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Linear and non-linear beam dynamics of the ISIS-II FETS vertical FFA using Zgoubi

Presented at the 2022 Workshop on Fixed Field alternating gradient Accelerators (FFA2022)

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September, 29th 2022





Linear and non-linear beam dynamics of the ISIS-II FETS vertical FFA using Zgoubi

Scaling law:
$$B = B_0 e^{kz}$$
, $(k = \frac{1}{B_z} \frac{\partial B_z}{\partial z})$

- \rightarrow Vertical orbit excursion
- → Strongly non-linear magnetic fields
- \rightarrow Strongly coupled optics

With the fringe field function g(x), the median plane field components are: $B_{y0}(x, 0, z) = 0$, $B_{z0}(x, 0, z) = B_0 e^{kz} g(x)$, $B_{x0}(x, 0, z) = \frac{B_0}{k} e^{kz} \frac{dg}{dx}$.

Linear and non-linear beam dynamics of the ISIS-II FETS vertical FFA using Zgoubi

The Zgoubi ray-tracing code:

- Often used to study horizontal FFAs
- Particle tracking in complex field maps
- Possibility to use its Python interface Zgoubidoo

- Presentation of the lattice and methods
 - Baseline lattice with arctan fringe field
 - Fieldmap construction and tracking with Zgoubi

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 - Closed orbit search
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 - Closed orbit search
 - Tune comparison with neighboring cells contribution
- Non-linear beam dynamics Dynamical aperture
 - Possible simple definition/computation for the 2D Dynamical aperture
 - Methods/suggestions for calculating a dynamic aperture for 4D motion in a highly nonlinear and coupled lattice

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Presentation of the lattice and methods

Baseline arctan lattice & Methods to track with Zgoubi in adapted field map

ULB Lattice and methods – Arctan baseline lattice

FETS-vFFA baseline lattice:

• Realistic **fringe field** fall off in Arctan(z)

 $g(z) = \left[\arctan\left(\frac{z + M/2}{L_m}\right) - \arctan\left(\frac{z - M/2}{L_m}\right) \right] / \pi$

• Parameters given by S. Machida:

Energy	3 to 12 MeV		
Repetition	50 Hz		
Number of proton per pulse	3.4 x 10 ¹¹		
Focusing	FDF triplet		
Circumference	28 m		
Number of cell	10		
Total cell length	2.8 m		
Bd and Bf core length (M)	0.50 m		
Straight length	1.24 m		
Distance between Bd centre and Bf centre	0.53 m		
Horizontal displacement between Bd and Bf	+/- 0 mm		
Fringe field parameter (L)	0.15 m		
Bd/Bf ratio (nominal)	1.15		
m-value (nominal)	1.31		
Orbit excursion	0.53 m		
Tune (qu, qv, nominal)	0.243445 / 0.120023 (0756555 / 0.120023)		
Dynamic aperture (normalised)	60 pi / 70 pi		
Nominal 100 % emittance (normalised)	10 pi		



Figures of S. Machida

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ULB Lattice and methods – Arctan baseline lattice

- Search for **closed orbits** at different energies **for a single cell field map** without taking into account the residual field of neighboring cells:
 - Need to relax the $p_0 = 0$ (vertical angle) condition to find something stable
 - p_0 !=0 at the cell ends \rightarrow Not a closed orbit on the entire ring
- The **vertical orbit excursion** is still visible and is similar to the one obtained with other codes



• The magnetic field scales with energy



ULB Lattice and methods – Arctan baseline lattice

Significant neighboring cells influence due to the important residual field at the cell ends

 \rightarrow The residual field in the lattice in arctan is ± 0.1kG in all directions.

 \rightarrow The transverse residual field explains the non-zero vertical angle (p_0 !=0 at the cell ends)



ULB Lattice and methods – Methods

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Methods to account for the residual field of neighboring cells:

1) Linear superposition by extending field maps

- Extent the integration zone to have field maps overlapping
- Valid for small residual fields to ensure that the trajectory deviation due to this field is limited
- In the arctan lattice, the **residual field is not negligible**
 - ightarrow Method not valid for this lattice

ULB Lattice and methods – Methods

Methods to account for the residual field of neighboring cells:

HIGHBORING CELL MICHBORING CELL

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2) Actual superposition of field maps

 Superposition of field maps with the same meshes with Zgoubi → Construct the 'left' and 'right' field maps due to rotated neighboring cells

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ULB Lattice and methods – Methods

Methods to account for the residual field of neighboring cells:

RUXELLE 8 2) I. 0 MAIN CELL LIBR чц UNIVERSIT FIELD MAP DUE TO RIGHT FIELD MAP DUE TO LEFT NEIGHBORING CELL **NEIGHBORING CELL** NEIGHBORING NEIGHBORING / CELL CE) 2) I. Longitudina, additional extensio. $L = L_{cell}$ High rigidity particles on the **DRIFT** NEIGHBORING FIELDMAP transverse mesh points Zgoubi gives $(B_{x'}, B_{y'}, B_{z'})$ →interpolate on longitudinal mesh points and transform into Horizontal additional (B_x, B_v, B_z) extension

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2) Actual superposition of field maps

- Superposition of field maps with the same meshes with Zgoubi \rightarrow Construct the 'left' and 'right' field maps due to rotated neighboring cells
- Construction of truncated field maps with 'finite' Ι. longitudinal/horizontal extents using Zgoubi
 - Influence of the horizontal/longitudinal field map • extent on the orbit
 - \rightarrow Convergence to neighboring field maps that cover the entire orbit region

ULB Lattice and methods – Methods

Methods to account for the residual field of neighboring cells:

2) II. MAIN CELL FIELD MAP DUE TO RIGHT FIELD MAP DUE TO LEFT **NEIGHBORING CELL** NEIGHBORING CELL NEIGHBORING CELL CELL 2) II. MAIN CELL NEIGHBORING FIELDNIAD For each point of the mesh: $(x, y, z) \rightarrow (x', y', z')$ Compute $B' = (B_{r'}, B_{y'}, B_{z'})$ $(B_{x'}, B_{y'}, B_{z'}) \rightarrow (B_x, B_y, B_z).$ (inverse rotation)

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2) Actual superposition of field maps

- Superposition of field maps with the same meshes with Zgoubi → Construct the 'left' and 'right' field maps due to rotated neighboring cells
- I. Construction of **truncated field maps** with 'finite' longitudinal/horizontal extents **using Zgoubi**
 - Influence of the horizontal/longitudinal field map extent on the orbit
 - → Convergence to neighboring field maps that cover the entire orbit region
- II. Computation of the neighboring rotated fields with analytical expressions → No border anymore (called 'infinite' or 'perfect' map in the following slides)
 - → Neighboring field maps that cover the entire region of the orbit

ULB Fieldmap construction – Superposition

• After the construction, we use a Zgoubi Tosca option to superpose the field map with its neighboring field maps

ULB Lattice and methods – Limits of integration

Integration limits fixed by the **polar** character of the machine:

- We integrate between some *droites de coupures* (ddc) in Zgoubi, which are integration limits. It may not match the cartesian field map edges.
 - Entry oblique ddc: $tan(\theta)X+Y=0$
 - Exit vertical dcc: X = 280
- We have **no map overlapping** with well-chosen cut lines
- Perfect trajectory continuity between neighboring cells

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Linear beam dynamics

Closed orbit search & Computation of the tunes

- Complete recovery of the orbit for a horizontal extension of 180cm: we should see a convergence
- →No more change in the optics if we further increase the map

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- For smaller horizontal extension, we cross the closed orbit at different places -> More or less influence of neighboring cells

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ULB Closed orbits for different neighboring fieldmap extents

Orbit convergence with the field map extension from rotated truncated field maps computed with Zgoubi to « perfect » field map analytically computed

0.5

1

0

 The closed orbit given by S. Machida (IBG) is put at an arbitrary vertical coordinate to compare the shape and the vertical extension

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- Orbit convergence towards IBG orbit, which is the same as the orbit found with the perfect field map
- Vertical excursion of the orbit with an increasing neighboring field map extension

0.36 Horiz. Ext. = 85cm Horiz. Ext. = 90cm 0.34 Horiz, Ext. = 100cm Horiz. Ext. = 110cm Horiz. Ext. = 120cm 0.32 Horiz. Ext. = 130cm Horiz. Ext. = 140cm 0.3 Horiz. Ext. = 150cm Horiz. Ext. = 170cm Z(m) Horiz. Ext. = 180cm 0.28 Horiz. Ext. = 220cm closed orbit IBG 0.26 Extension = Infinite 0.24 0.22

2

1.5

S (m)

2.5

Convergence to the total influence of neighboring cells

ULB Linear beam dynamics - Computation of the tunes

Convergence of the tunes:

 Tunes computed with the perfect map are (0.24362, 0.119732), compared with IBG tunes: (0.243445, 0.12002)

Computation of the lattice functions:

 Correct computation of the lattice function with the Lebedev and Bogacz parametrization

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Non-linear beam dynamics

2D Dynamic aperture - 4D Dynamic aperture

Tune (qu, qv, nominal)	0.2434	45 / 0.120023	(0756555 / 0.1	20023)
ynamic aperture (normalised)		60 pi /	′ 70 pi	

Normalized dynamic aperture calculated with Zgoubi in eigenplanes on 100 turns

- Assumed ellipse-shaped phase space
- Small "stable" region, then a more diffuse region, but where the particles are not lost
- Islands \rightarrow fixed points of order 4

Normalized dynamic aperture calculated with Zgoubi in eigenplanes on 1000 turns

- The phase space is even more "diffuse"
- We still observe the islands (4th order fixed points)
- The 2D-dynamic aperture (assuming ellipse formula) in u and v spaces decrease

Normalized dynamic aperture calculated with Zgoubi in eigenplanes on 1000 turns

- Tracking of 3 particles into the islands
- The tune is 0.25, as expected

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Dynamic aperture – Non-linearity considerations

The magnetic field in vFFA is highly non-linear:

 Non-elliptical shapes in the linearly decoupled phase spaces → The "2D-DA" need to be refined to take into account the non-linearity

Dynamic aperture – Non-linearity considerations

The **magnetic field** in vFFA is **highly non-linear**:

 Non-elliptical shapes in the linearly decoupled phase spaces → The "2D-DA" need to be refined to take into account the non-linearity

- Different computation methods for the 2D-DA:
 - Assumed ellipse-shaped phase space
 - Computation of the average of linear invariants
 - Integration around the innermost points of the phase space area

Non-elliptical shape; 2D-DA

Dynamic aperture – Non-linearity considerations

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- Non-elliptical shapes in the linearly decoupled phase spaces → The "2D-DA" need to be refined to take into account the non-linearity
- The non-linearity couples the 'linearly decoupled' planes → Need to account for the non-linear coupling and to define/compute a more general "4D-DA"
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Non-elliptical shape; 2D-DA

ULB Dynamic aperture – Non-linearity considerations

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- Non-elliptical shapes in the linearly decoupled phase spaces → The "2D-DA" need to be refined to take into account the non-linearity
- The non-linearity couples the 'linearly decoupled' planes \rightarrow Need to account for the **non-linear** coupling and to define/compute a more general "4D-DA"
- Different <u>computation methods for the 2D-DA</u>:
 - Assumed ellipse-shaped phase space
 - Computation of the average of linear invariants
 - **Integration** around the innermost points of the phase space area

More robust definition for the 4D betatron motion:

- Literature review for the 2D and 4D betatron motion
- Computation of 2D-DA for different amplitude ratios between decoupled planes
- Computation of a "theoretical" 4D-DA (average distance) to compare working points
- Computation of a "practical" DA based on precautionary principle for better interpretation

Non-elliptical shape; 2D-DA

Non-linear coupling; "4D-DA"

Non-elliptical shape; 2D-DA

Different computation methods

Possible definitions for 2D-DA:

Assume ellipse-shaped space

 Computation of the Courant-Snyder invariant at a given point

Average of linear invariants

• Average of the emittances computed with the ellipse formula for each blue point in the cloud (each turn)

Integration over phase space

• Points cloud in $R - \theta$ form and integrate around the innermost points of the phase space area

Dynamic aperture 2D – FETS-vFFA

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#turns

- More robust definition for the 4D motion: Non-linear coupling; "4D-DA
 - Litterature review
 - 2D-DA (various decoupled planes amplitude ratios)
 - "Theoretical" 4D-DA (average distance)
 - "Practical" DA (precautionary principle)
- Examples of papers that extensively discuss Dynamic Aperture:
 - E. Todesco and M. Giovannozzi, 'Dynamic aperture estimates and phase-space distortions in nonlinear betatron motion', Phys. Rev. E 53(4), 4067 (1996).
 - M. Giovannozzi and E. Todesco, Numerical methods to estimate the dynamic aperture, Part. Accel. 54, 203 (1996).
 - S. Tygier, et al., 'The PyZgoubi framework and the simulation of dynamic aperture in fixed-field alternating-gradient accelerators', Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment, 775, pp. 15–26 (2015).
 - Giovannozzi, M., Scandale, W. and Todesco, E. (1996) 'Prediction of Long-Term Stability in Large Hadron Colliders', LHC Project Report 45-Rev.
 - Bojtár, L. (2020) 'Frequency analysis and dynamic aperture studies in a low energy antiproton ring with realistic 3D magnetic fields', *Physical Review Accelerators and Beams*, 23(10), p. 104002.

Dynamic aperture – General considerations

- « General » definition: stability domain particles bounded after N turns
- Dependent on N, and N depends on the application
 →How many turns do we need for the full acceleration cycle?
 →For hadron storage ring: predict « long-term » stability
- **Dependent on the motion** we look at:
 - **<u>2D betatron motion</u>**, without coupling
 - \rightarrow Stability domain = phase space **area** of initial conditions that survive N turns

 \rightarrow Border between stable/unstable motion (1D KAM torus) \rightarrow stability domain enclosed by the last connected stable invariant curve

- **<u>4D betatron motion</u>**, including coupling
 - \rightarrow Stability domain = phase space **volume** of initial conditions that survive N turns
 - ightarrow Volume may be irregular/have holes but generally not the case
 - ightarrow Dynamic aperture: radius of the hypersphere with the same volume as the stability domain

Dynamic aperture – 2D betatron motion

- Methods without averaging precautionary principle:
 - With the linear definition, the ellipse depends on the direction because of the phase distortion
 - \rightarrow Choose the smaller possible ellipse.

- Methods that give an average distance to stability border:
 - Direct integration: $\int \int \chi(x,p_x) \, dx \, dp_x \quad A_{\vartheta} = \int_0^{2\pi} \int_0^{r(\vartheta)} r \, dr \, d\vartheta = \frac{1}{2} \int_0^{2\pi} [r(\vartheta)]^2 \, d\vartheta \quad r_{\vartheta} = \left| \frac{A_{\vartheta}}{\pi} \right|^{1/2}$
 - Scan on all phase space variables needed
 - Integration over the dynamics $\frac{1}{2\pi} \int_0^{2\pi} [r(\vartheta)]^2 d\vartheta \rightarrow \lim_{N \to \infty} \frac{1}{N} \sum_{n=1}^N [r^{(n)}(\overline{\vartheta})]^2$
 - Fix θ and replace the space average with average over the N iterates
 - Uniform distribution of the phases of the iterates needed
 - Normal form method: $\rho(\overline{\vartheta}) = |\Psi(r(\overline{\vartheta}) \cos\overline{\vartheta}, r(\overline{\vartheta}) \sin\overline{\vartheta})|^2$
 - Compute the NL invariant with the truncated inverse conjugating function
 - Not valid close to a resonance

Dynamic aperture – 4D betatron motion

- Coupling between 'linearly decoupled' planes due to non-linearities
 - \rightarrow Ratio between the amplitudes in the different planes; Use of α such that $x = r \cos(\alpha)$ and $y = r \sin(\alpha)$
- «Fast DA estimates», commonly used : $\frac{x}{y} = 1$ ($\alpha = 45^{o}$) with $\theta_{\chi} = \theta_{y} = 0$

→Unprecise results, can not be used with non-negligible phase space distortion or 'ratio-dependent' dynamics

Methods that give an average distance to stability border:

Dynamic aperture – 4D motion – « roadmap »

Methods for 4D motion dynamic aperture computation in a strongly non-linear and coupled lattice

- Computation of 2D-DA (ε₁, ε₂) for different amplitude ratios between the decoupled planes to account for the non-linearity
 - Evolution of the $\epsilon_1 \ et \ \epsilon_2$ invariants as a function of α and computation of the average of ϵ_1 and ϵ_2 ; ϵ_1 and ϵ_2 can be related to measurable parameters with appropriate lattice functions (linear coupling).
- Computation of a "theoretical" 4D-DA estimate (average distance to the stability border) to compare working points
 - Direct integration: scan all the phase space variables
 - Integration over the dynamics

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- The 4th order islands of stability are taken into account, and a metric representing the « filling factor » is being defined to account for the topology of the phase space
- Computation of a "practical" DA based on precautionary principle for better interpretation
 - Launch an elliptical shape bunch \rightarrow Every particle needs to survive N turns
 - ightarrow Define the DA in the coupled space

ULB Summary

- Fieldmaps for the FETS-vFFA arctan lattice have been generated, and particles tracked with Zgoubi
 - There is an **important influence of neighboring cells** due to the significant cell ends residual fields
 - ightarrow The actual superposition of field maps is needed
- The detailed study of the linear transverse beam dynamics has been performed
 - The **influence** of the neighboring cell **field map extents** on the orbit and tunes have been studied, and **convergence** towards the 'infinite' map has been found
 - The tunes and lattice functions have been computed and are similar to those obtained with other codes
- An in-depth study of the non-linear dynamics of the lattice is in progress
 - The **2D normalized dynamic apertures** were calculated in the decoupled planes with Zgoubi, without non-linearities; **Islands of stability** appear in phase space.
 - Different **definitions and methods to calculate the DA** exist in the literature for **2D and 4D** betatron motion; it includes definitions based on a **phase space variable average** or **the 'precautionary principle'**.
 - Methods to study completely the 4D motion dynamic aperture in a strongly non-linear and coupled lattice have been suggested. An in-depth study of dynamic aperture in the full 4D phase space is still in progress.