Useful things to know about accelerators – part l

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Science and Technology Facilities Council

Accelerators - A Window on Nature

- Particle accelerators provide the source for most high energy physics experiments
 - Provide high luminosity, high energy beams for colliders
 - Provide high brightness beams for secondary particle production
 - Also key technology for life sciences, engineering, chemistry
- How do they work?
 - How can we get to high energy?
 - How can we keep the beam in the accelerator?
 - How can we get to high luminosity?
- What are the main HEP facilities in the world today?
- What might HEP facilities look like in the future?



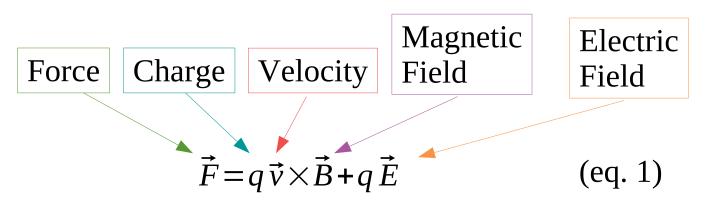
Accelerator Components

- Most accelerators share similar components
- Main components of an accelerator
 - Bending dipoles
 - Focussing quadrupoles
 - Acceleration RF cavities
- Also
 - Vacuum
 - Diagnostics
 - Targets for secondary particle production
- First Lecture: Derive basic theory of accelerator physics
- Second Lecture: Discuss accelerator equipment and techniques



Lorentz force law

Fundamental equation for particles moving through fields

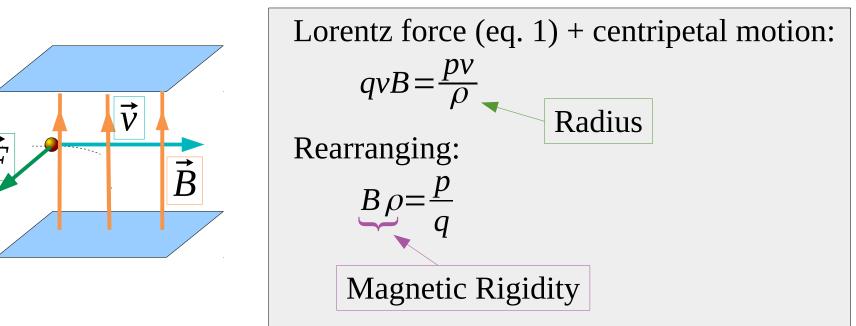


- Magnetic force is perpendicular to velocity
 - Magnetic field conserves energy
- Electric force is weaker by factor velocity
 - Magnets are better for bending and focussing



Magnetic Rigidity and Bending

- Simplest magnet "dipole"
 - Uniform magnetic field perpendicular to beam direction



- Constant force \rightarrow constant curvature \rightarrow circular motion
- Magnetic rigidity parameterises momentum
- Charge-to-mass ratio important when accelerating multiple particle species



Worked example – LHC

- If we wanted to accelerate, say, 7 TeV particles, what bending radius is required?
- Maximum dipole field around 8.3 T

$$B\rho = \frac{p}{q}$$

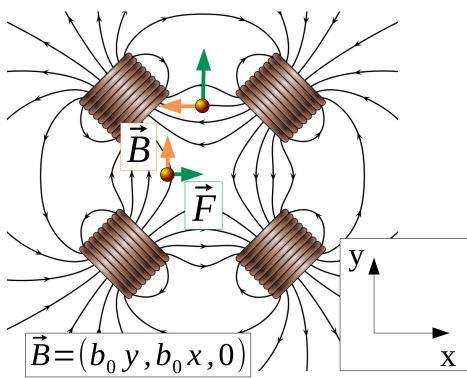
$$\rho = \frac{p}{qB} = \frac{7}{0.3 \times 8.3} = 2.8 \text{ km}$$

- Nb: LHC radius ~ 4.1 km
 - Need space for detectors, etc



Quadrupole magnets

- If we only had bending magnets, particles would soon be lost from the accelerator
- Need to keep the particles in the accelerator using focussing elements
 - Usually use quadrupoles
- Field stronger away from beam centre
 - Like a spring or pendulum
 - Simple harmonic motion
- Overall focussing by alternating the gradient





Quadrupole field – horizontal (1)

- For a particle moving near to the z-axis $\vec{F} = q \vec{v} \times \vec{B} + q \vec{E}$ $\vec{B} = (b_0 y, b_0 x, 0)$
- Considering only p_x for now

$$\frac{dp_x}{dt} = q \frac{dz}{dt} B_y$$

Use the chain rule

$$\frac{dp_x}{dt} = \frac{dp_x}{dz}\frac{dz}{dt}$$

Combining these equations:

$$\frac{dp_x}{dz} = q b_0 x$$



Quadrupole field – horizontal (2)

$$\frac{dp_x}{dz} = q b_0 x \qquad \textcircled{\circ}$$

Definition of x-component of momentum

$$p_x = m \gamma v_x = m \gamma \frac{dz}{dt} \frac{dx}{dz} = p_z \frac{dx}{dz}$$

Substitute this definition into 🙂 gives

$$p_z \frac{d^2 x}{dz^2} = q b_0 x$$

 Rearrange and wrap up constant terms in focussing strength k

$$\frac{d^2x}{dz^2} - kx = 0$$



Quadrupole field – vertical

Lorentz force law with quadrupole field definition

$$\frac{dp_y}{dt} = -q b_0 v_z y$$

Use chain rule and eliminate v₂

$$p_z \frac{d^2 y}{dz^2} = -q b_0 y$$

 Rearrange and wrap up constant terms in defocussing strength k

$$\frac{d^2 y}{dz^2} + k y = 0$$



Solutions

Motion is governed by $\frac{d^2 x}{dz^2} - kx = 0 \qquad \frac{d^2 y}{dz^2} + ky = 0$

- This is simple harmonic motion solutions are of form $x = x_0 \cos(\sqrt{k} z) + \frac{dx_0}{dz} \frac{1}{\sqrt{k}} \sin(\sqrt{k} z)$
- Taking derivative

$$\frac{dx}{dz} = -x_0 \sqrt{k} \sin(\sqrt{k} z) + \frac{dx_0}{dz} \cos(\sqrt{k} z)$$

For y

$$y = y_0 \cosh(\sqrt{k} z) + \frac{dy_0}{dz} \frac{1}{\sqrt{k}} \sinh(\sqrt{k} z)$$
$$\frac{dy}{dz} = y_0 \sqrt{k} \sinh(\sqrt{k} z) + \frac{dy_0}{dz} \cosh(\sqrt{k} z)$$



Transfer Matrix

Just thinking about x, the particles move according to

$$x_1 = x_0 \cos(\sqrt{k} z) + \frac{dx_0}{dz} \sin(\sqrt{k} z)$$
$$\frac{dx_1}{dz} = -x_0 \sqrt{k} \sin(\sqrt{k} z) + \frac{dx_0}{dz} \sqrt{k} \cos(\sqrt{k} z)$$

We can rewrite this as a matrix

$$\begin{pmatrix} x_1 \\ \frac{dx_1}{dz} \end{pmatrix} = \begin{pmatrix} \cos(\sqrt{k}z) & \frac{1}{\sqrt{k}}\sin(\sqrt{k}z) \\ -\sqrt{k}\sin(\sqrt{k}z) & \cos(\sqrt{k}z) \end{pmatrix} \begin{pmatrix} x_0 \\ \frac{dx_0}{dz} \end{pmatrix}$$

This matrix is known as the quadrupole's transfer matrix

$$\underline{u}_1 = \boldsymbol{M}_{\mathbf{01}} \underline{u}_0$$





- Exercise what is the transfer matrix for a drift space, that is a region with no fields at all?
 - What is the force acting on the particle?
 - What is x(z) in terms of dx₀/dz and x₀
 - What is dx/dz in terms of dx₀/dz
 - Now write that as a matrix



- Exercise what is the transfer matrix for a drift space?
 - What is the force acting on the particle?
 - No force
 - What is x(z) in terms of dx₀/dz and x₀

$$x = x_0 + \frac{dx_0}{dz}z$$

What is dx/dz in terms of dx₀/dz

$$\frac{dx}{dz} = \frac{dx_0}{dz}$$

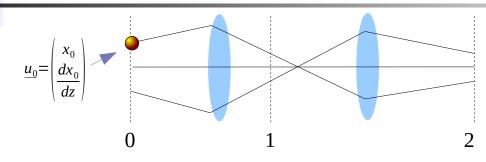
Now write that as a matrix

$$\begin{pmatrix} x \\ \frac{dx}{dz} \end{pmatrix} = \begin{pmatrix} 1 & z \\ 0 & 1 \end{pmatrix} \begin{pmatrix} x_0 \\ \frac{dx_0}{dz} \end{pmatrix}$$





Transfer Lines



- Transfer matrix defines transport through a region
- Transfer matrices can be combined by multiplication
- Say we have transfer matrices like:

$$\underline{u_1} = \boldsymbol{M}_{01} \underline{u_0}$$
$$\underline{u_2} = \boldsymbol{M}_{12} \underline{u_1}$$

Then

$$\underline{u}_2 = \boldsymbol{M}_{12} \boldsymbol{M}_{01} \underline{u}_0$$

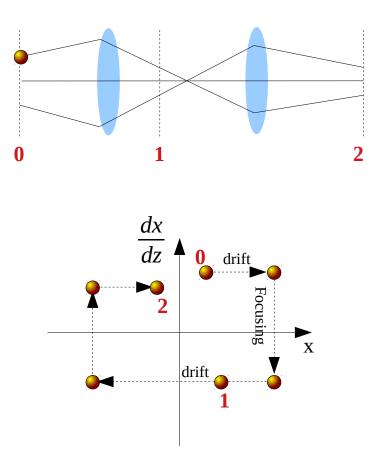
• i.e. we can define a combined transfer matrix like

$$M_{02} = M_{12} M_{01}$$

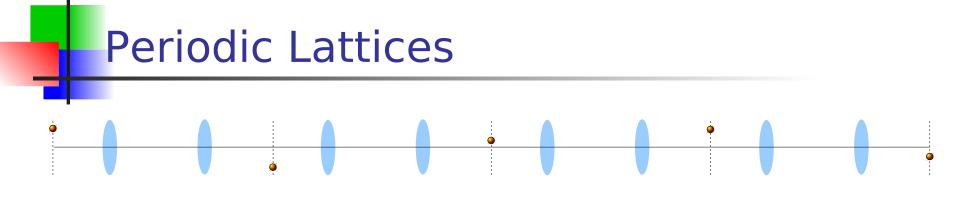


Phase space

 Another instructive way to look at beam optics is by considering the phase space







 Following n identical cells or turns in a ring with one-turn matrix M

$$\underline{u}_n = \boldsymbol{M}^n \underline{u}_0$$

Rewrite

 $M = I \cos \mu + J \sin \mu$ $J = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} \text{ with } \gamma \beta - \alpha^2 = 1 \text{ and } I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$ So $J^2 = -I$

And

$$M^n = I \cos(n\mu) + J \sin(n\mu)$$



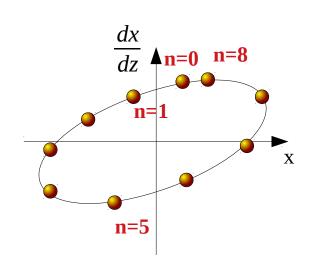
Periodic Lattices

What does this mean?

 $\boldsymbol{M}^{\boldsymbol{n}} = \boldsymbol{I}\cos(n\,\mu) + \boldsymbol{J}\sin(n\,\mu)$

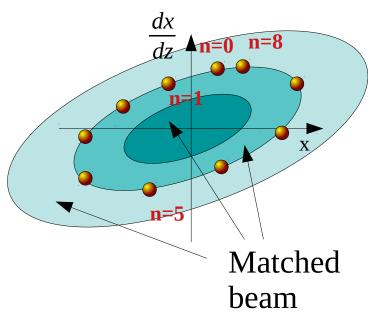
- Particles move around an ellipse in phase space if Trace(M) < 2
- µ is the "phase advance"
 - Sometimes use "tune" ... $2\pi v = \mu$
- α, β and γ are "Twiss parameters"
 - Tell us the alignment of the ellipse
- Each particle sits on ellipse area ε the particle's amplitude





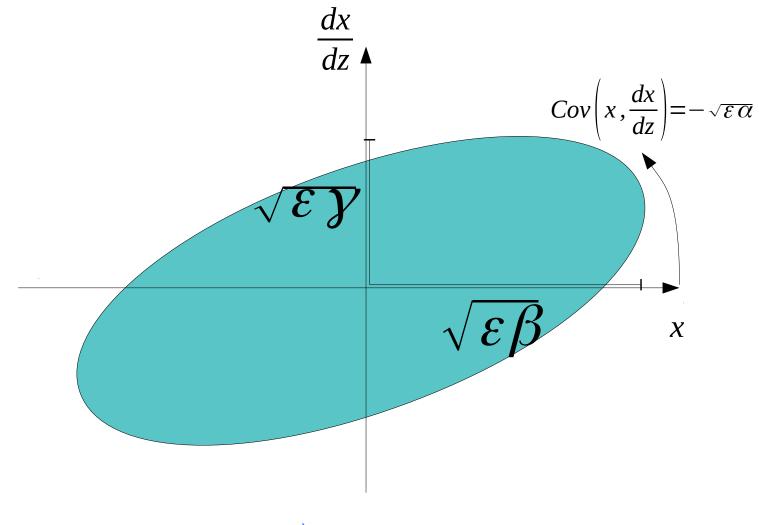
Periodic Lattices and beams

- Beam is composed of many particles
 - Particles occupy a region in phase space
- "Emittance" is area occupied by the entire beam
- Sometimes classify "RMS emittance"
 - Area occupied by ellipse 1 RMS distance from beam centre
- Low emittance is crucial for
 - High luminosity
 - Low losses





Beam ellipse





Emittance Growth

- Ideally emittance is conserved, but this is not always the case
- Long list of effects that can cause emittance growth
 - Beam mismatch
 - Scattering off residual gas
 - Scattering off particles in the same beam
 - Scattering off particles in other beams (e.g. in collider)
 - Space charge
 - Resonances



Transverse Space Charge 1

- Consider a circular beam of radius a having uniform density
- $\rho(r) = q \frac{I}{\beta_{rol} c \pi a^2} \qquad r < a$ Quote field around a cylinder of charge/curre $E(r) = \frac{1}{2 \pi \varepsilon_0} \frac{I}{\beta_{rel} c} \frac{r}{a^2}$ $B_{\phi}(r) = \frac{1}{2 \pi \varepsilon_0} \frac{I}{c^2} \frac{r}{a^2}$ Apply Lorentz force law $\vec{F} = q \vec{v} \times \vec{B} + q \vec{E}$ $F_{r} = q \vec{v} \times \vec{B} + q \vec{E} = \frac{1}{2 \pi \varepsilon_{0}} \frac{r}{a^{2}} \left(\frac{I}{\beta_{rel} c} - \frac{I}{\beta_{rel} c} \beta_{rel}^{2} \right) = \frac{1}{2 \pi \varepsilon_{0}} \frac{r}{a^{2}} \left(\frac{I}{\gamma^{2} \beta_{rel} c} \right)$



Transverse Space Charge 2

Force is defocusing

$$\frac{d^2 x}{dz^2} - (k - K_{sc}) x = 0 \quad \text{with}$$

Treat SC as a perturbation

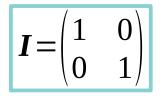
$$M_{p} = M M_{sc}$$
$$M = I \cos \mu + J \sin \mu$$
$$M_{sc} = \begin{pmatrix} 1 & 0 \\ -K_{sc} & 1 \end{pmatrix}$$

- Change of beam size (β)
- Change of tune
 - Drive the beam onto resonances \rightarrow ruin the acceptance
- Tune \rightarrow look at Trace of M_p

$$Tr(\boldsymbol{M}_{p}) = 2\cos(\mu) + \alpha\sin(\mu) - \alpha\sin(\mu) + \beta K\sin(\mu)$$



$$K_{sc} = \frac{1}{2 \pi \varepsilon_0} \frac{1}{a^2} \left(\frac{I}{\gamma_{rel}^3 \beta_{rel}^2 c} \right)$$



$$\boldsymbol{J} = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix}$$

Transverse Space Charge 3

- Consider just the $trace(M_p)$ $Tr(M_p)=2\cos(\mu)+\beta K\sin(\mu)$
- Consider compound angle formula $\cos(\mu + \delta\mu) = \cos(\mu)\cos(\delta\mu) + \sin(\mu)\sin(\delta\mu)$ $\cos(\mu + \delta\mu) \simeq \cos(\mu) + \sin(\mu)\sin(\delta\mu)$
- Looking at the tune

$$\delta v = \frac{\delta \mu}{2\pi} = \frac{\beta K}{4\pi}$$
$$\delta v = \frac{r_0 N}{2\pi \epsilon \beta_{rel}^2 \gamma_{rel}^3}$$

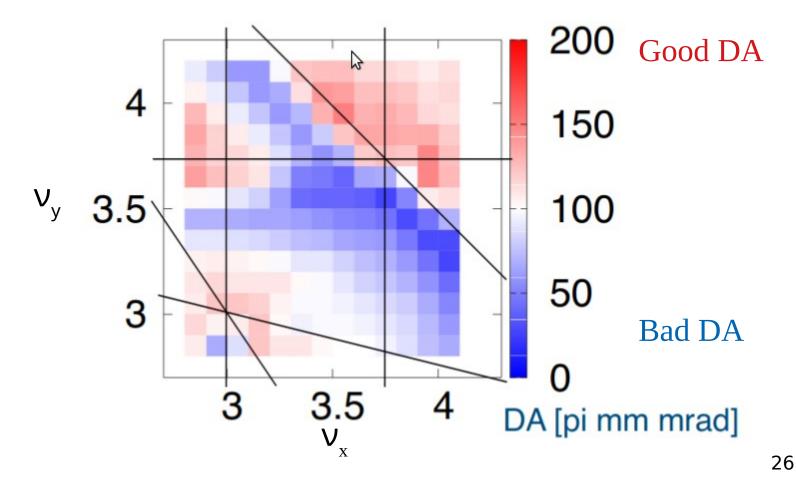
$$K_{sc} = \frac{1}{2 \pi \varepsilon_0} \frac{1}{a^2} \left(\frac{I}{\gamma_{rel}^3 \beta_{rel}^2 c} \right)$$

$$\sigma(x) = \sqrt{\beta \epsilon}$$



Resonances

- The beam does not behave well when $j v_x + k v_y = N$
 - Imperfections in the field get amplified
 - Depending on the surrounding lattice (i.e. the imperfections)



Emittance Reduction (Cooling)

- Several techniques to reduce emittance
 - Synchrotron radiation cooling
 - Stochastic cooling
 - Laser cooling
 - Electron cooling
 - Ionisation cooling
- Fundamental principle is to remove "heat" from the beam using a neighbouring heat sink
 - Comoving electron beam → electron cooling
 - Comoving laser → laser cooling
 - Emission of synchrotron radiation
 - Photon emission caused by (principally) electrons bending in magnetic field





- What is behaviour of particles in phase space if
 - Trace(M) < 2
 - Trace(M) = 2
 - Trace(M) > 2



- What is behaviour of particles in phase space if
 - Trace(M) < 2
 - Motion is an ellipse
 - Trace(M) = 2
 - x → +/- x
 - Trace(M) > 2
 - Motion is a hyperbola

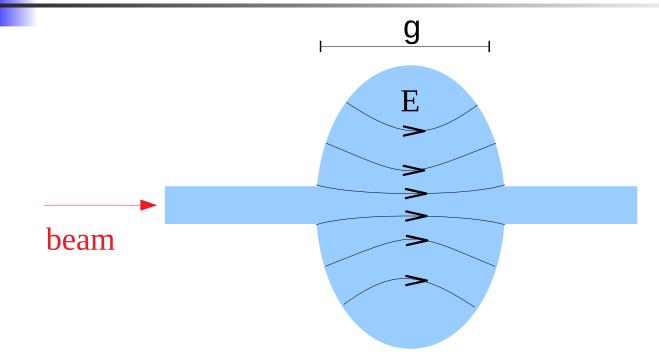


Longitudinal Dynamics

- So much for transverse motion (i.e. x and y planes)
- What about energy and acceleration?
- Electrostatic acceleration limited by breakdown potential
 - Change in energy is given by voltage differential
 - High voltage differentials cause breakdown (sparks)
 - Practically limits electrostatic acceleration to few MeV
- To accelerate beyond MeV require oscillating electric field
- RF Cavities



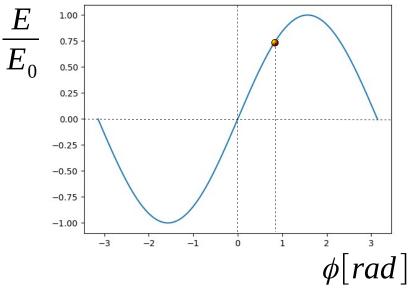
RF cavity field



- RF cavity holds a resonating EM wave
- Recall Lorentz force law $\vec{F} = q \vec{v} \times \vec{B} + q \vec{E}$
- Force is in direction of motion energy changes!



RF cavity field



In RF cavity

 $\vec{E} = E_0 \sin(\omega t + \phi)$

- Energy change of synchronous particle crossing at ϕ_s $\delta W = qTgE_0 \sin(\phi_s)$
 - T is factor to allow for phase to vary a bit during crossing
 - g is the gap length

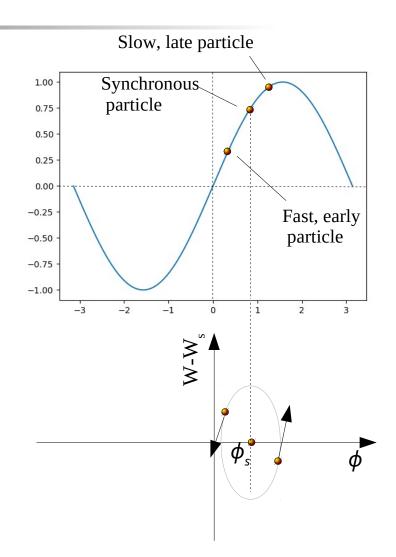


Phase stability

 Particle crossing at phase φ relative to synchronous particle

 $\delta W = q T g E_0 \sin(\phi + \phi_s)$

- Particle arriving early
 - Fast
 - t negative
 - Gets smaller energy kick
 - Ends up relatively slower
- Particle arriving late
 - Slow
 - t positive
 - Gets bigger energy kick
 - Ends up relatively faster
- Phase stability!

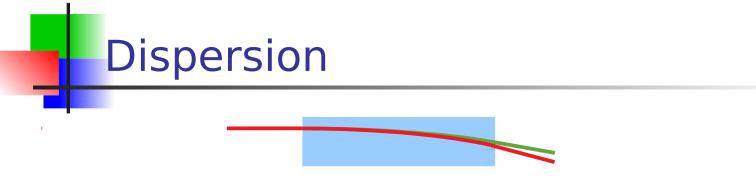




Dealing with momentum spread

- Momentum spread introduces a few effects
 - Dispersion
 - Chromaticity
 - Momentum compaction
- Dispersion:
 - Off-momentum particles follow a different trajectory
- Momentum compaction (rings):
 - Different path length yields different time of flight
- Chromaticity:
 - Off-momentum particles get a different focussing strength





Recall the definition of magnetic rigidity

$$B\rho = \frac{p}{q}$$

- Particles having different momentum (p) get different radius of curvature
 - Introduce dispersion D

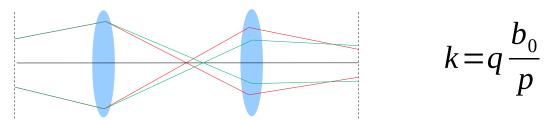
$$D = p \frac{dx}{dp}$$

 Which is another optical function that we must make periodic



Chromaticity

 Chromaticity arises because quadrupoles focus differently for different momenta



- This often limits the degree of focussing at a collision point
 - Limits luminosity
- Can deliberately enhance/reduce chromaticity by
 - Introduce a dispersion
 - Using a magnet with variable focussing strength across the aperture - "sextupole"





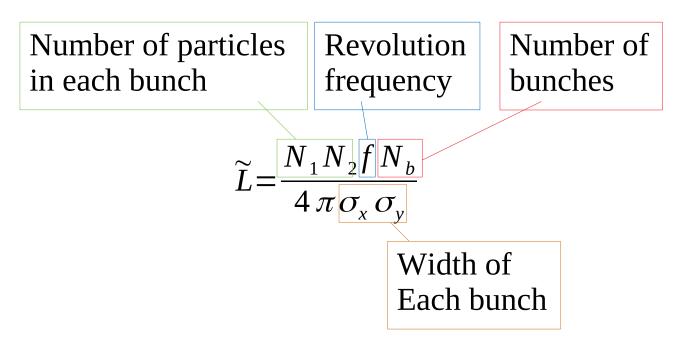
Review

- Dipoles are used to bend a beam rigidity is $B\rho = \frac{p}{2}$
- Quadrupoles are used to focus a beam: $k = q \frac{b_0}{r}$
- Beam in each of x and y can be characterised by 3 Twiss parameters and an emittance
- Lattices can be characterised by a phase advance
- RF cavities are used to accelerate the beam
- Introducing momentum spread, one can also define a dispersion (and its derivative with respect to z)



Finally... luminosity

- Luminosity defines the number of interactions in a collider per unit time for a given cross section
- Luminosity will increase if
 - Beam is narrower
 - Current is higher





What dictates luminosity?

$$\widetilde{L} = \frac{N_1 N_2 f N_b}{4 \pi \sigma_x \sigma_y}$$

- Typically
 - Number of particles → space charge
 - Revolution frequency → ring circumference
 - Number of bunches → RF frequency
 - Beam width $\rightarrow \sqrt{\varepsilon \beta}$
 - Emittance (cooling?)
 - Twiss beta (final focus and chromaticity)



Next lecture...

- Accelerator equipment
- Types of accelerator
- Current facilities
- Future facilities

