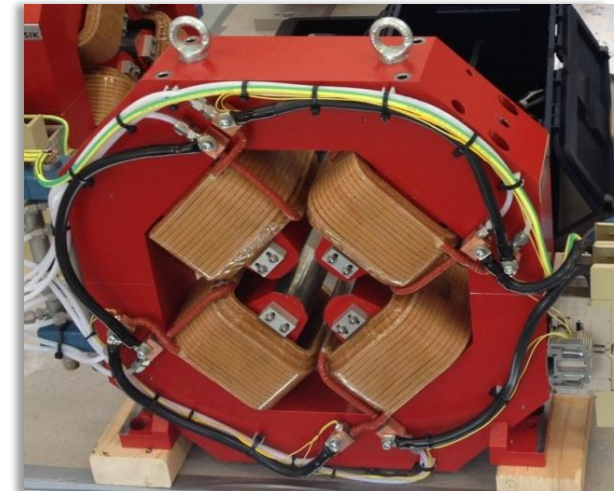
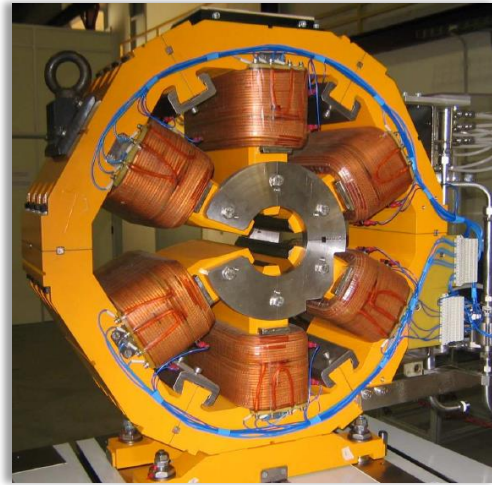
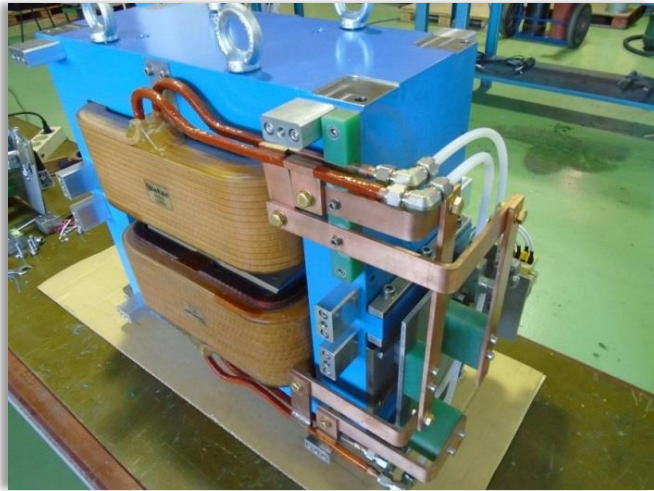
Particle accelerators use magnets that exhibit **multipole fields**.

- Dipoles are used for **beam steering**.
- Quadrupoles are used for **planar focusing**.



- These are pure multipole fields.
  - Quadrupole and sextupole fields in normal and skew orientation.
- **Real fields are not pure** and contain higher order multipole fields (allowed and forbidden).

# Realistic constraints on accelerator components

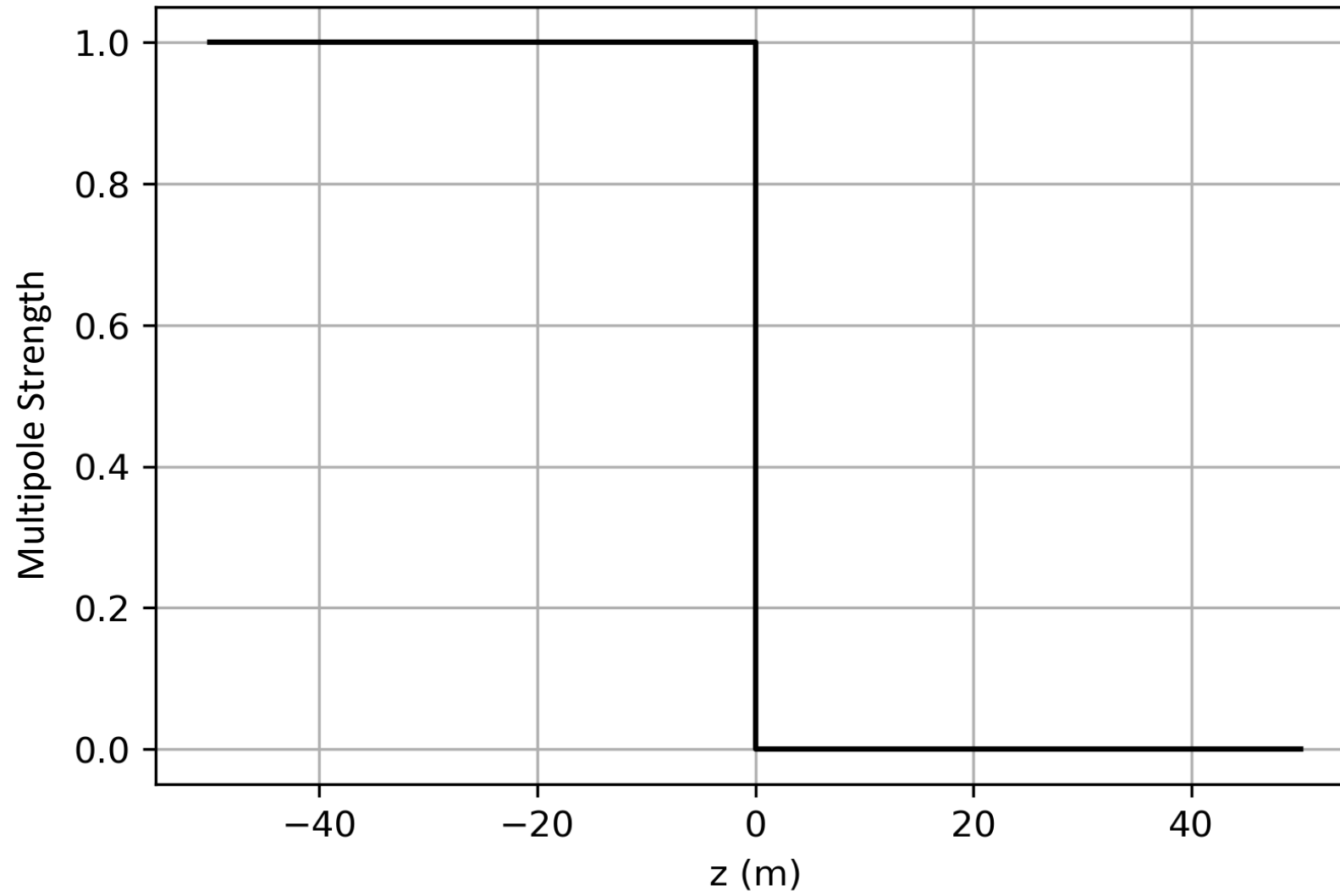


- Real magnets must have **finite dimensions**.
- In **theory**, to create pure multipole fields, **magnets should not be bound, have infinite permeability** and conductors placed at **infinity**.

# Maxwell's equations, fringe fields and modelling

- Maxwell's equations **must** be satisfied!
- We are interested in the **fringe field** regions at the **edge** of a magnet's influence.
- The “roll-off” will be referred to often, this is the **multipole strength along the magnetic axis**.
- Often, in simulations, **hard-edged field models** are assumed.
- These **do not** satisfy Maxwell's Equations.

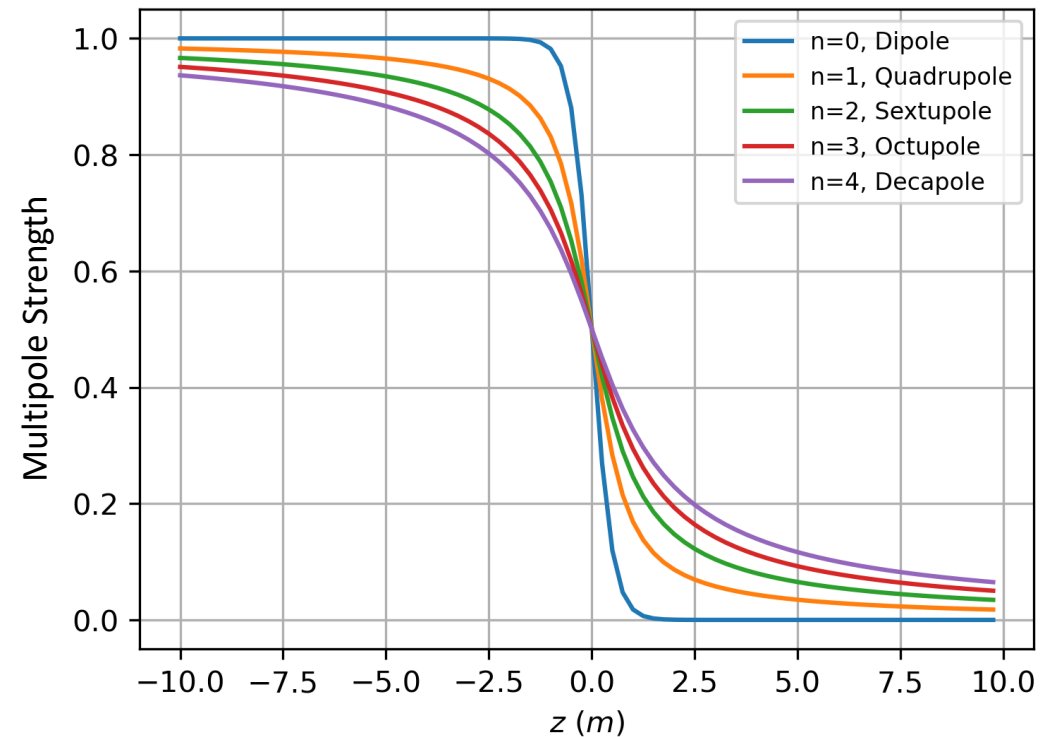
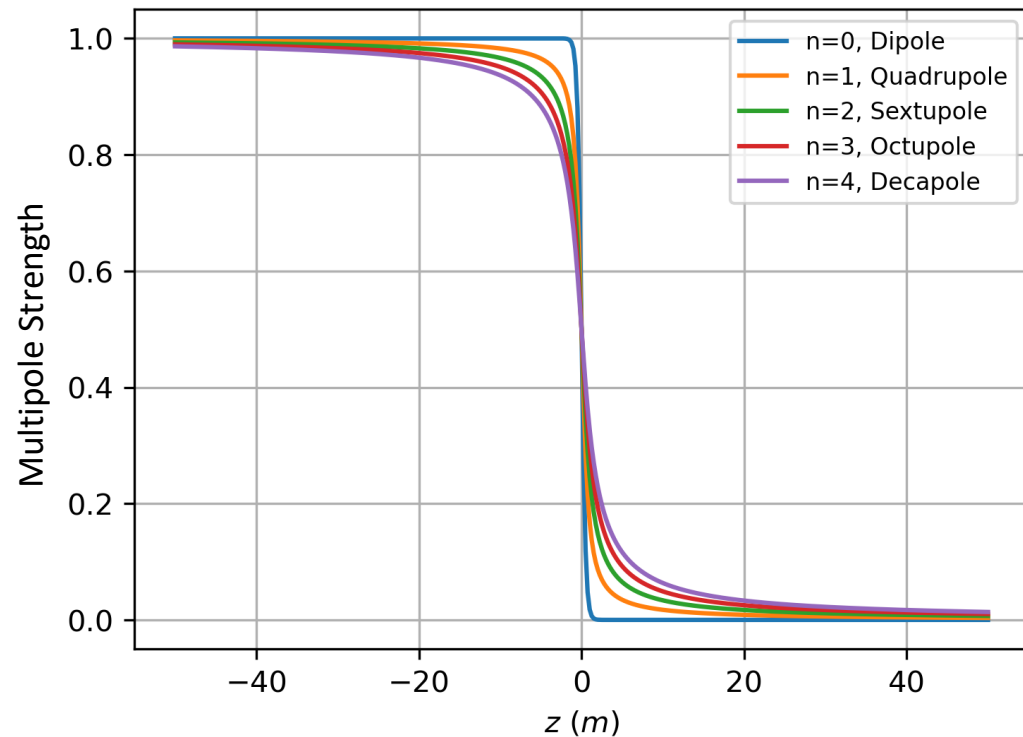
# Hard-Edge Approximation



# The Enge roll-off

- We can satisfy Maxwell's equations with the Enge roll-off.
- This is a more realistic approach (the use of the Enge is commonplace).

Enge Roll-Off For N-Order Multipoles



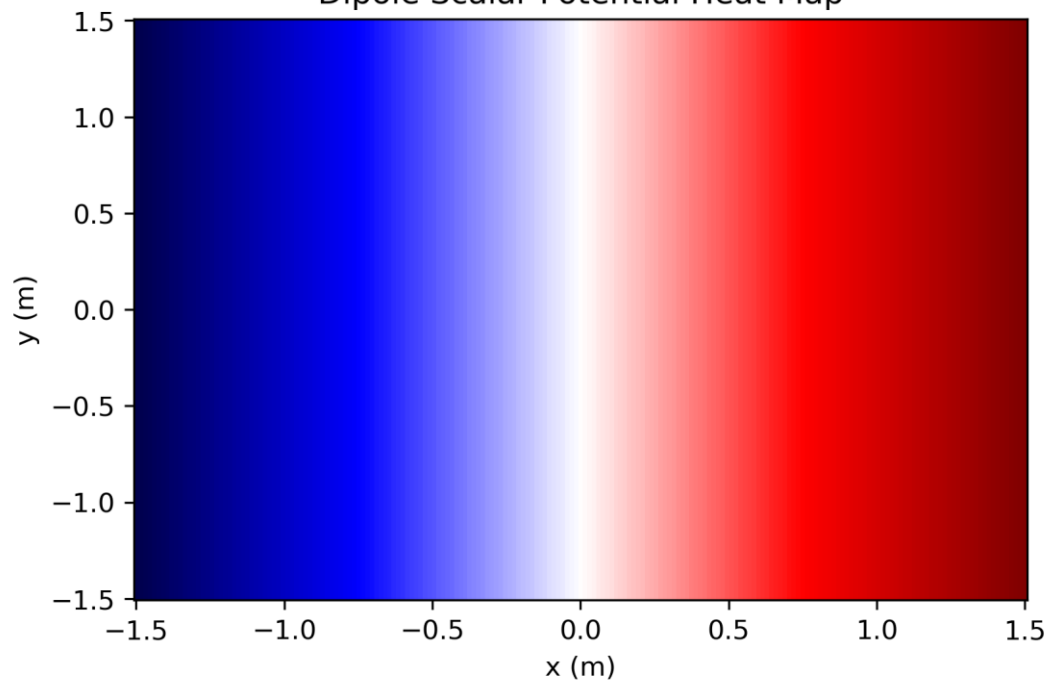
# Creating a magnet

Can we customise a magnet for a specific dynamical system?

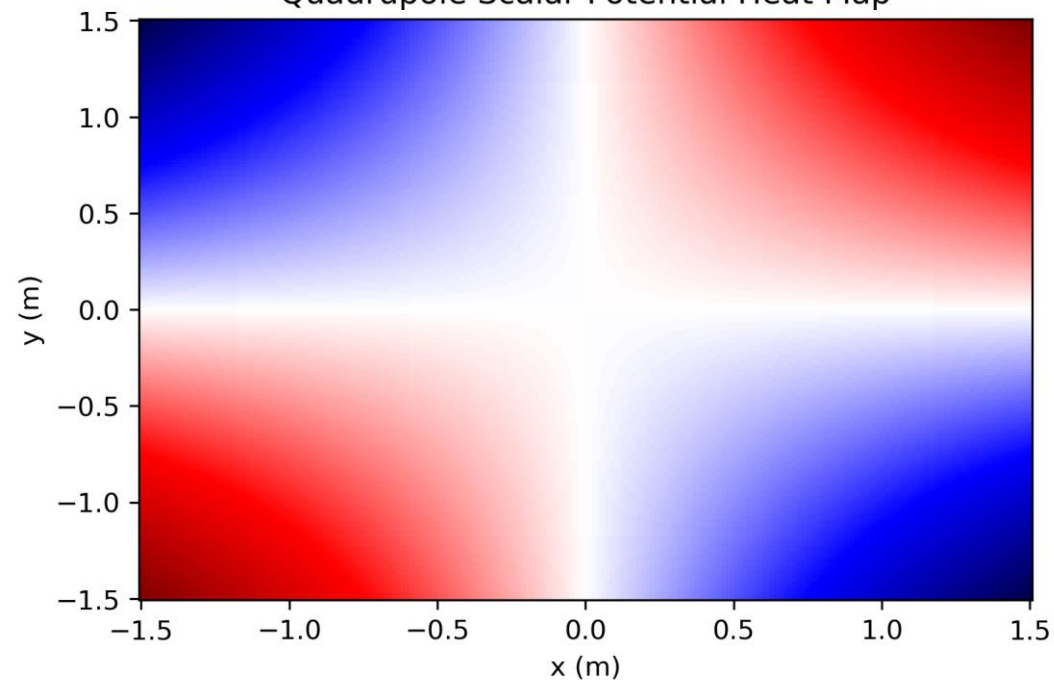
- Define the roll-off required for specific dynamics.
- Approximate the roll-off as **Fourier** series (for later generalisation).
- Produce the scalar potential analytically\*.
- Surfaces of **constant scalar potential** define the **pole face shape** of the magnet.
  - This is the case if we assume the magnet has **infinite permeability**.
- Create the magnet (e.g. in OPERA)!

\*Muratori, B. Wolski, A, Jones, J. Analytical Expressions for Fringe Fields in Multipole Magnets. (2014). Physical Review Special Topics – Accelerators and Beams. DOI: 10.1103/PhysRevSTAB.18.064001.

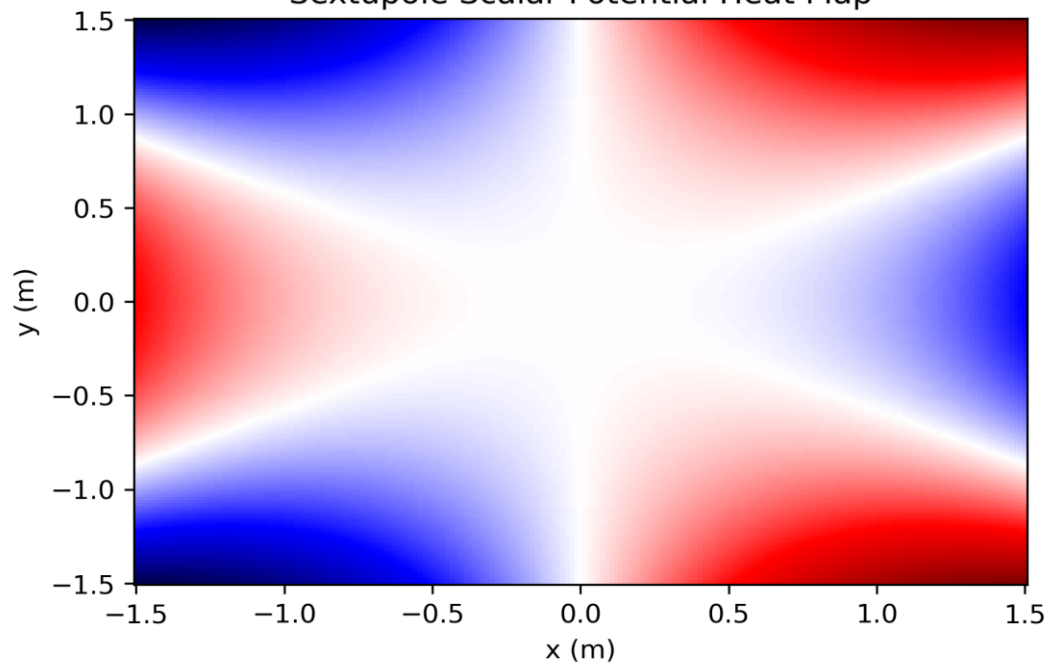
Dipole Scalar Potential Heat Map



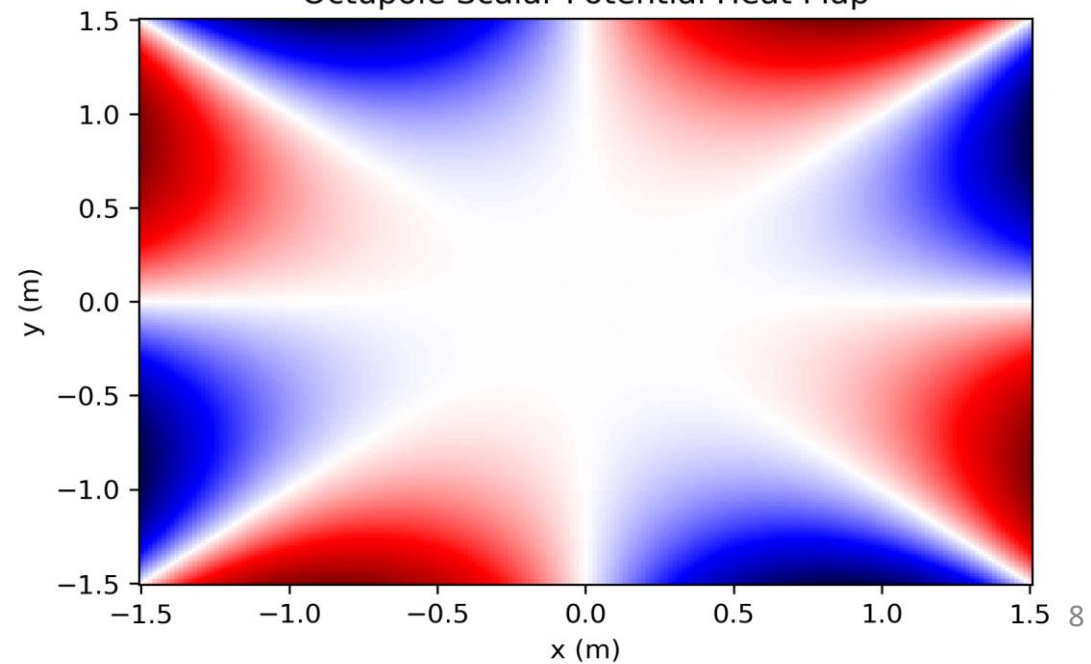
Quadrupole Scalar Potential Heat Map



Sextupole Scalar Potential Heat Map

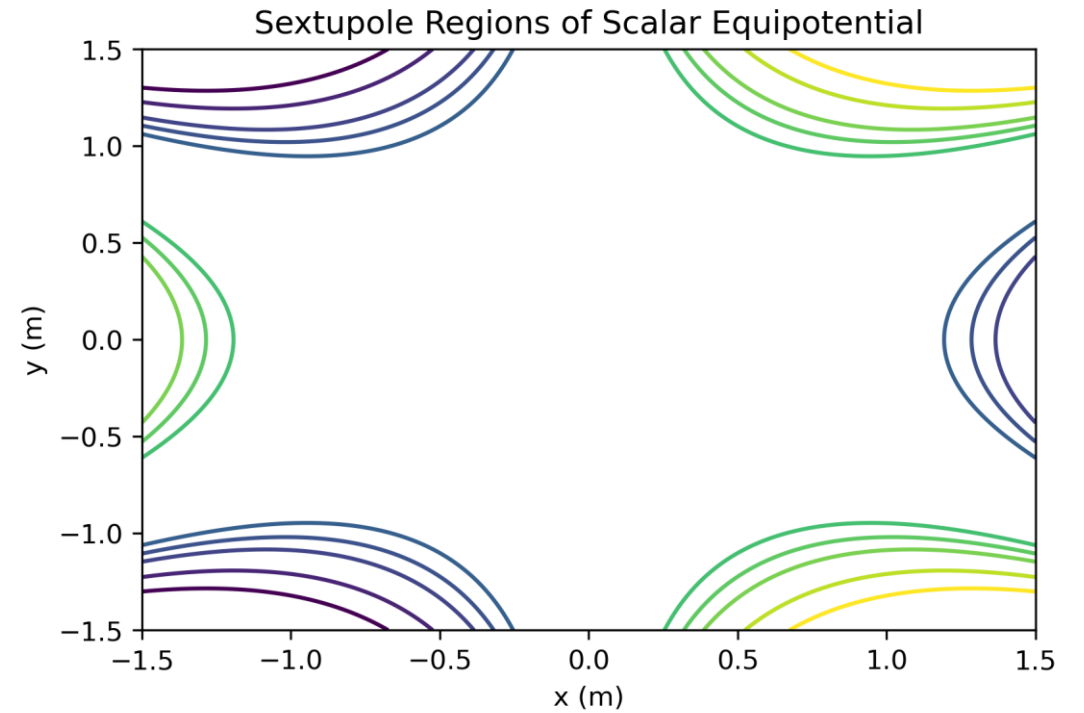
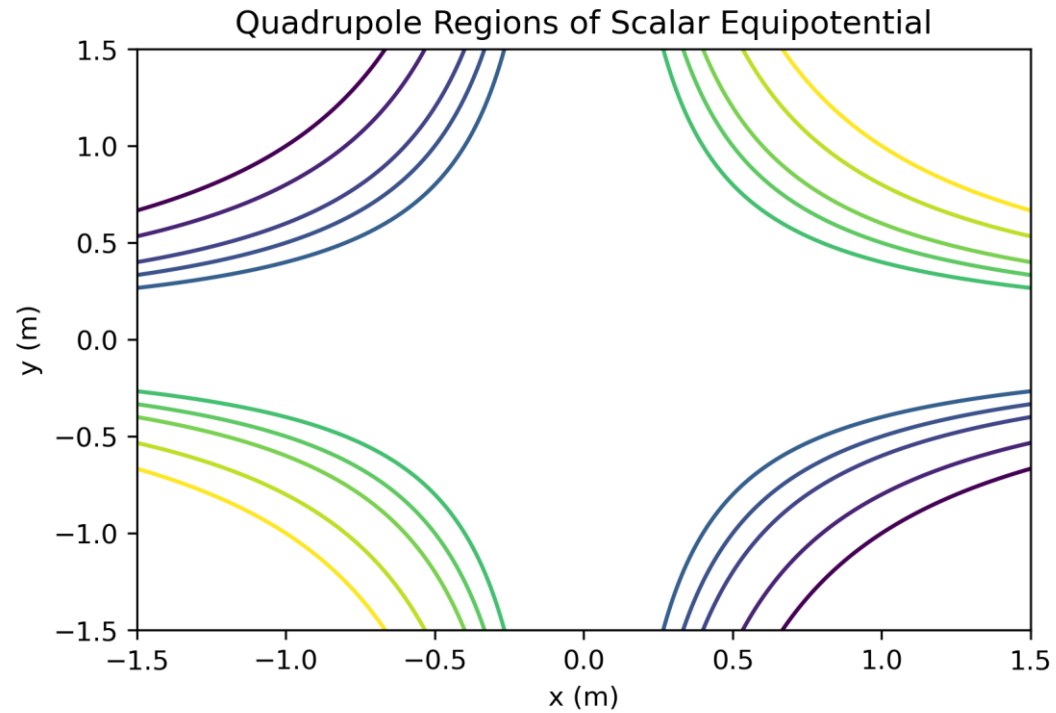


Octupole Scalar Potential Heat Map

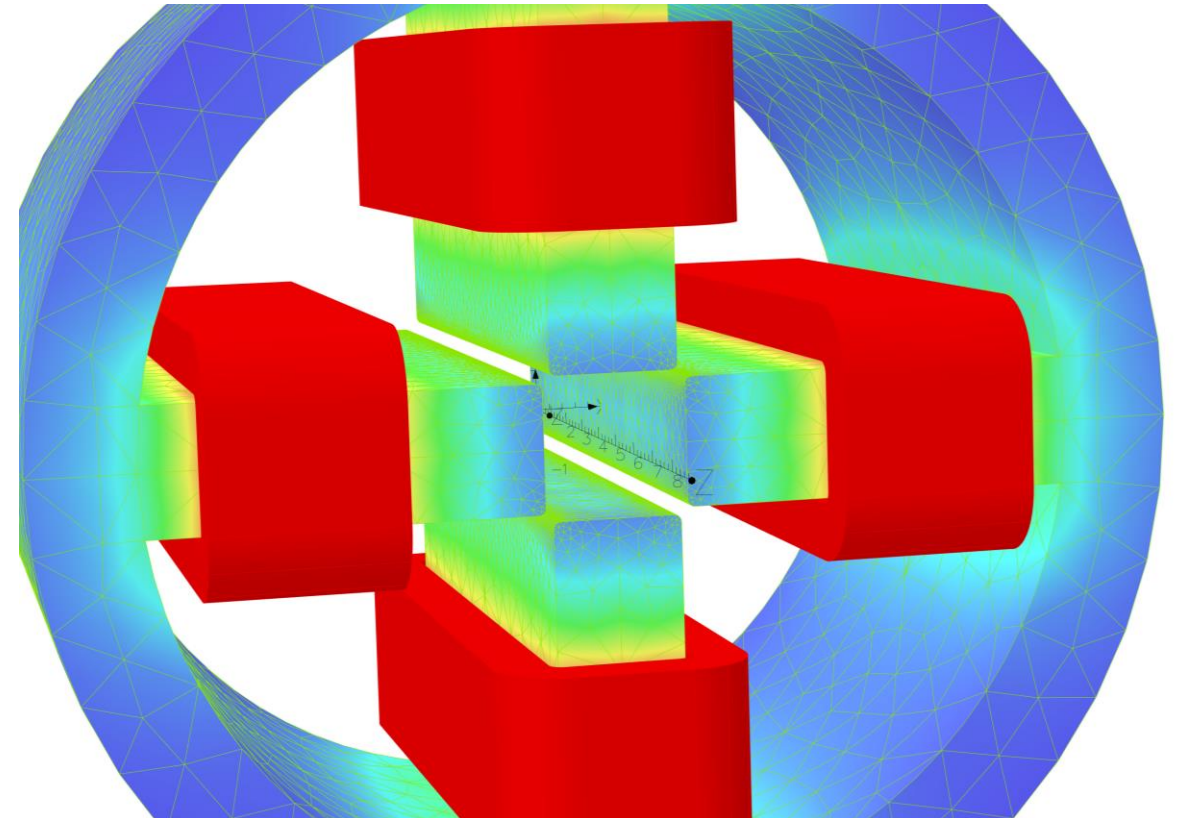
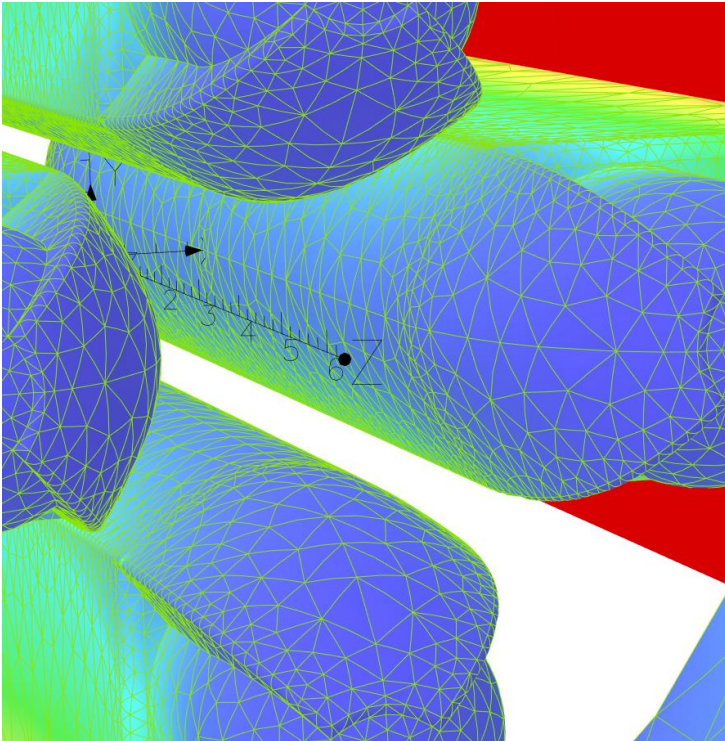




# Scalar equipotential



# Create and test the magnets



\*Work in progress.

# Progress so far

- We've now progressed from the **roll-off function (multipole strength)** to the **fields**, the **scalar potential** and finally to the **pole face shape**.
- Process is slow, a **faster solution** from roll-off to pole face would have significant benefits.

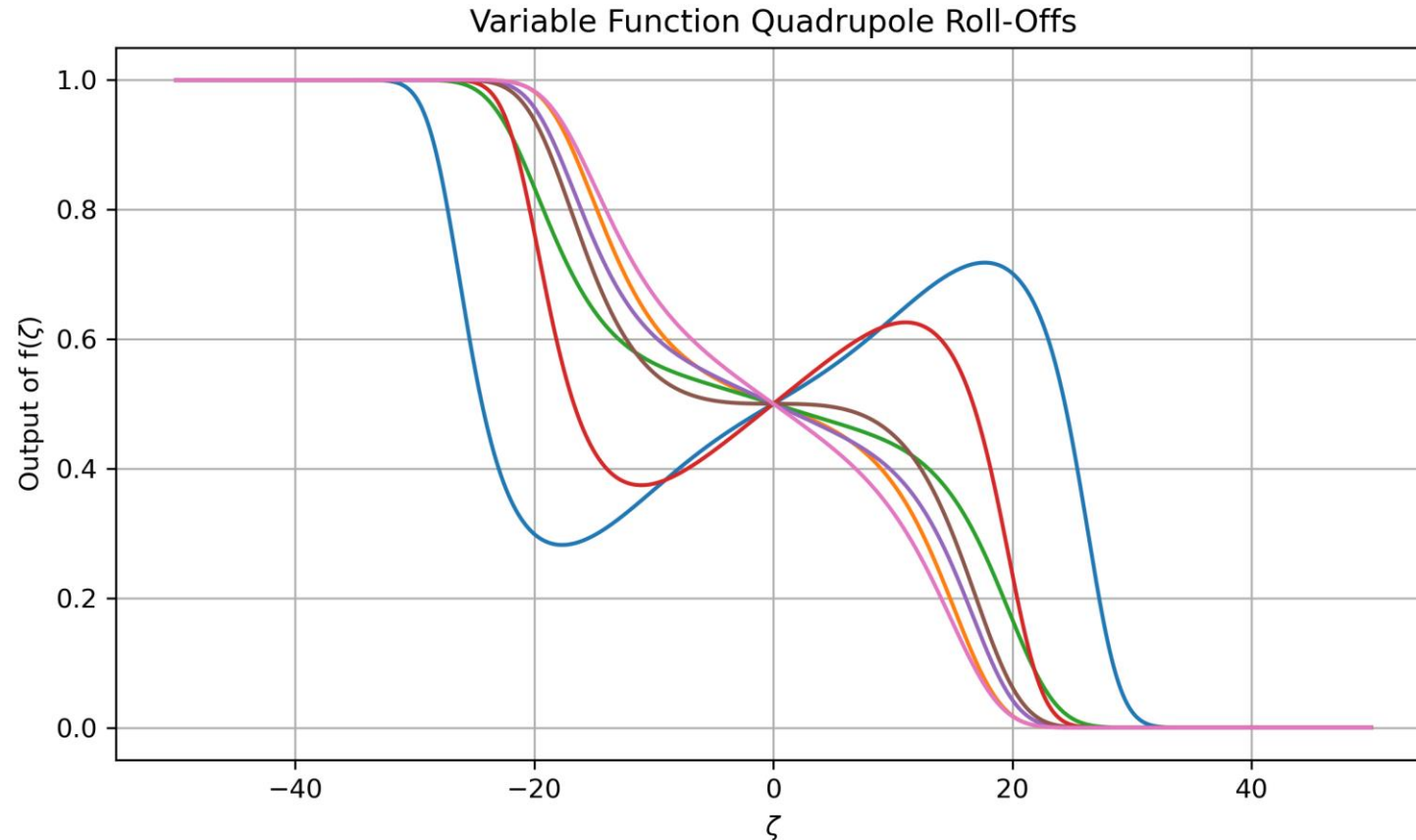
# Machine learning applications

The goal, use **machine learning** to move directly between the **roll-off function** and the **pole face shape**.

1. Modify the **Enge function** to produce many field roll-offs that take **different forms**.
2. Create the associated **scalar potential** for each given **roll-off**.
3. Train a **neural network** to learn and apply the underlying patterns of this relationship.

Using **machine learning saves time** and allows for **quick computation of roll-offs beyond the Enge**.

# The variable Enge-type roll-off



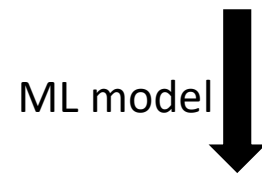
The shape of these roll-offs are manipulated by altering three Enge coefficients, A, B and C.

# Order of machine learning processes

Engage coefficients, A, B and C.

# Order of machine learning processes

Enge coefficients, A, B and C.



Roll-off shape (field gradient at entrance / exit).

# Order of machine learning processes

Enge coefficients, A, B and C.



Roll-off shape (field gradient at entrance / exit).

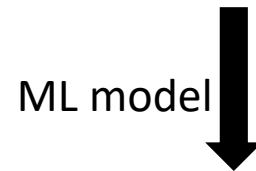


# Order of machine learning processes

Enge coefficients, A, B and C.



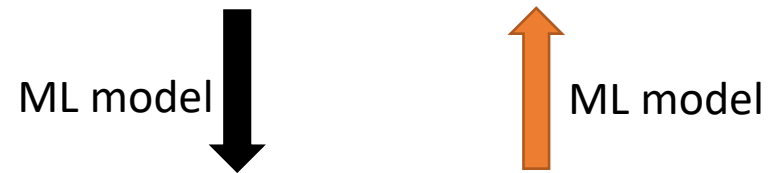
Roll-off shape (field gradient at entrance / exit).



Scalar potential.

# Order of machine learning processes

Enge coefficients, A, B and C.



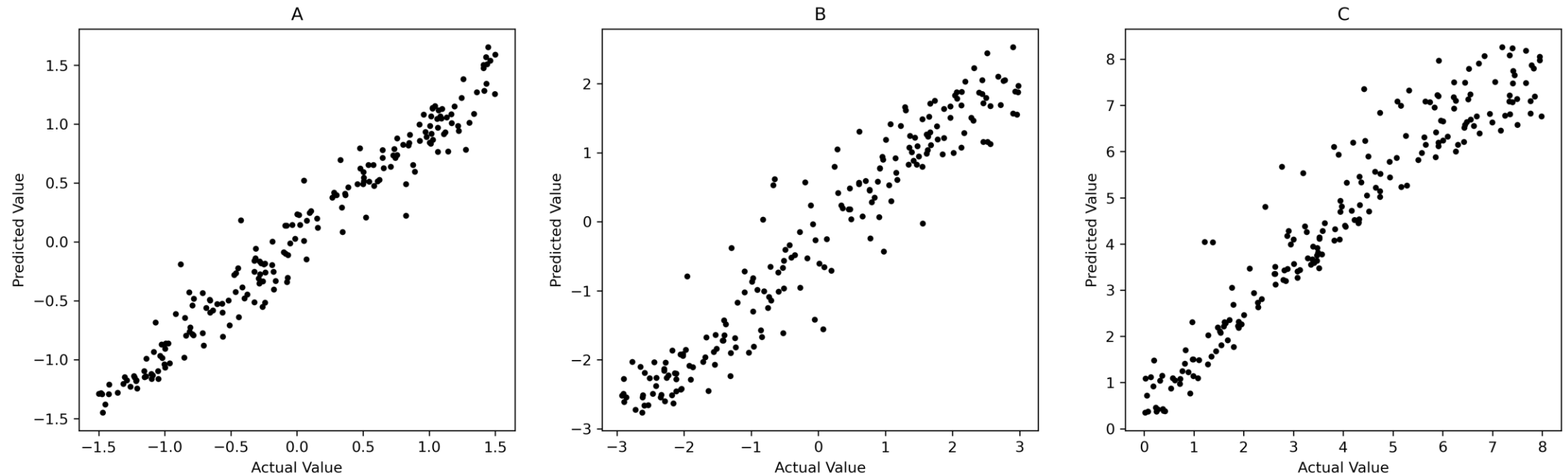
Roll-off shape (field gradient at entrance / exit).



Scalar potential.

# Predicting Enge coefficients from the roll-off

Enge Coefficients From a Trained Network: First Iteration



\*Work in progress.

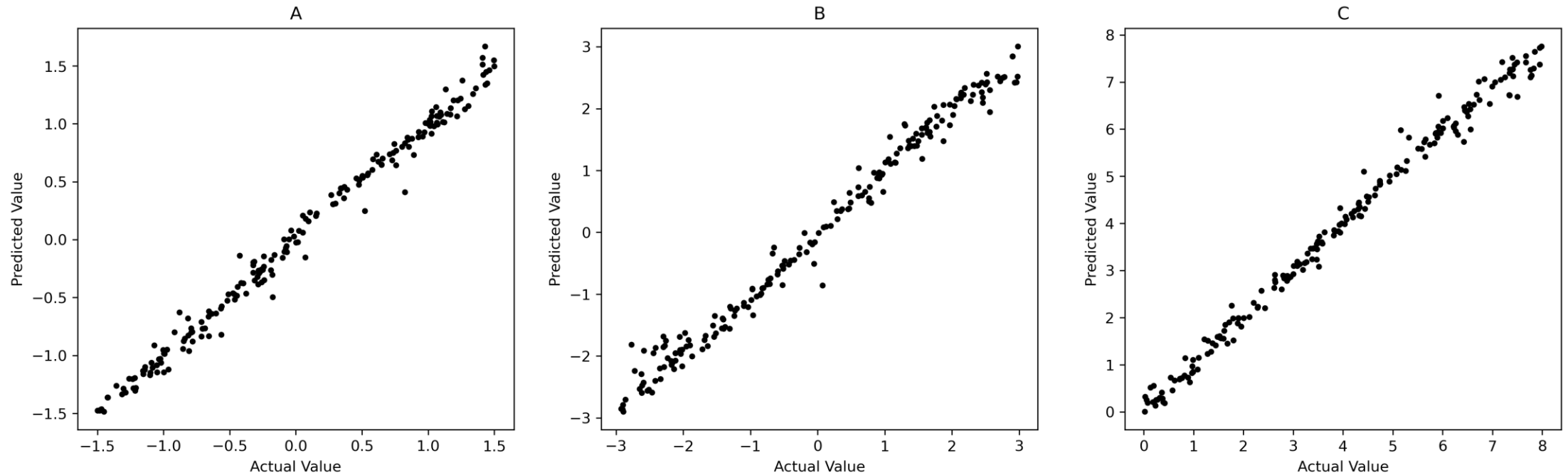
Enge coefficients, A, B and C.



Roll-off shape (field gradient at entrance / exit).

# Predicting Enge coefficients from the roll-off

Enge Coefficients From a Trained Network: Second Iteration



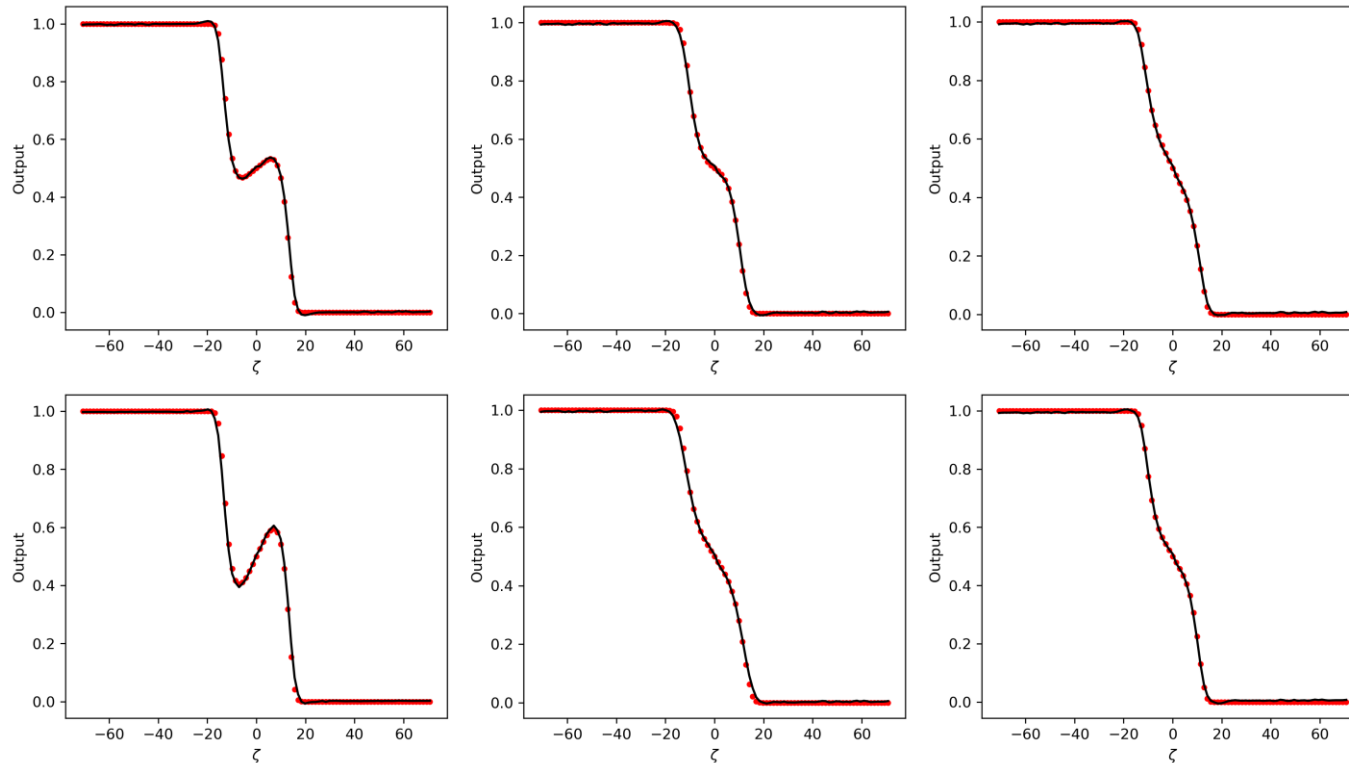
\*Work in progress.

Enge coefficients, A, B and C.



Roll-off shape (field gradient at entrance / exit).

# Predicting the roll-off from the coefficients



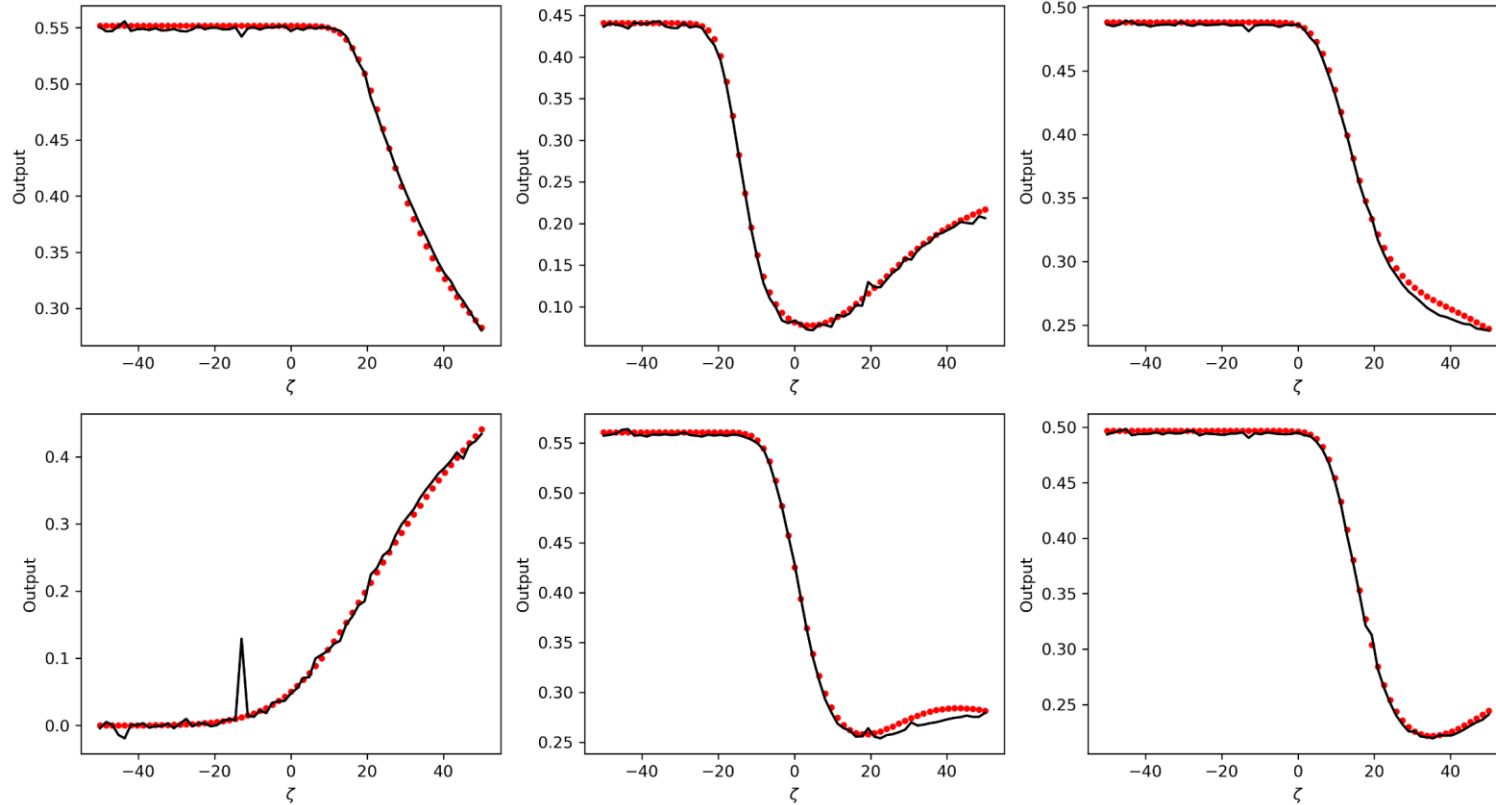
\*Work in progress.

Engge coefficients, A, B and C.



Roll-off shape (field gradient at entrance / exit).

# Predicting the roll-off from the scalar potential



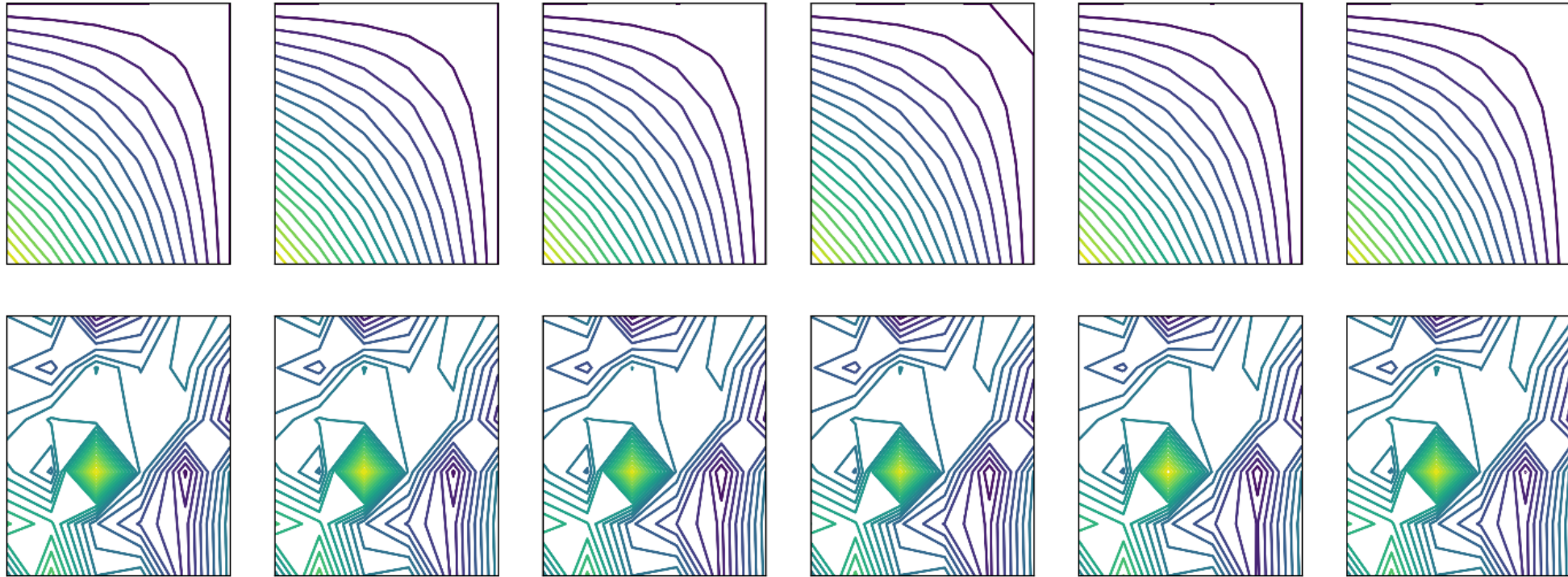
\*Work in progress.

Roll-off shape (field gradient at entrance / exit).



Scalar potential.

# Optimising the scalar potential model



\*Work in progress.

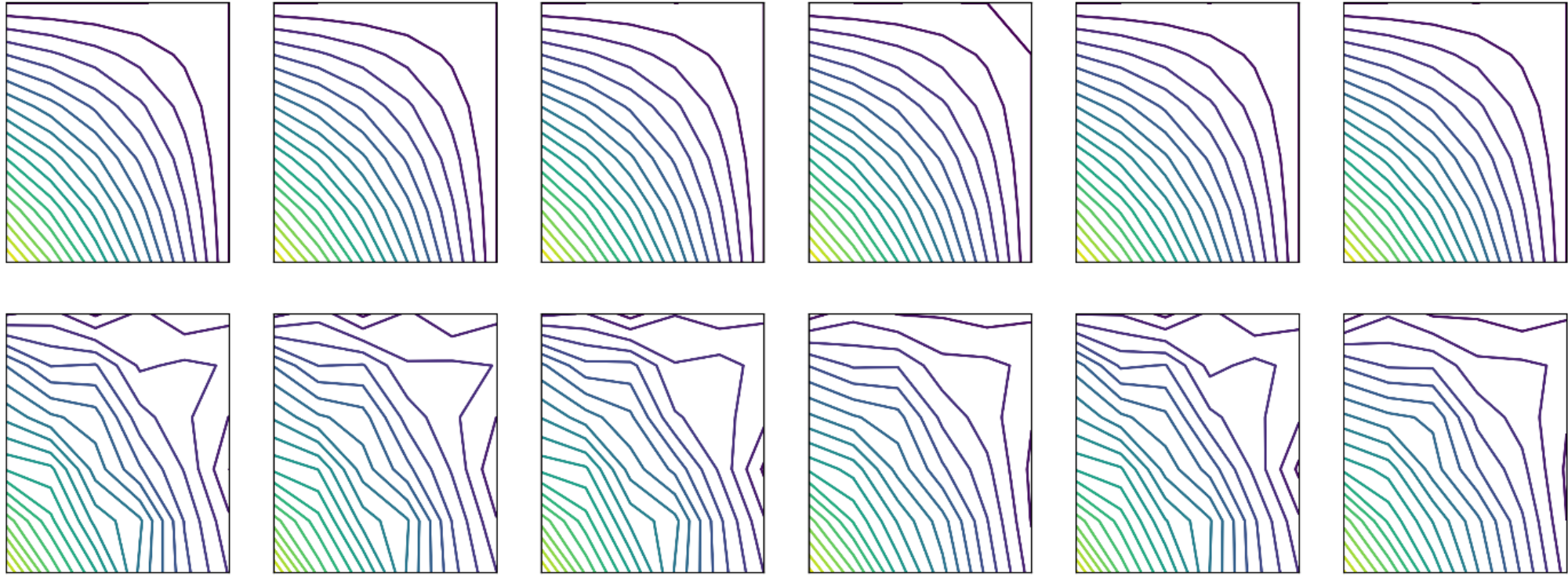
Roll-off shape (field gradient at entrance / exit).

ML model



Scalar potential.

# After some more optimisation



\*Work in progress.

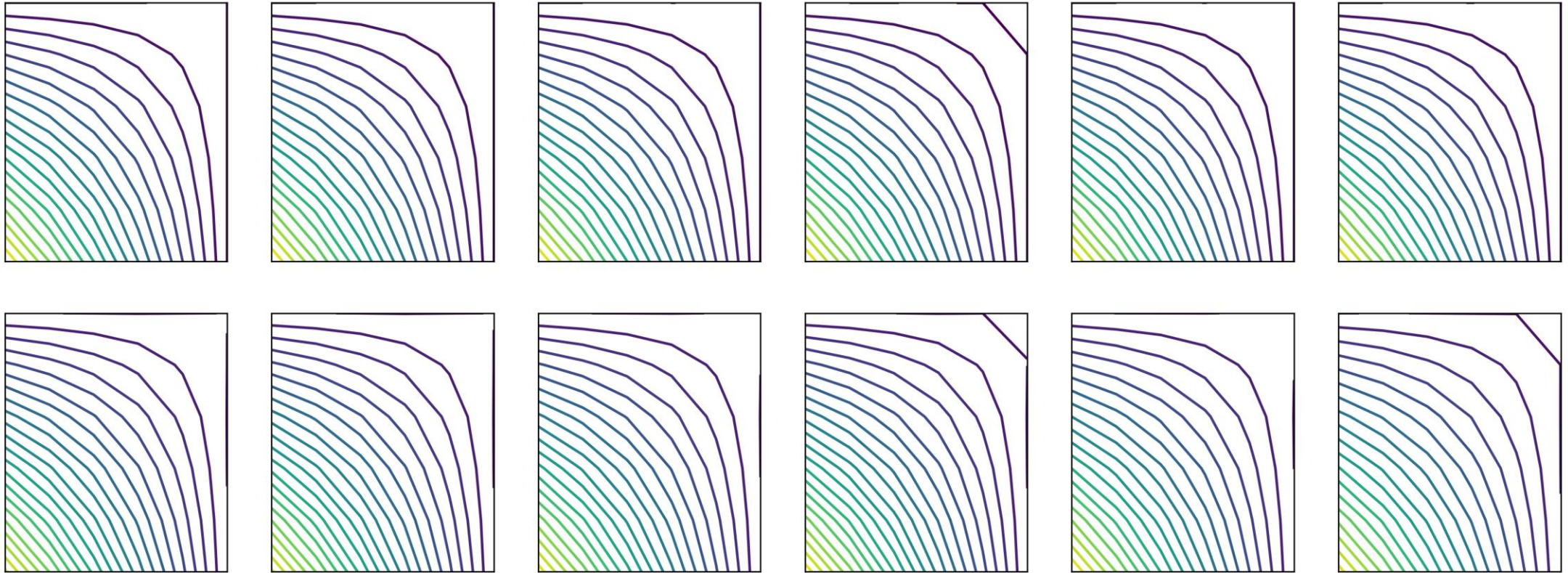
Roll-off shape (field gradient at entrance / exit).

ML model



Scalar potential.





\*Work in progress.

Roll-off shape (field gradient at entrance / exit).

ML model



Scalar potential.

# Beam dynamics and future study

The next step is to model the **beam dynamics**.

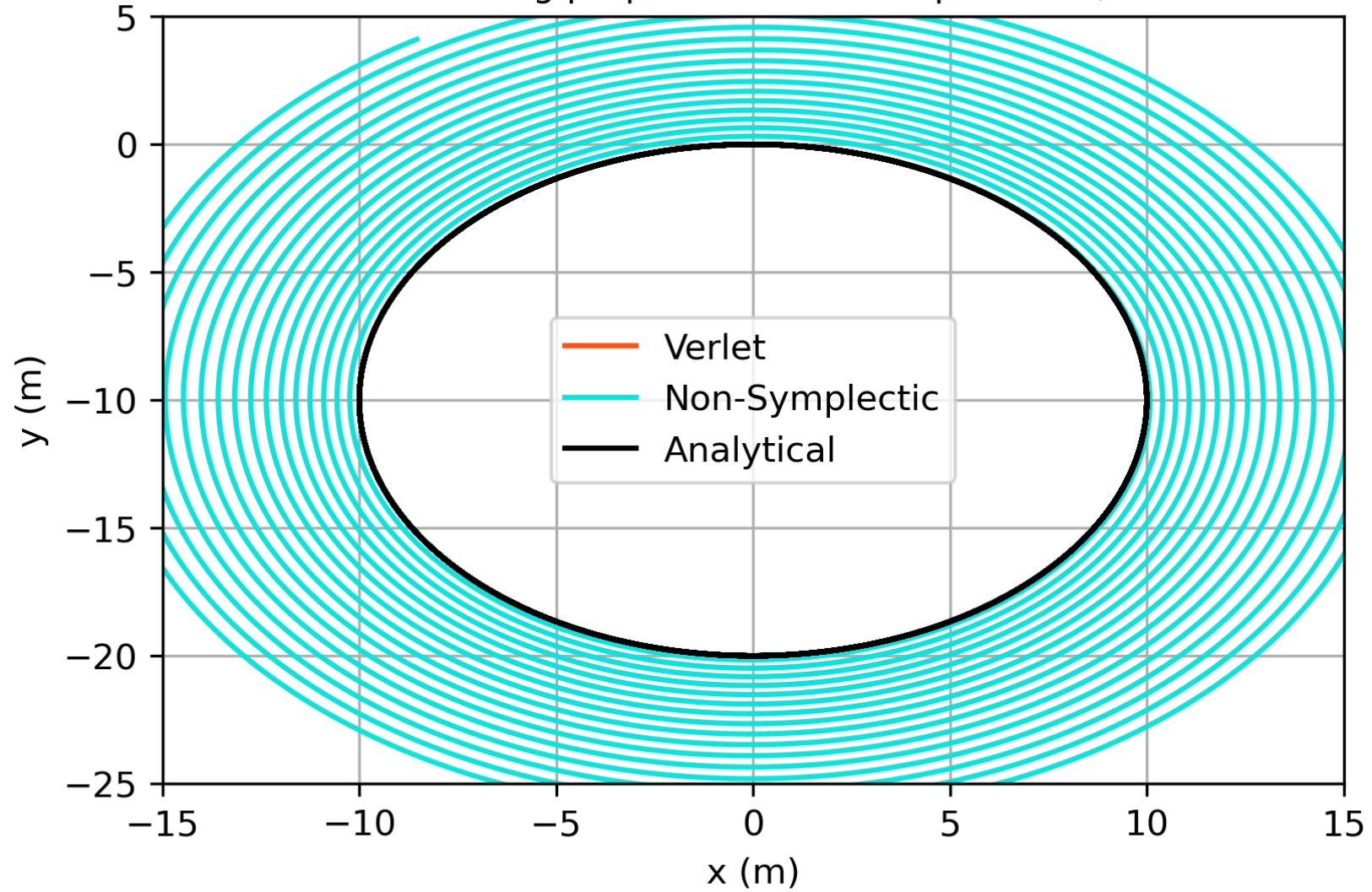
Including beam dynamics in the set of machine learning tools will allow rapid optimisation of magnet geometry for specified dynamical properties.

Modelling the beam dynamics requires implementation of an integrator (particle tracker).

# Magnetic velocity Verlet method

- The **force exerted** by a magnetic field on a charged particle **is velocity dependent**. Think back to the **Lorentz force** equation.
- An integrator has been specifically created for this situation.
  - “A Magnetic Velocity Verlet Method.” Chambliss, A. Franklin, J (2020).
- The magnetic **velocity Verlet** method **is symplectic**.

Particle travelling perpendicular to a dipole field,  $\delta t = 0.1$



\*Work in progress.

# Future study

- Use of the particle tracker is yet to be fully implemented and optimised.
- The **transfer functions** through the fringe fields are to be constructed from the particle tracker.
- Ultimate goal – **create a magnet quickly that satisfies specified beam dynamics.**

# Thank you for your time!

A special thank you to

Andrzej Wolski

Bruno Muratori

Tessa Charles

Alex Bainbridge

Amelia Pollard



# Appendix A – Non symplectic integrator

$$x_{n+1} = x_n + \dot{x}_n \Delta t$$

$$a_n = \frac{q}{m} (\dot{x}_n \times B)$$

$$\dot{x}_{n+1} = \dot{x}_n + a_n \Delta t$$



# Appendix B – Verlet integrator

$$x_{n+1} = x_n + \dot{x}_n \Delta t + \frac{1}{2m} F_n \Delta t^2$$

$$\dot{x}_{n+1} = \dot{x}_n + \frac{1}{2m} F_{n+1} \Delta t^2$$

We don't have access to the force value at the new position without velocity (Lorentz force).

# Appendix C – Velocity Verlet integrator

Consider,

$$\dot{x}_{n+1} = \dot{x}_n + \frac{q\Delta t}{2m} ((\dot{x}_n \times B_n) + (\dot{x}_{n+1} \times B_{n+1}))\Delta t^2$$

It can be shown that,

$$\dot{x}_{n+1} = \frac{1}{1 + \left(\frac{q\Delta t}{2m}\right)^2 C \cdot C} \left[ d + \frac{q\Delta t}{2m} d \times C + \left(\frac{q\Delta t}{2m}\right)^2 C(d \cdot C) \right]$$

Where,

$$C = B_{n+1}$$

$$d = \left( \dot{x}_n + \frac{q\Delta t}{2m} \dot{x}_n \times B_n \right)$$