

# Event Generators for Collider Physics

Christian Gütschow

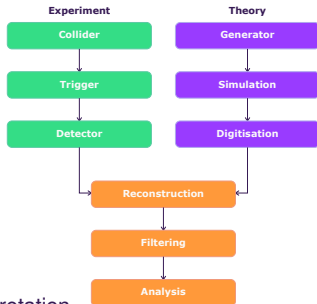
RAL Advanced Graduate Lectures

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# Introduction

# Overview

- Event generation is the simulation of the fundamental particle-collision process.
  - Theory predicts amplitudes (not observables!), detectors measure tracks and calorimeter deposits.
  - Monte Carlo **event generators bridge this gap**
- Indispensable in HEP – from experiment design to interpretation.
- This lecture focuses on the parton-shower generators that produce fully exclusive hadron-level events, ready to be fed into the detector simulation.
- These tools are often treated as black boxes – rarely a good idea!
- A little extra understanding goes a long way, so let's see what we can do . . .



## The dream (and why it fails)

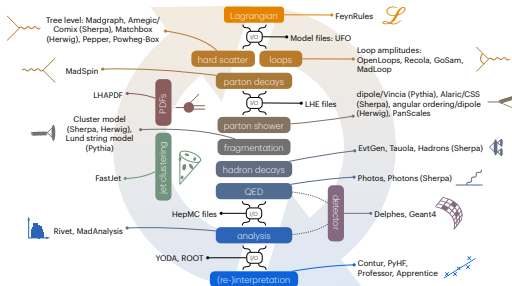
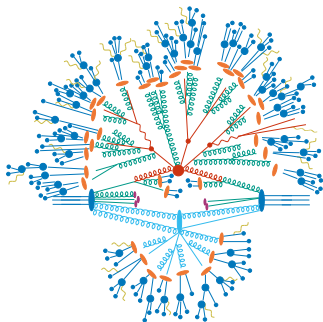
- In an ideal world, we'd compute the full event from first principles:
  - Write down the squared amplitude from initial beams to every final-state particle hitting the detector
  - Integrate over phase space, sample, done.
- Quantum field theory gives us the framework – so why not **just do it**?

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- In an ideal world, we'd compute the full event from first principles:
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- Quantum field theory gives us the framework – so why not **just do it**?
  
- Three obstacles force us back to reality:
  - ✗ **Non-perturbative effects**: the transition to and from hadrons isn't calculable from first principles (beyond even lattice QCD).
  - ✗ **Perturbative tractability**: the number of terms grows factorially – we're limited to  $< 10$  final-state partons, but need hundreds.
  - ✗ **Phase-space dimensionality**:  $\sim 3N_{fs} = \mathcal{O}(1000)$  dimensions, with probability concentrated in unknown islands.

# Principled approximation

 [arXiv:2605.16036]



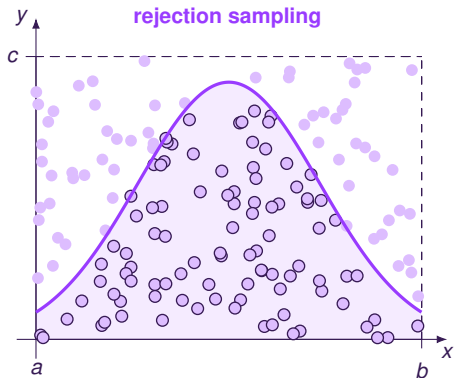
→ But good approximations are not far away:

- Most partonic multiplicity can be emulated by a Markov-chain **parton shower** – iterated splittings in place of full correlations.
- NP effects yield to modelling too: factorised beam structure, beam remnants, multiple interactions, hadronisation.

✓ Realistic events become possible – via a **modular, multi-stage chain**, each stage addressing one piece of the physics.

# From theory to partons

# Monte Carlo integration

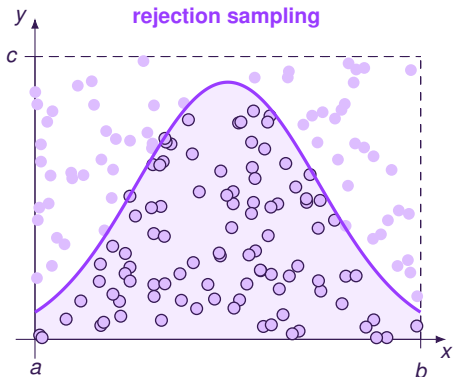


● rejected (miss):  $y > f(x)$

○ accepted (hit):  $y < f(x)$

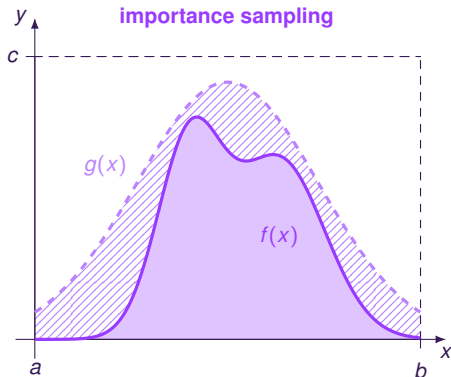
$$\int_a^b f(x) dx \approx \frac{N_{\text{hits}}}{N_{\text{total}}} \times c(b-a)$$

# Monte Carlo integration



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$$\int_a^b f(x) dx \approx \frac{N_{\text{hits}}}{N_{\text{total}}} \times c(b-a)$$



construct simple  $g(x) > f(x)$  with  $\frac{f(x)}{g(x)} < 1$   
draw  $x$  from  $g(x)$ , reject if  $\frac{f(x)}{g(x)} < R \in [0, 1]$

VEGAS decomposes into multiple channels:

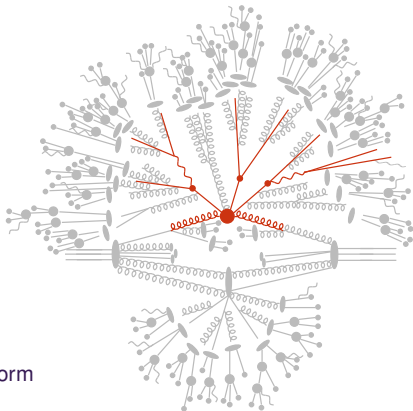
$$g(x) = \sum_i a_i g_i(x)$$

# Hard scattering and phase space

- Core of all event generation is an integral of PDFs and partonic matrix elements over phase space:

$$d\sigma_n^{\text{LO}} = f_a^{h_1}(x_a, \mu_f) f_b^{h_2}(x_b, \mu_f) \times \frac{1}{2\hat{s}} |\mathcal{M}_{ab \rightarrow n}(\Phi_n)|^2 \times d\Phi_n(\hat{s}; p_1, \dots, p_n)$$

- Integral is done by Monte Carlo sampling
  - integration uncertainty scales as  $1/\sqrt{N}$
- Easy to spend all CPU on atypical points!
- Better strategy to perform a Jacobian transform and use importance/adaptive sampling or even ML to 'learn the space'.
- MC integration can take months, unless massively parallelised.
  - Save integration as 'gridpack' for event generation.



## Beyond leading order

→ Extra partonic emissions at tree-level increase the final state multiplicity and change the event kinematics directly.

→ An NLO cross-section has three parts:

$$d\sigma_n^{\text{NLO}} = d\sigma_n^{(0)} + d\sigma_n^{(1)} + d\sigma_{n+1}^{(0)}$$

→ The LO prediction for the process  $\mathcal{P}_{n+1}$  is exactly the real-emission contribution entering the NLO correction to  $\mathcal{P}_n$ .

→ Kinematic cuts (hatched region) are required to ensure that this extra radiation remains resolved.

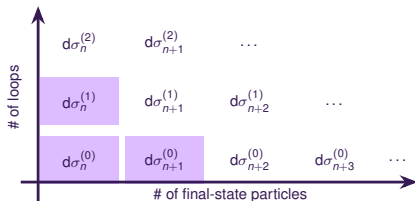
→ Many NLO frameworks, but only few automated

→ aMC@NLO and Sherpa fully automated

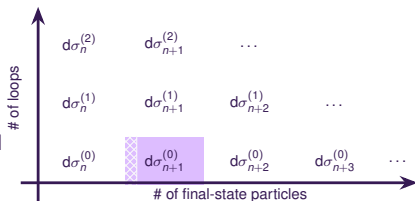
→ POWHEG-BOX is a framework to assist manual implementation

→ Virtual terms from dedicated libraries (OpenLoops, BlackHat, NJETS, GoSAM, ...)

$\mathcal{P}_n @ \text{NLO}$



$\mathcal{P}_{n+1} @ \text{LO}$



## Splitting functions and resummation

- Partons are not jets: fixed-order partonic events are only comparable to data for IR-safe observables.
- Soft and collinear phase-space ( $1 \text{ GeV} < p_T \lesssim 20 \text{ GeV}$ ) receives large resummation corrections from multiple QCD emissions.
- Analytic resummation is process-specific but factorises in the collinear limit: cross-section for process  $\sigma_0$  with parton  $i$  to be accompanied by a collinear parton  $j$  with momentum fraction  $z$ :

$$d\sigma \approx \sigma_0 \sum_{\text{partons } i} \frac{\alpha_s}{2\pi} \frac{d\theta^2}{\theta^2} dz P_{ji}(z, \phi) d\phi$$

where  $\theta$  is the angle between  $i$  and  $j$  and

$$P_{qq}(z) = C_F \frac{1+z^2}{1-z}$$

$$P_{gq}(z) = C_F \frac{1+(1-z)^2}{z}$$

$$P_{gg}(z) = C_A \frac{z^4 + 1 + (1-z)^4}{z(1-z)}$$

$$P_{qg}(z) = T_R(z^2 + (1-z)^2)$$

- These are the spin-averaged QCD collinear splitting functions, or DGLAP kernels.

## Splitting functions and resummation

→ If the emission probability is proportional to

$$C \frac{d\theta}{\theta} dz$$

then in a sequence of **strongly ordered** soft/collinear radiations, the probability is given by

$$\prod_{i=0}^n C \frac{d\theta_i}{\theta_i} dz_i$$

→ Quiz time: If integration is the natural continuous limit of a sum:

$$\sum \rightarrow \int$$

What's the equivalent continuum limit of a product?

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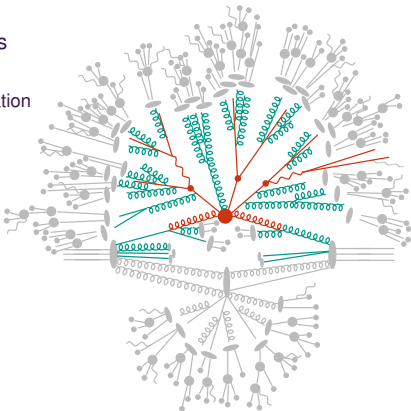
What's the equivalent continuum limit of a product?

$$\prod = \exp \left\{ \sum \ln \right\} \rightarrow \exp \left\{ \int \ln \right\}$$

- Every scale-step has a differential probability of no-emission; the **Sudakov exponentiation** multiplies them together.

# Parton shower

- Parton showers are Markov chain algorithms based on the QCD splitting functions
  - process-independent, approximate resummation
- From splitting functions, one can calculate the probability of no emission above  $q_0$ :  
$$\Delta(q_0, Q) = \exp \left\{ -C' \int^Q \frac{d\theta}{\theta} \int dz P_{ji}(z) \right\}$$
The famous **Sudakov form factor**.
- Main showers used:
  - Pythia8 (most widely used shower model in ATLAS/CMS)
  - Herwig7 (angularly ordered shower with coherent soft emissions)
  - Sherpa's CSS (dipole-style shower used internally)
- Controls event topology and jet multiplicities



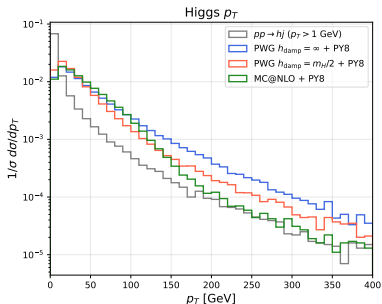
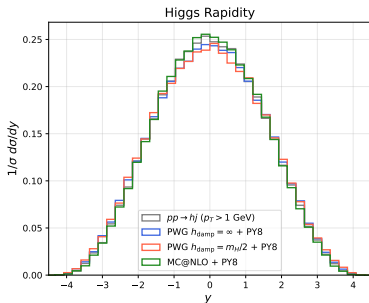
# Matching

- NLO includes the Born, one-loop, and one-real emission amplitudes (and all their interferences), and smoothly matches the real emissions between shower and ME:

$$\sigma^{\text{NLO}} = \int d\Phi_n [B(\Phi_n) + V(\Phi_n)] + \int d\Phi_{n+1} R(\Phi_{n+1}),$$

- Main benefit: NLO scale stability – ensures  $n + 1$  emission shapes now also stable without disrupting formal accuracy of parton shower.
  - rule of thumb: shapes from real, normalisation from virtual
- Matching methods split the  $(n + 1)$ -particle real-emission correction  $R$  into an infrared ( $S$ ) and a hard ( $H$ ) part:  $R = S + H$ .
- MC@NLO defines a **subtractive** procedure for  $H = R - S$ , identifying  $S$  with the splitting kernels
  - process specific, prone to negative weights
- POWHEG uses a **multiplicative** procedure to define  $S = FR$  such that  $H = (1 - F)R$  where  $F$  is defined to approach 1 in the infrared-singular limits.
  - shower independent, closer to all-positive weights (but not actually positive only)

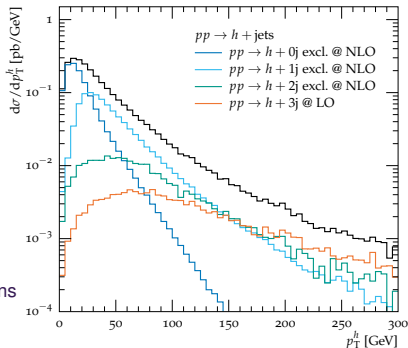
# Comparing matching schemes in Higgs production



- Higgs rapidity already well described by LO  $pp \rightarrow hj$ , no sensitivity to extra emissions, but Higgs needs to recoil against a jet for non-zero  $p_T$ .
- At low energies, matched distributions feature the so-called “Sudakov peak”, corresponding to the recoil against shower emissions.
- Beyond this, MC@NLO begins to follow the  $pp \rightarrow hj@LO$  calculation, POWHEG is very sensitive to the choice of  $h_{\text{damp}}$  in  $F = h_{\text{damp}}^2 / (p_T^2 + h_{\text{damp}}^2)$ 
  - can lead to significant effects away from the shower-dominated region

# Multi-jet merging

- Consistently combine matrix-elements for a range of final-state multiplicities.
  - Let the shower freely populate the soft region below a **merging cut**  $Q_{\text{cut}}$ , leaving the hard domain above this cut to the matrix elements.
  - Cluster the partons according to the inverse splitting probabilities of the shower.
  - Evaluate no-emission probability between successive splittings, making these contributions exclusive in the resolved emissions.
- CKKW(-L) and MLM schemes designed to replace the shower's collinear splitting functions with proper matrix elements in the relevant (hard) phase space
- Extended to automatic merging of many NLO and/or LO multi-leg matrix elements in MEPS@NLO (Sherpa) and FxFx (MG5\_aMC@NLO).
- **Huge book-keeping exercise!**

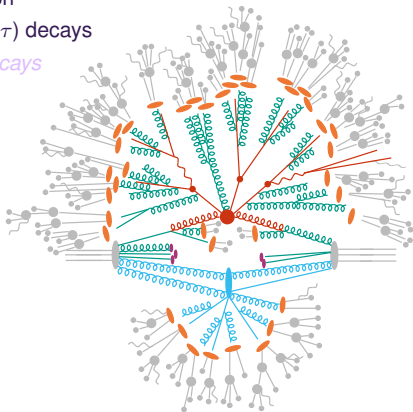


# Non-perturbative modelling

# From partons to hadrons

Fragmentation =  $\left\{ \begin{array}{l} \text{(Final-state) parton showers} \\ \text{Hadronisation} \\ \text{Prompt hadron (and } \tau \text{) decays} \\ \textit{Non-prompt decays} \end{array} \right.$

- Hadronisation = non-perturbative transition from partons to hadrons around  $\Lambda_{\text{QCD}}$
- Underlying event = beam remnants + multiple parton interactions (MPI)
- Main models used:
  - Pythia8: Lund string model + MPI model (interleaved with ISR/FSR)
  - Herwig7: Cluster model + MPI
  - Sherpa: Cluster hadronisation + UE model; plugin to Pythia8 for Lund-style hadronisation
- Requires **model tuning** to MB/UE observables (ATLAS A14/AZ, CMS CP5/CH3 etc.)



# Hadronisation models



## → Lund string (Pythia):

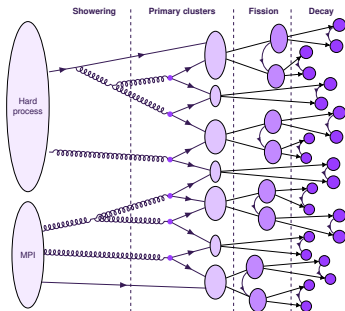
- Inspired by linear scaling of QCD potential at large distances
- Break colour strings to produce new quark pairs; gluons form kinks in strings
- Lorentz invariance and left-right symmetry give Lund Symmetric Fragmentation Function
- Kinematics well-described, but flavour (especially baryons) not natural

## → Cluster hadronisation (Herwig, Sherpa):

- Colour preconfinement, seen in colour-connected neighbour parton mass spectrum
- Non-perturbative  $g \rightarrow q\bar{q}$ , then cluster colour singlets: requires finite gluon constituent mass
- Clusters treated as meson resonances

## → Both models also contain colour reconnection heuristics

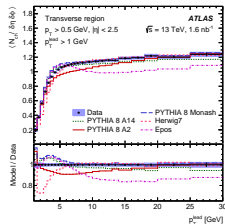
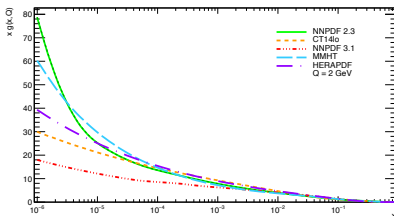
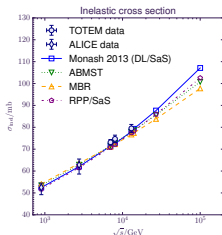
## → Requires **model tuning!**



# Underlying event and MPI

🔗 [arXiv:1808.07224]

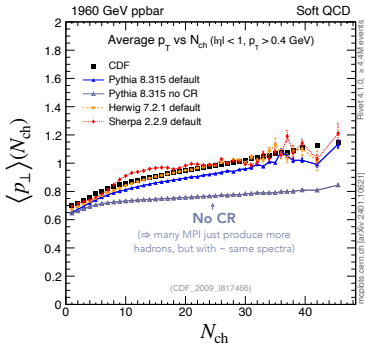
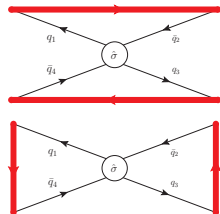
🔗 [arXiv:1701.05390]



- Inelastic cross-section rises strongly with centre-of-mass energy due to low- $x$  PDFs
- We can think of two colliding hadrons as two beams of partons colliding
  - the probability for a high- $p_{\perp}$  interaction is low ( → calculate using perturbation theory)
  - the probability for low- $p_{\perp}$  scatterings is very large (dominated by  $t$ -channel gluon exchanges)
- Sample  $\langle n_{\text{parton-parton}} \rangle \sim \sigma_{\text{parton-parton}} / \sigma_{\text{hadron-hadron}}$  interactions from a Poisson
- Effective matter overlap between the two hadrons characterised by “overlap factors”
  - small  $\langle n_{\text{parton-parton}} \rangle$  in peripheral collisions, large  $\langle n_{\text{parton-parton}} \rangle$  in central collisions
- **Pedestal effect** in UE as “minimum bias” sample transitions to a “maximum bias” one
  - Not possible to be more central than fully central!
- Requires **model tuning!**

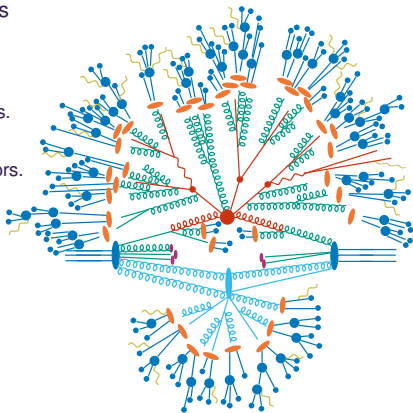
# Colour reconnection

- Correlations in colour space could cause non-trivial rearrangements of the hadronising systems.
- Most modern CR models seek to minimise a measure of the total potential energy of the hadronising system.
  - e.g. total string length in string CR models or (summed) cluster masses in cluster CR models
- Such minimisations typically result in fewer but more energetic hadrons.
- Shower radiation patterns could also be induced by colour connections to partons outside the resonance-decay systems (e.g. top quarks).



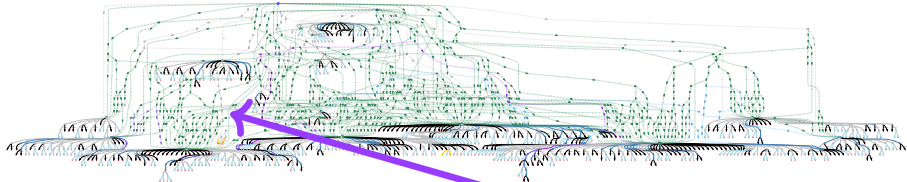
# Hadron/tau decay modelling

- Accurate modelling of hadron decays can be complex, but in many cases sufficiently accurate results can be obtained with tabulated decay modes and branching ratios.
  - Decay kinematics sampled (in the decaying particle's rest frame) from the appropriate dimensionality of phase-space configurations.
  - Models with more specific matrix-element influences can be implemented as form factors.
- Main tools used:
  - Pythia/Herwig/Sherpa's internal decayers
  - EvtGen (used for  $B$ -hadron decays with custom decay tables on top of Pythia/Herwig-based samples)
  - Tauola (optional for  $\tau$  decays with polarisation modelling used on top of Pythia/Herwig-based samples)
- Spin correlations sometimes retained depending on interface

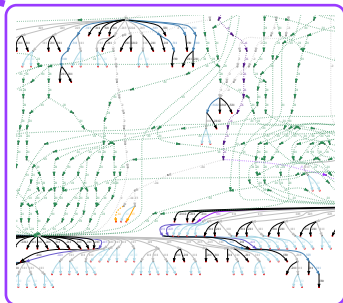


# Using Monte Carlo events

# The HepMC event graph

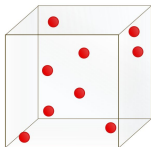


- A MC event graph showing all the particles (lines) and interactions (vertices) in the Pythia event record.
- Guess the event type?  
(Hint: 2212 = proton, 1–6 = quarks, 11–16 = leptons, 21–25 = bosons)
- Uff . . . how to safely analyse *that*?
  - There is no unambiguous truth in quantum mechanics!
  - Event records are half physics, half debug info – and zero indication of amplitude interference!
  - Kinematic frames aren't defined until the final state, momentum not necessarily conserved at all vertices.
  - **Beware!**

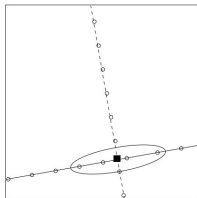
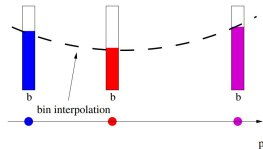


# Model tuning

- Tuning was historically a mix of brute force and expert intuition.
- Professor/Apprentice accelerates convergence by building surrogate models:
  - Sample parameter points  $p_n$  from a hypercube/sphere.
  - Generate MC run-sets (for beams, processes, etc.) at each  $p_n$ .
  - Run jobs in parallel on batch/grid systems to produce histograms.
  - Fit surrogate models  $\hat{b}(p)$  for each histogram bin using the MC outputs – traditionally 3rd/4th-order polynomials, though other forms are possible (e.g. rational approximants via Apprentice).
  - Construct a surrogate goodness-of-fit from these models and optimise the parameters efficiently.
- Expert knowledge is still essential – but surrogates dramatically reduce the cost of scanning high-dimensional parameter spaces.
- Machine learning? Sure – but if polynomials (possibly after a variable transformation) capture the behaviour well, they remain simple, transparent, and robust.



[Andy Buckley]



# Tactics for tuning

## → Factorise the parameter space

- Traditionally split hadron flavours/spectra, jet structure, event topologies, underlying event – typically  $\mathcal{O}(10)$  parameters
- Approximate but practical; reduces dimensionality and stabilises fits.
- Possible to automate grouping via mutual sensitivities / parameter-bin influence matrices.

## → Weighting, observable balance, and uncertainties

- Some data types dominate the fit unless reweighted — balance is essential.
- Models cannot fully describe all bins: examine envelopes, sensitivity maps, parameter ranges, and consider per-bin weighting.
- Custom goodness-of-fit is common, but regularisation weakens strict statistical interpretation.
- Even “ $\chi^2$ ” is non-classical in practice: eigentunes, empirical tolerances... still room for more principled approaches.

## → Future directions

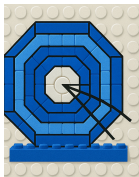
- Improved treatment of heavy flavour; incorporating matching/merging systematics.
- Systematic-uncertainty handling via event weights.
- More flexible surrogate models and robust optimisation strategies.

# Fiducial cross-sections: simple idea, big impact

→ A surprisingly powerful idea

- 👍 Define cross-sections using observable final-state particles, within a clearly specified kinematic region.
- 👍 Makes it easy to reproduce key plots, enabling real understanding, catching issues early, and improving MCs.
- 👍 Establishes a shared language between theory and experiment – essential for tuning, fits, and reinterpretation.

[fiducial]



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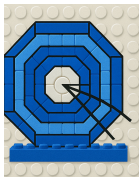
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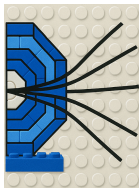
→ But it's tempting to cheat . . .

- 🗨️ Partons, bosons, and other “truth” objects in the event record look easy to use.
- 🗨️ In practice, they're often ambiguous, model-dependent, or even non-existent (e.g. in higher-order simulations).
- ⚠️ Focus on physical final states (hadrons, leptons, photons) as the reliable basis for comparisons.

[fiducial]



[extrapolated]



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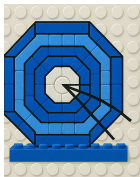
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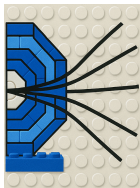
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😊 Who knows – **nature might still surprise us!**

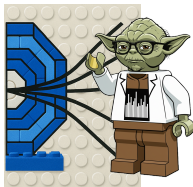
[fiducial]



[extrapolated]



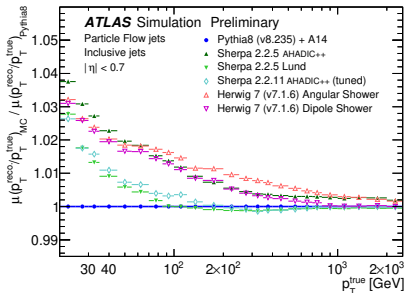
[nature]



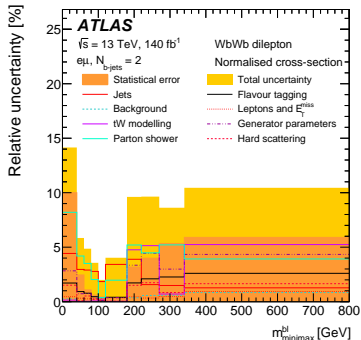


# Generator uncertainties

- ☺ Main systematics: scale variation, PDF choice, matching/merging systematics
- ☹ Alternative generators often used to estimate modelling systematics (2-point systematics)
- ☹ Tuning systematics: UE/hadronisation model dependence (often neglected)

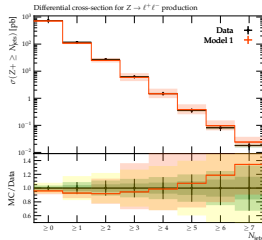
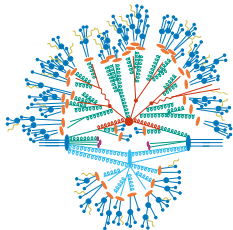


🔗 [ATL-PHYS-PUB-2022-021]



🔗 [arXiv:2506.14700]

# Reweighting techniques



→ Compute alternative event weights during generation for:

- Renormalisation/factorisation scale variations (e.g.  $\mu_r, \mu_f = 0.5, 1, 2 \times$  central)
- PDF error sets (e.g. eigenvectors or replicas from CT18, NNPDF, MSHT etc.)
- NLO electroweak virtual corrections

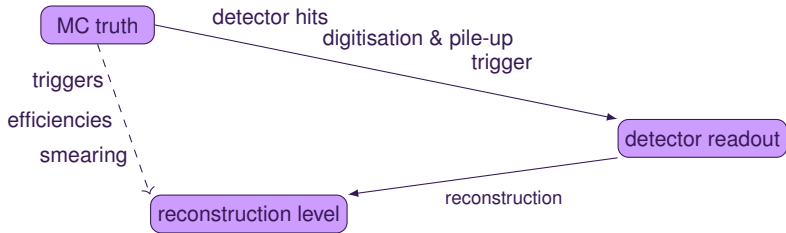
→ Advantages:

- Exact same phase-space sampling for all variations → reduced statistical noise
- No need to rerun generation → huge cost saving

# Monte Carlo production at scale

# Event generators in experiment data processing

- Typical experimental use of generators is to feed their output into a detector simulation (e.g. based on Geant4)
- Then apply the same reconstruction & analysis as for data:

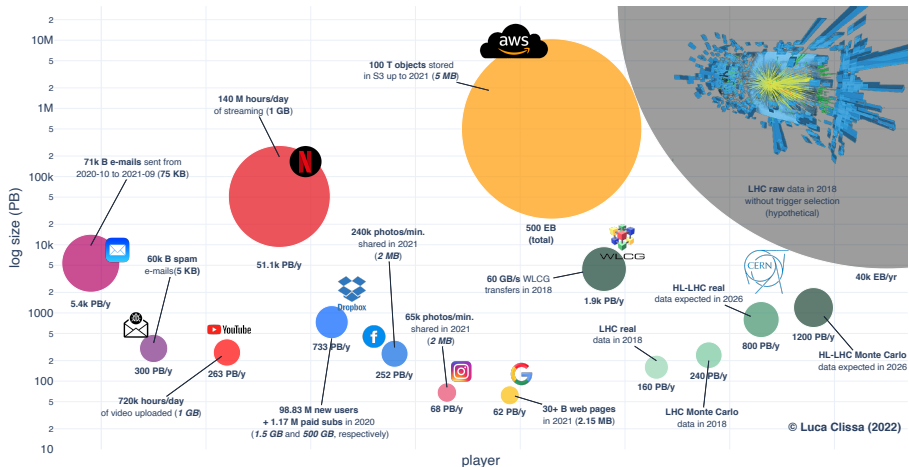


- The generator bit of this chain was historically cheap.
  - CPU/memory requirements much less than detector geometry, *B*-field stepping, material interaction and secondaries.

**⚠ Not true these days!**

# LHC = Big Data

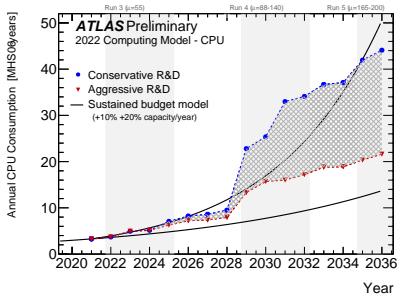
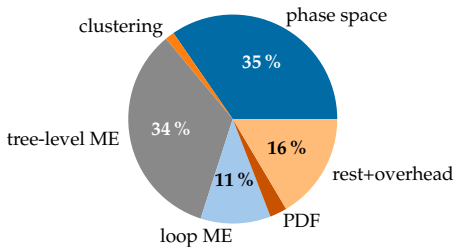
[\[arXiv:2202.07659\]](https://arxiv.org/abs/2202.07659)



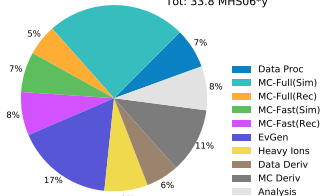
# Expected computing requirements

- Projected evolution of computing resources saw cost of event generation on par with detector simulation
- LHC measurements in danger of being limited by Monte Carlo statistics

$$pp \rightarrow e^+e^- + 0,1,2j @ \text{NLO} + 3,4,5j @ \text{LO}$$



ATLAS Preliminary  
2022 Computing Model - CPU: 2031, Conservative R&D  
Tot: 33.8 MHS06\*y

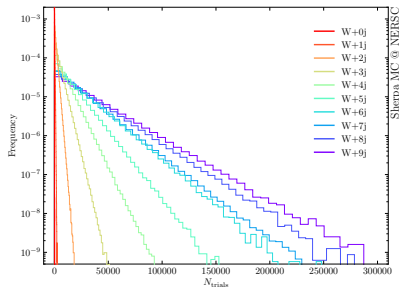
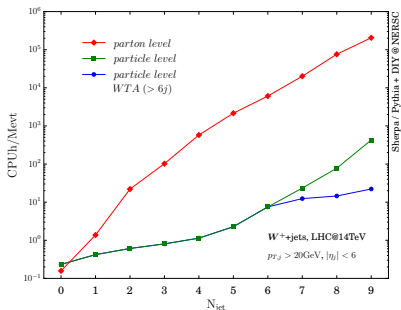


[CERN-LHCC-2022-005]

$$d\sigma_{pp \rightarrow n} = f_a(x_a, \mu_f^2) f_b(x_b, \mu_f^2) \times |\mathcal{M}_{ab \rightarrow n}|^2 \times d\Phi_n(p_1, \dots, p_n)$$

# Understanding the cost of precision Monte Carlo

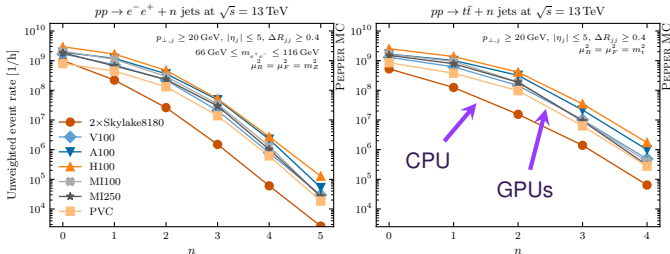
- Cost of multi-leg production **dominated by parton level** matrix elements.
  - Number of diagrams grows factorially with every additional emission (at best exponentially when exploiting recursions a la Comix).
- High-multiplicity matrix elements lend themselves well to GPU-based bulk production.
- Low-multiplicity matrix elements cheaper to regenerate entirely than to store on disk.



🔗 [PRD 100 (2019) 1]

# Architecture-based acceleration

- Computing architectures are shifting rapidly, driven in part by AI-driven hardware trends
  - HPC is entering the exascale era, requiring both scalability and portability across diverse platforms
- 👉 Pepper and MadGraph4GPU are designed with these goals in mind, using:
  - MPI for parallelisation, HDF5 for I/O, and Kokkos for performance portability
- These tools parallelise the entire parton-level event generation chain
  - Supports both multi-core CPUs and a wide range of GPU architectures



🔗 [SciPost Phys. 17 (2024) 3, 081]

# Final thoughts

# Open-source tools & community infrastructure

👥 Monte Carlo (support) tools are community-driven and open-source

→ Transparent, inspectable, and modifiable code

🔧 Widely used and maintained by both theorists and experimentalists:

→ FastJet: Standard library for jet clustering in collider physics

→ LHAPDF: Standardised interface for accessing and interpolating PDF sets

→ HepMC: Common event record format for storing full generator output

→ Rivet: Portable, experiment-independent analysis toolkit for generator validation

→ YODA: Histogramming and statistical accumulator library used by Rivet

→ Contur: Constraint setting from LHC measurements and beyond using Rivet

→ Professor/Apprentice: Generator tuning framework using fast parameterisation

→ HEPData: Archival and retrieval of experimental data for reuse in MC validation

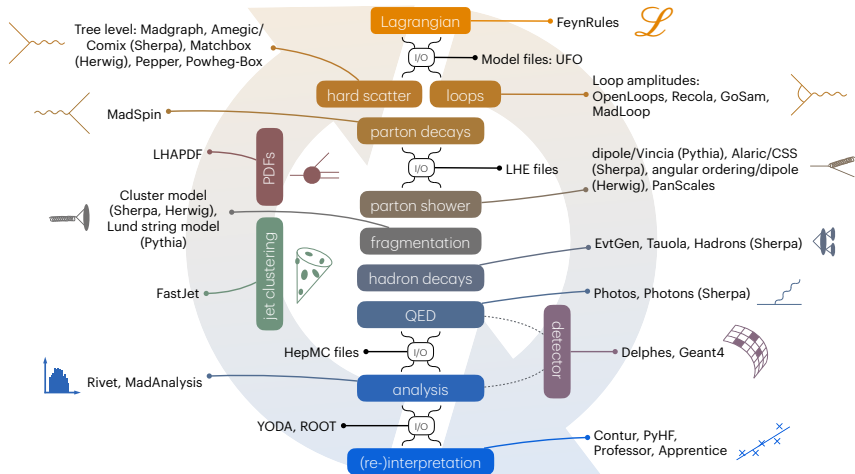
🔄 Shared infrastructure = less duplication, better validation, reproducibility

👤 Open tools = easier training and broader contributions

🗨️ Cross-talk between experiments, theorists, and developers is essential.



# The Monte Carlo Ecosystem in High-Energy Physics



[\[arXiv:2605.16036\]](https://arxiv.org/abs/2605.16036)

# Summary

- Event generation is **principled approximation**: a modular chain turning theory into detector-ready events
  - **Hard process**: matrix elements at LO/NLO, matched and merged with the shower
  - **Parton shower**: resums the dominant soft/collinear radiation
  - **Hadronisation + UE**: non-perturbative models (string, cluster), tuned to data
  - **Decays, detector, analysis**: the path to something comparable with measurement
- Every stage carries **choices and uncertainties** – not a black box!
  - Know which knobs you're turning, and which sources of uncertainty they probe.
- These tools are long-lived, shared infrastructure – you'll use them for years, and you can contribute back!
- Further reading:
  - **MCnet review** (2011) [🔗 \[arxiv:1101.2599\]](#)
  - **Höche review** (2014) [🔗 \[arxiv:1411.4085\]](#)
  - **Practical Collider Physics** (2021) [🔗 \[Buckley, White \(×2\)\]](#)
  - **MC primer** (2026) [🔗 \[arxiv:2605.16036\]](#)

# MCnet short-term studentships

🎓 3-month PhD placements with MC developer teams, in collaboration with the LHC Physics Centre at CERN



👥 What is it?

- A chance to work closely with MC generator developers on a focused project.
- Similar to past MCnet-funded short-term placements — but now open more broadly.
- Not limited to MCnet member nodes — external projects and students welcome.

👤 Who can apply?

- PhD students interested in event generators, theory/phenomenology, or experimental MC use.
- Especially suited for students looking to deepen their expertise in LHC simulation tools.

🔗 More info & available projects [🔗 \[click me\]](#)

- New rounds of placements will be announced 2–3 times per year.
- Feel free to reach out with ideas for new projects – external proposals encouraged!

**Backup**

# QED corrections

- Collinear QED radiation affecting lepton isolation and mass resolution
- Main tools used:
  - Pythia8 / Herwig7 include basic QED FSR in their shower, good enough for many purposes
  - Photos used as afterburner to Herwig/Pythia-based precision samples
  - Sherpa's YFS-style soft photon resummation
- Important for precision Drell-Yan and  $4\ell$  production etc.
- Can shift kinematic distributions subtly – affects fiducial phase-space matching!

