



# National Quantum Computing Centre



Nobel Prize in Physics 2025  
RAL presentation

Dr Jonathan Burnett

*Deputy Director of Research –  
National Quantum Computing Centre*

# Nobel prize in Physics 2025

“For the discovery of macroscopic quantum mechanical tunnelling and energy quantisation in an electric circuit”



John Clarke (born 1942)  
Professor Emeritus at UC Berkeley



Michel H. Devoret (born 1953)  
Professor Emeritus at Yale  
Professor at UC Santa Barbara  
Chief Scientist at Google AI



John M. Martinis (born 1958)  
Professor Emeritus at UC Santa Barbara  
Team-lead Google AI  
Co-Founder & CTO at QoLab

+Many awards and honours between them

# Nobel prize in Physics 2025

“For the discovery of macroscopic quantum mechanical tunnelling and energy quantisation in an electric circuit”

Headline paper(s) –

- Devoret, Martinis & Clarke PRL 55 (1985)
- Martinis, Devoret & Clarke, PRL 55 (1985)

Accompanying papers

- Devoret, Martinis, Esteve & Clarke, PRL 53 (1984)
- Martinis, Devoret & Clarke, PRB 35 (1987)
- Devoret, Esteve, Martinis, Cleland & Clarke, PRB 36 (1987)
- Clarke, Cleland, Devoret, Esteve & Martinis, Science 239 (1988)

VOLUME 55, NUMBER 18

PHYSICAL REVIEW LETTERS

28 OCTOBER 1985

## Measurements of Macroscopic Quantum Tunneling out of the Zero-Voltage State of a Current-Biased Josephson Junction

Michel H. Devoret,<sup>(a)</sup> John M. Martinis, and John Clarke

*Department of Physics, University of California, Berkeley, California 94720, and Materials and Molecular Research Division, Lawrence Berkeley Laboratory, Berkeley, California 94720*  
(Received 26 July 1985)

The escape rate of an underdamped ( $Q \approx 30$ ), current-biased Josephson junction from the zero-voltage state has been measured. The relevant parameters of the junction were determined *in situ* in the thermal regime from the dependence of the escape rate on bias current and from resonant activation in the presence of microwaves. At low temperatures, the escape rate became independent of temperature with a value that, with no adjustable parameters, was in excellent agreement with the zero-temperature prediction for macroscopic quantum tunneling.

PACS numbers: 74.50.+r, 03.65.-w, 05.30.-d, 05.40.+j

The observation of macroscopic quantum tunneling is regarded as a test of whether quantum mechanics is valid for macroscopic variables, a fundamental question<sup>1</sup> that has only recently been addressed experimentally. The necessary conditions for the observation of macroscopic quantum tunneling can be realized in the current-biased Josephson tunnel junction, where the phase difference between the two superconductors is the macroscopic variable, and the tunneling occurs from the zero-voltage state to the nonzero-voltage state. Previous experiments on a current-biased Josephson junction<sup>2-4</sup> or on a superconducting ring interrupted by a Josephson junction<sup>5-7</sup> have yielded results that have been interpreted as being consistent with the theoretical predictions for macroscopic quantum tunneling. In this Letter, we present results of

In the zero-voltage state, the plasma frequency  $\omega_p/2\pi$  of small oscillations of the particle at the bottom of the well is  $\omega_p = (2\pi I_0/C\Phi_0)^{1/2}[(1 - (I/I_0)^2)^{1/4}]$ , while the damping factor is  $Q = \omega_p RC$ .

In the thermal regime ( $k_B T \gg \hbar\omega_p$ ), the escape of the particle from the well occurs via thermal activation at a rate<sup>10</sup>

$$\Gamma_t = a_t(\omega_p/2\pi)\exp(-\Delta U/k_B T), \quad (1)$$

where  $a_t = 4/[(1 + Qk_B T/1.8\Delta U)^{1/2} + 1]^2$  is of the order of unity in our experiment,  $k_B$  is Boltzmann's constant, and  $T$  is the temperature. In the quantum regime ( $k_B T \ll \hbar\omega_p$ ), to lowest order in  $1/Q$  the escape is predicted to occur via macroscopic quantum tunneling at a rate<sup>11</sup>



# Outline

## Background physics

- Quantum tunnelling
- When might we expect quantum behaviour
- Superconductivity
- Josephson effect

## Walk through the experiment

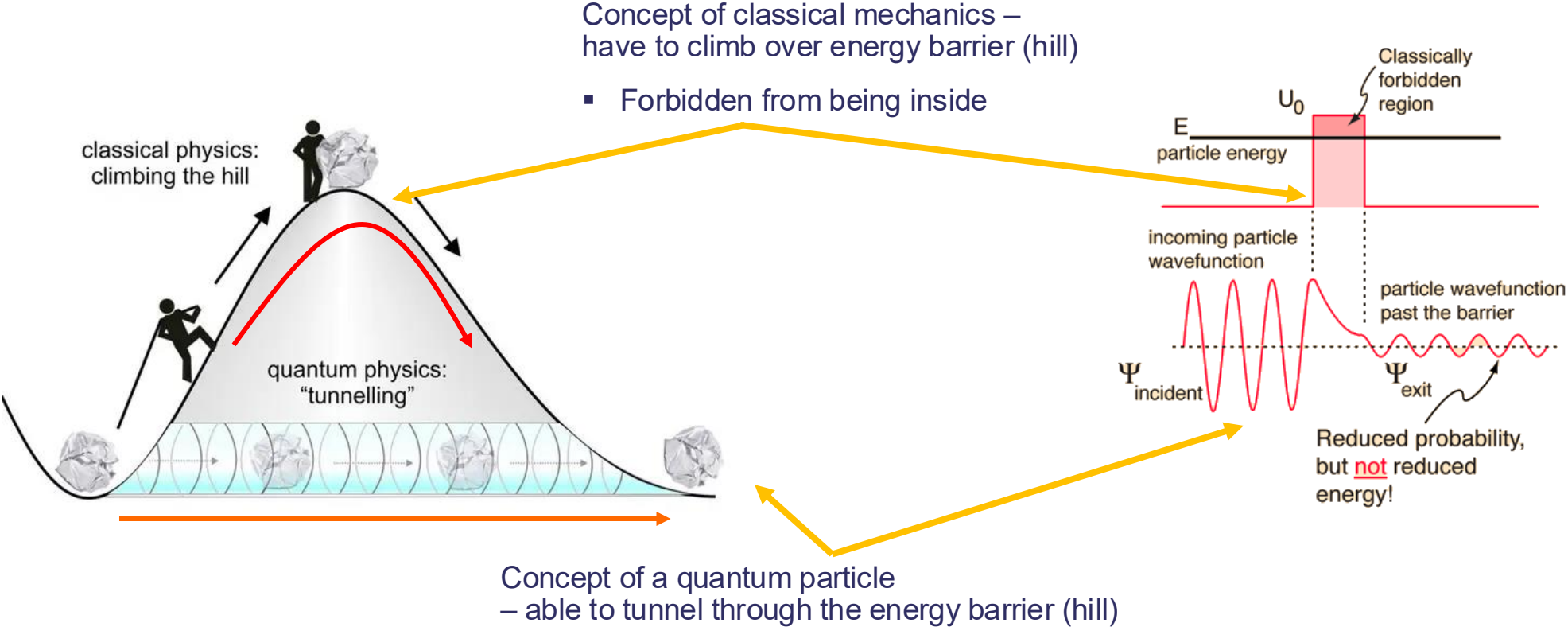
- Experimental setup
- Results – tunnelling
- Results – quantisation

## So what followed

## Quantum Computing in the UK



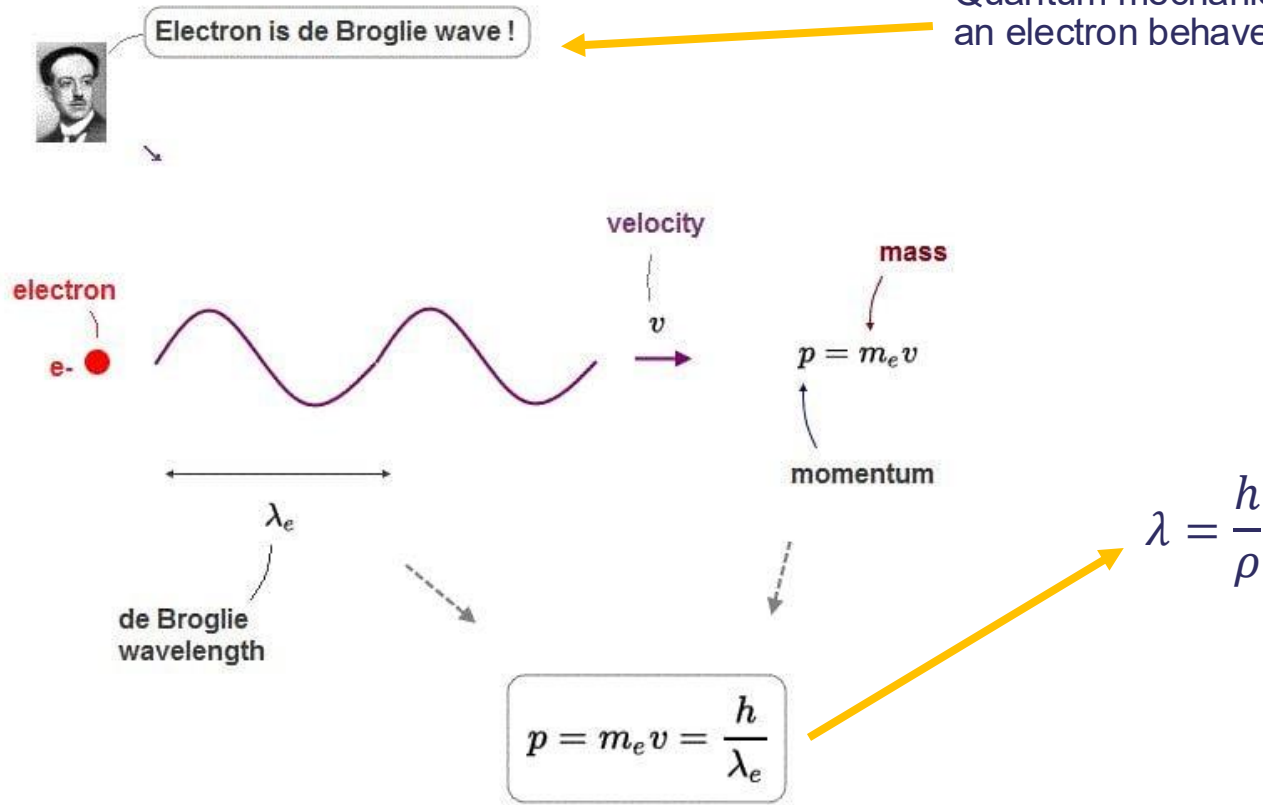
# Quantum tunnelling



- Probability of quantum particle existing on far side (tunnelled)

# When does something behave in a quantum way?

Quantum mechanical property of wave-particle duality, an electron behaves both as a **particle** and as a **wave**



Lets compare an electron to a cannon-ball  
Electron (mass  $9 \times 10^{-31}$  kg, velocity  $2 \times 10^6$  ms<sup>-1</sup>)

Wavelength  $3 \times 10^{-10}$  m -> An atomic length-scale  
-> **quantum behaviour expected**

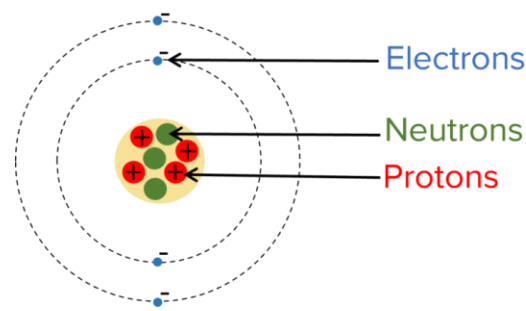
Cannon-ball (mass 10 kg, velocity  $3 \times 10^2$  ms<sup>-1</sup>)

Wavelength  $2 \times 10^{-36}$  m -> unphysically small  
-> **no quantum behaviour expected**



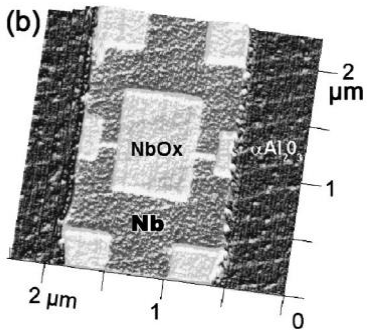
# Micro- and Macro- systems

Microscopic system e.g. an atom (or electron, or photon)



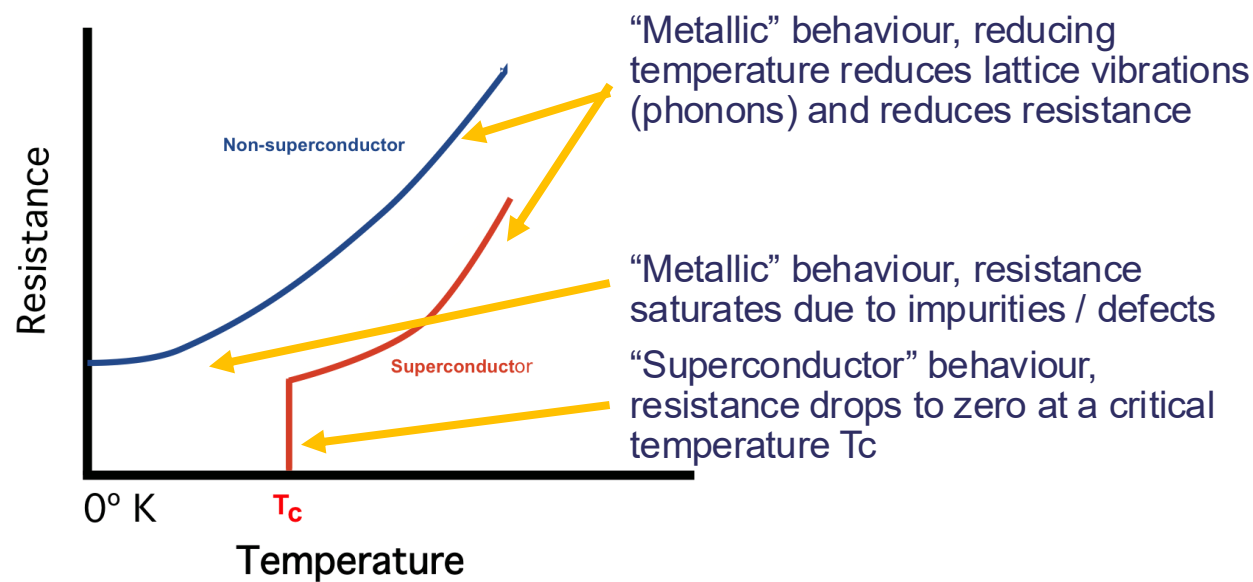
Can quantify properties  
- Mass, charge, nuclear spin etc.

Macroscopic system e.g. Niobium circuit



Hard to quantify “size”  
- How many atoms?  
- What about imperfections?  
Order-of magnitude estimates of properties

# Superconductivity



Superconductivity discovered in 1911 – K Onnes  
Theory descriptions – phenomenological Ginzberg-Landau 1950s  
Microscopic model Bardeen, Cooper and Schrieffer 1957





# Superconductivity

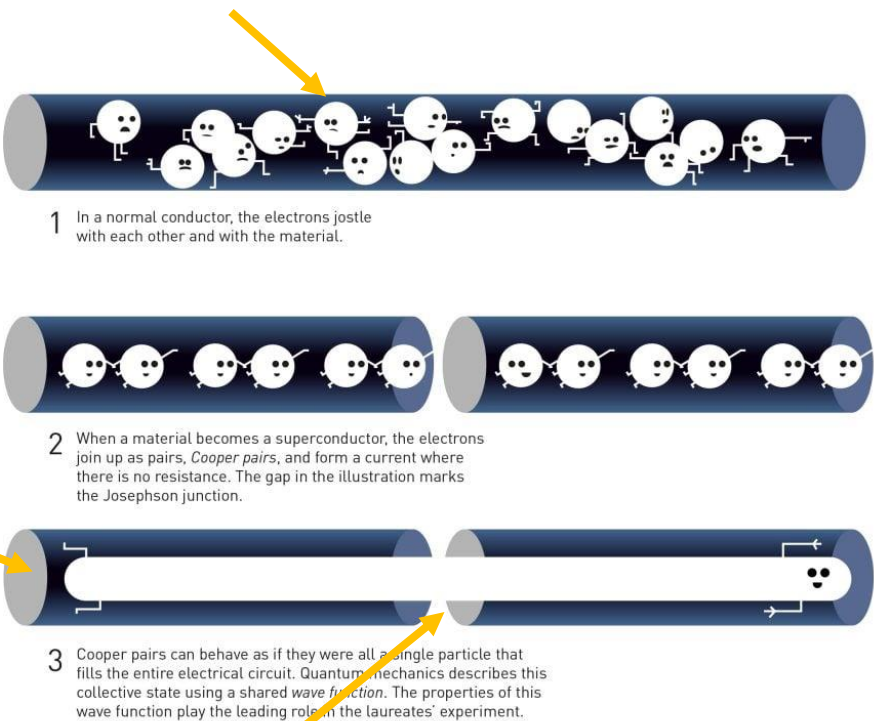
Electrons combine into Cooper Pairs which do not collide -> zero resistance

Cooper Pairs described by a collective quantum state (wavefunction phase  $\varphi_1$ )

## Theory motivation for experiment

It is not clear that this wavefunction (collection from many many electrons) is a quantum mechanical variable! - A.J. Leggett (1983)

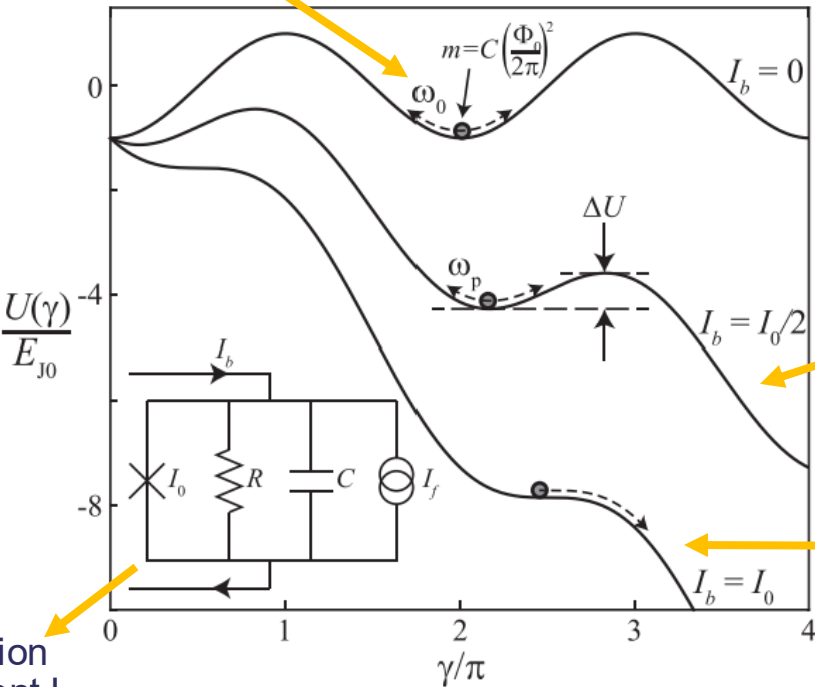
Electron collisions is what gives resistance



“Gap” represents a tunnel barrier called a Josephson junction, the wavefunction can tunnel through.

# Josephson junction circuit

Josephson junction circuit  
acts as a *particle in a well*



When Bias current  $I_b$  is 0,  
potential is a sinusoid

When Bias current  $0 < I_b < I_0$ ,  
potential is a tilted sinusoid

When Bias current  $I_b = I_0$ ,  
potential is effectively a slide

Height of the potential  
 $\Delta U$  scales with  $I_b$

Oscillation frequency  $\omega_p$   
also scales with  $I_b$

Josephson-junction  
biased with current  $I_b$

# Josephson junction circuit

Look at the transport characteristics (I-V)

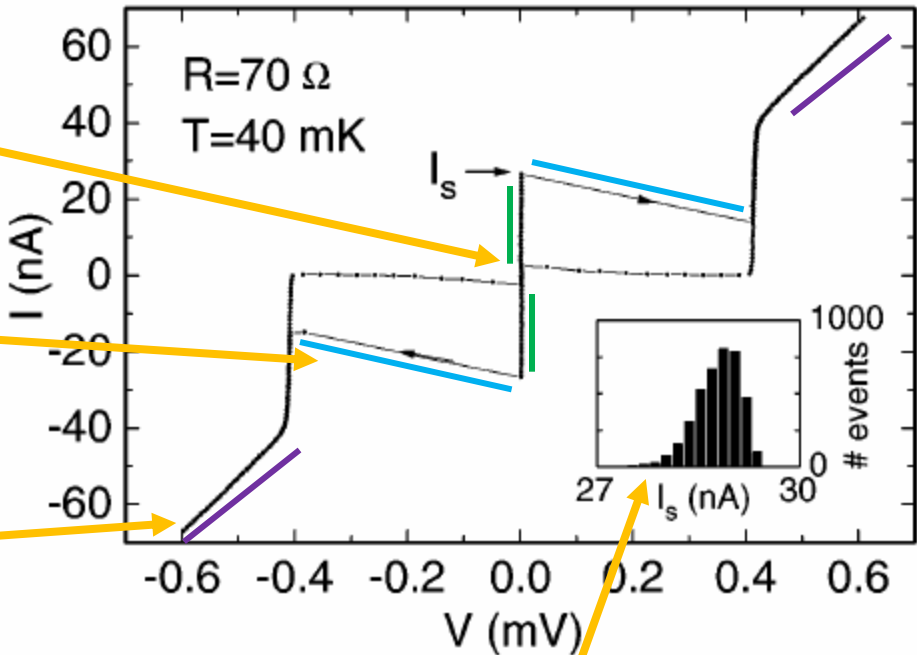
Start at zero current & ramp current

Observe a superconducting current (zero voltage state) is sustained up to a characteristic current  $I_s$

At the switching  $I_s$  current the junction switches to a voltage state where it has resistance

Junction acts a threshold detector

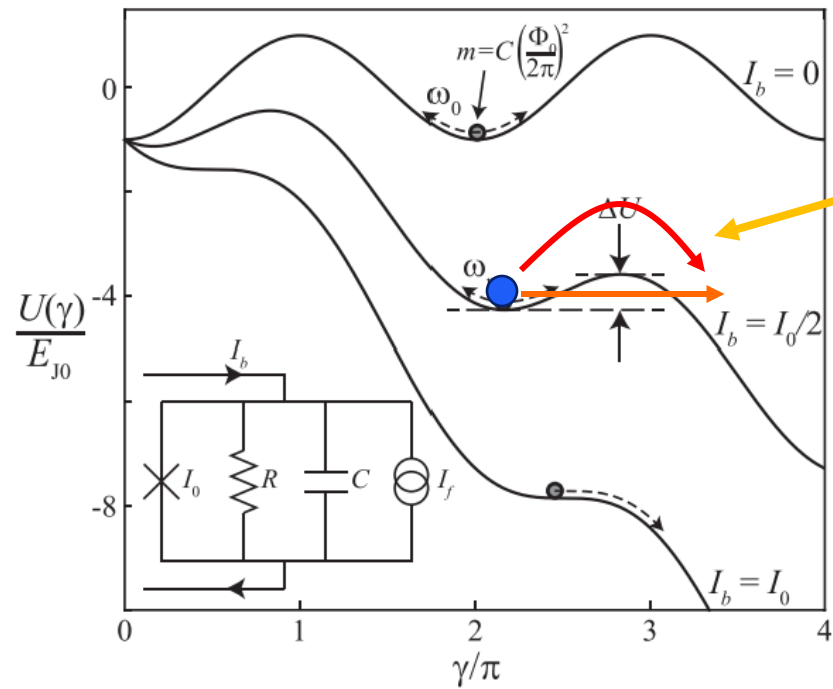
When large bias current is applied, the Josephson junction behaves like a normal conductor e.g.  $V=IR$



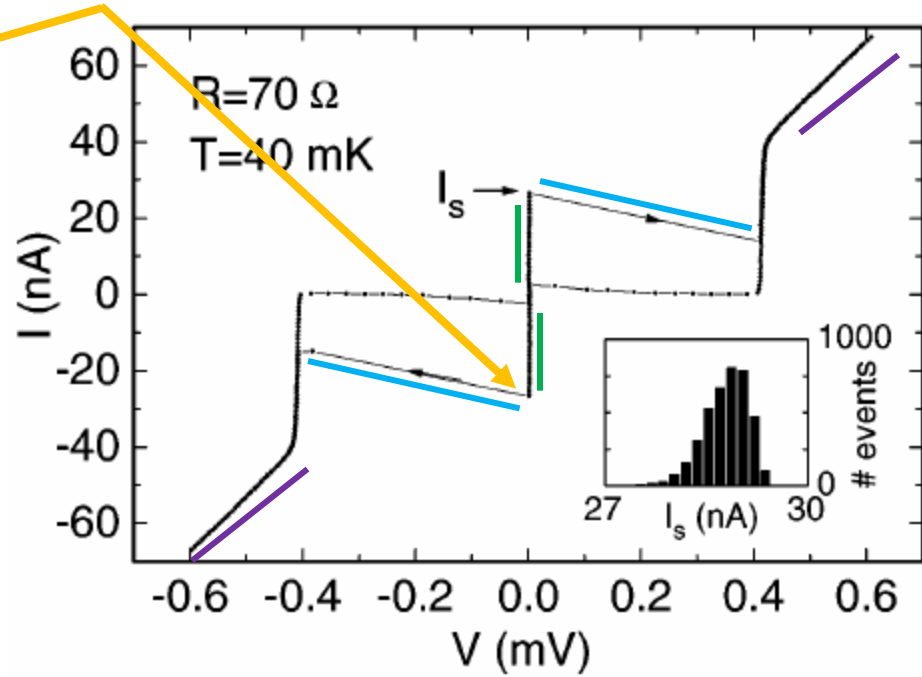
Through repeat and reset cycles, can look at the distribution of the switching current  $I_s$

# Josephson junction circuit

Junction acts a threshold detector with a switching event happening at  $I_s$  where it switches to a voltage state



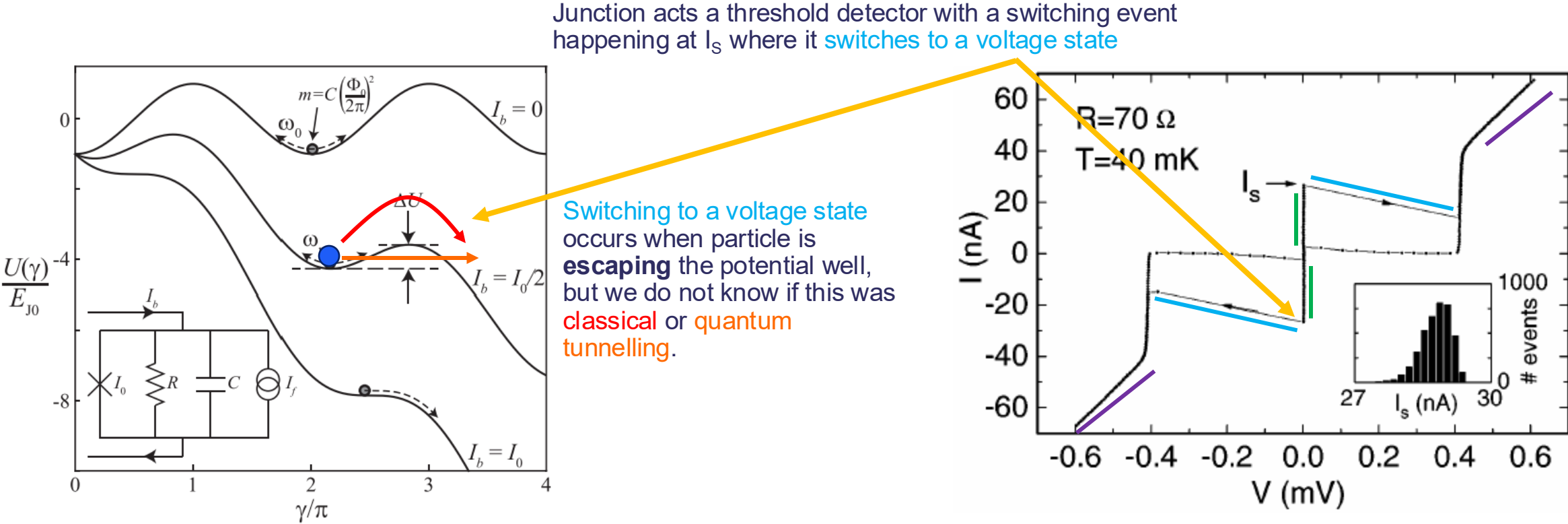
Switching to a voltage state occurs when particle is escaping the potential well, but we do not know if this was classical or quantum tunnelling.



# Experimental design

**Experiment 1**

Create a system with a current ramp, use the voltage state as a threshold detector, **vary temperature** and look for **non-classical behaviour** (i.e. quantum tunnelling)

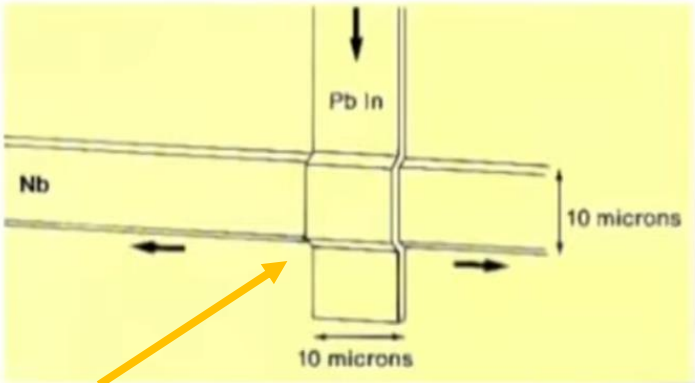


# Experimental design

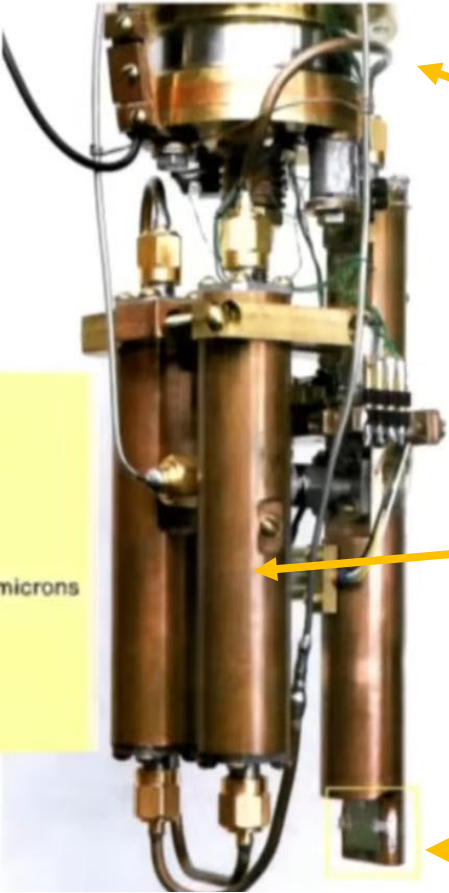
## Experiment 1

Create a system with a current ramp, use the voltage state as a threshold detector, **vary temperature** and look for **non-classical behaviour** (i.e. quantum tunnelling)

EXPERIMENTAL SETUP  
OF THE 1985 BERKELEY  
EXPERIMENT



Nb-NbOx-PbSn Josephson junction  
(Superconductor-Insulator-Superconductor)



Dilution fridge to reach 20mK  
(-273.13 Celsius)

Cu-Filters  
(explained later)

Sample mount  
(Josephson device sits here)

# Experimental design

## Experiment 1

Create a system with a current ramp, use the voltage state as a threshold detector, **vary temperature** and look for **non-classical behaviour** (i.e. quantum tunnelling)

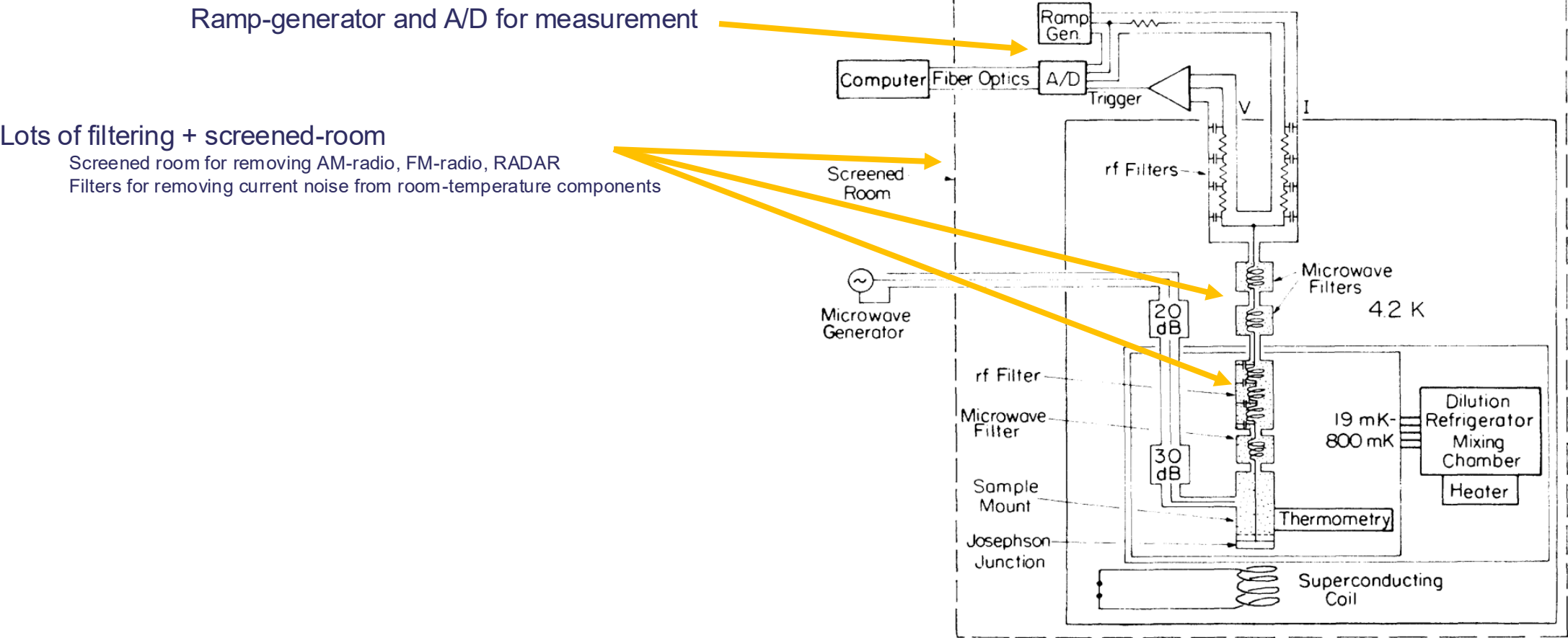


FIG. 3. Schematic drawing of apparatus.



# Experimental design

- Sample mount for the Josephson junction device
- Dilution fridge for cooling to 20mK
- Ramp-generator and A/D for measurement
- Lots of filtering + screened-room
  - Screened room for removing AM-radio, FM-radio, RADAR
  - Filters for removing current noise from room-temperature components

$$I_{RMS} = \sqrt{\frac{4 k_B T BW}{R}}$$

$$I_b = I_{RMS} + I_{RampGen}$$

Without filtering,  
→ 2μA !

With filtering  
→ 0.0001μA = 0.1 nA!

## Experiment 1

Create a system with a current ramp, use the voltage state as a threshold detector, **vary temperature** and look for **non-classical behaviour (i.e. quantum tunnelling)**

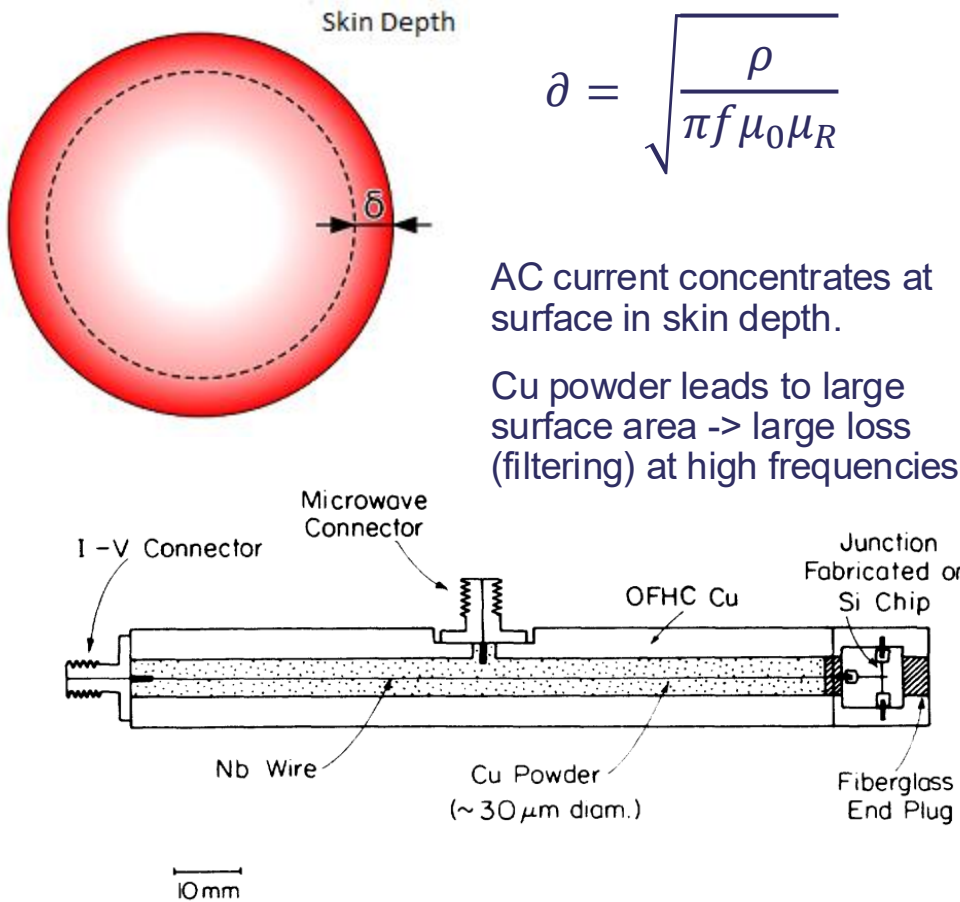


FIG. 4. Scale drawing of filter on which junction was mounted.

# Experimental design

Measuring the escape events by measuring switching current  
switching events while varying temperature (x-axis)

High temperature, find switching events proportional to temperature

- Switching is classical (thermally assisted)

Low temperature, find switching events saturate at non-zero

- Evidence for quantum tunnelling

**Had to demonstrate the saturation was not due to electrical noise!**

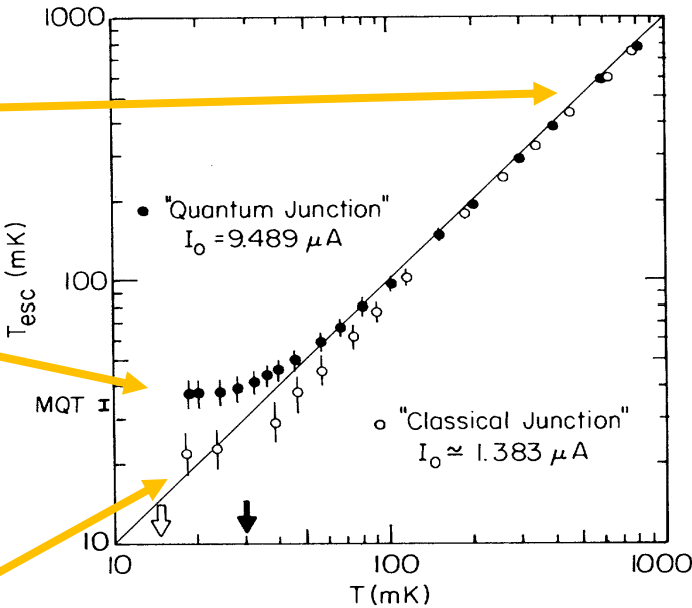
Applying magnetic field suppresses superconductivity -> Suppress the critical current

Suppress the transition to quantum behaviour to lower threshold and no longer observe the evidence for quantum tunnelling

- Self calibration that the plateau was not due to noise!

## Experiment 1

Create a system with a current ramp, use the voltage state as a threshold detector, **vary temperature** and look for **non-classical behaviour (i.e. quantum tunnelling)**



**Evidence for Macroscopic quantum tunnelling!**

# Experimental design

## Experiment 2

With quantum behaviour shown, now drive with RF and look for resonant enhancements which would be evidence for **quantised energy levels**

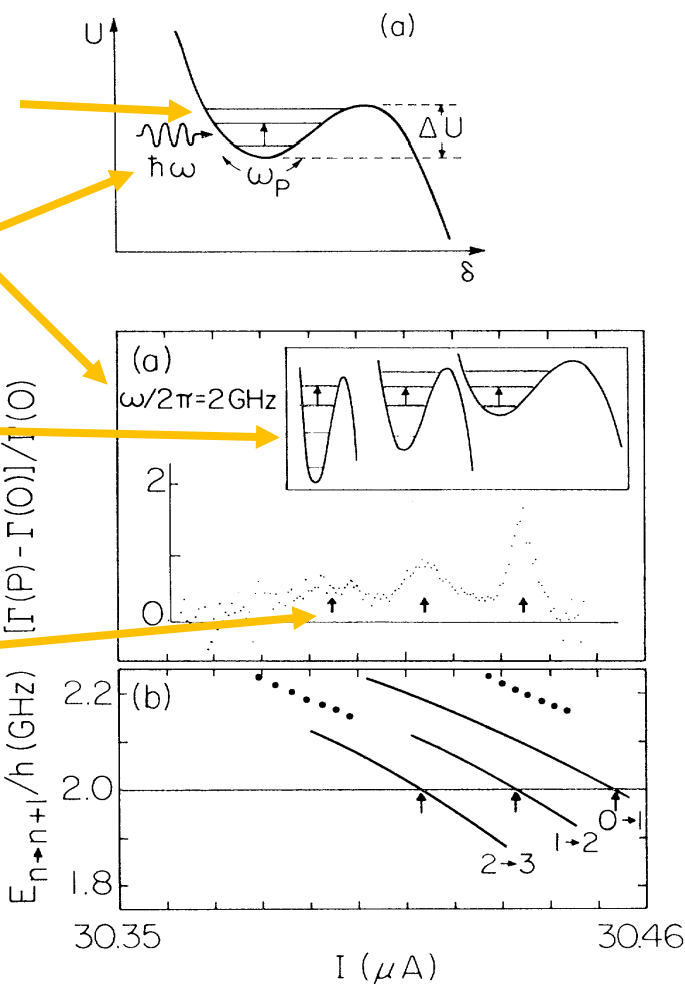
If the phase is a quantum system, would expect quantised energy levels.

Apply microwave radiation at 2 GHz

When sweeping  $I_b$  we change the shape of the potential well, -> which changes energy levels

When energy levels swept through the 2 GHz radiation, we then see peaks in the escape rate -> resonant enhancement

**Evidence for quantised energy levels!**



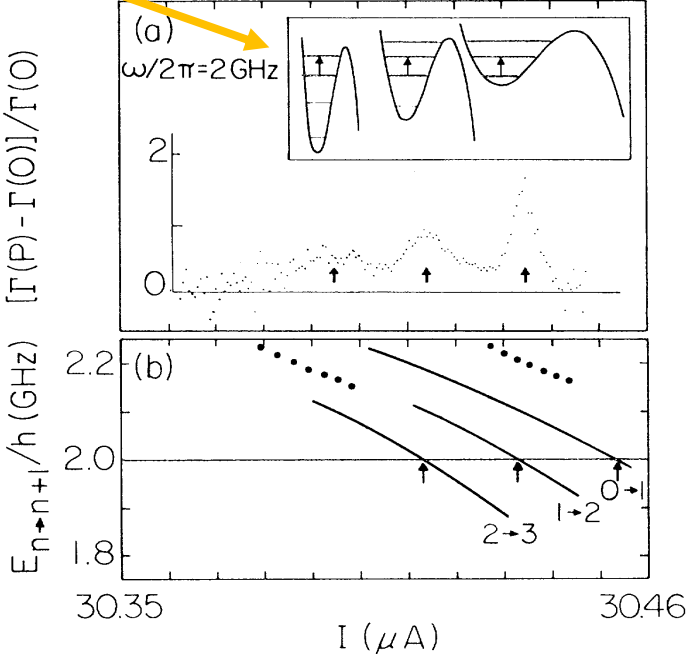
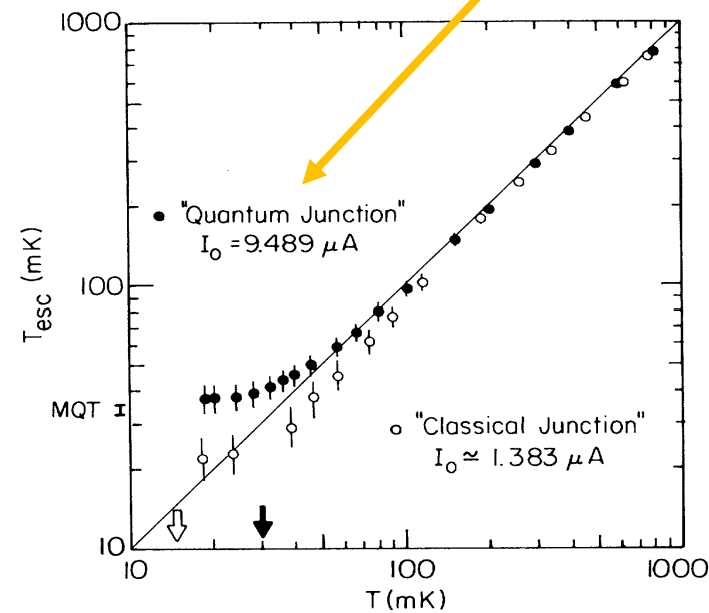
# Nobel prize in Physics 2025

“For the discovery of macroscopic quantum mechanical tunnelling and energy quantisation in an electric circuit”

Discovery of macroscopic quantum tunnelling in an electronic circuit

Discovery of energy quantisation in an electric circuit

An electric circuit acting as an **artificial atom**



An electric circuit acting as an *artificial atom*  
... so quantum computing?

# So what followed: I

VOLUME 58, NUMBER 9

PHYSICAL REVIEW LETTERS

2 MARCH 1987

Superconductivity at 93 K in a New Mixed-Phase Y-Ba-Cu-O Compound System at Ambient Pressure

M. K. Wu, J. R. Ashburn, and C. J. Torng  
*Department of Physics, University of Alabama, Huntsville, Alabama 35899*

and

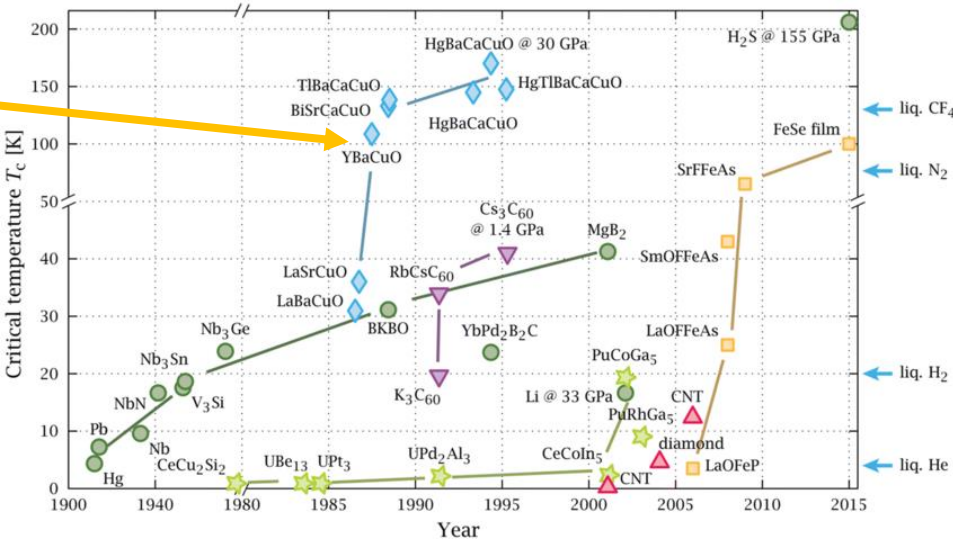
P. H. Hor, R. L. Meng, L. Gao, Z. J. Huang, Y. Q. Wang, and C. W. Chu<sup>(a)</sup>  
*Department of Physics and Space Vacuum Epitaxy Center, University of Houston, Houston, Texas 77004*  
(Received 6 February 1987; Revised manuscript received 18 February 1987)

A stable and reproducible superconductivity transition between 80 and 93 K has been unambiguously observed both resistively and magnetically in a new Y-Ba-Cu-O compound system at ambient pressure. An estimated upper critical field  $H_{c2}(0)$  between 80 and 180 T was obtained.

PACS numbers: 74.70.Ya

The search for high-temperature superconductivity and novel superconducting mechanisms is one of the most challenging tasks of condensed-matter physicists and material scientists. To obtain a superconducting state reaching beyond the technological and psychological

sharper transition. A transition width<sup>10</sup> of 2 K and an onset<sup>11</sup>  $T_c$  of 48.6 K were obtained at ambient pressure. Pressure<sup>8,12</sup> was found to enhance the  $T_c$  of the La-Ba-Cu-O system at a rate of greater than  $10^{-3}$  K bar<sup>-1</sup> and to raise the onset  $T_c$  to 57 K, with a “zero-resis-



Other results relating to high- $T_c$  superconductors dominated funding and research in condensed matter



# So what followed: II

VOLUME 75, NUMBER 25      PHYSICAL REVIEW LETTERS      18 DECEMBER 1995

**Demonstration of a Fundamental Quantum Logic Gate**

C. Monroe, D. M. Meekhof, B. E. King, W. M. Itano, and D. J. Wineland  
*National Institute of Standards and Technology, Boulder, Colorado 80303*  
(Received 14 July 1995)

We demonstrate the operation of a two-bit "controlled-NOT" quantum logic gate, which, in conjunction with simple single-bit operations, forms a universal quantum logic gate for quantum computation. The two quantum bits are stored in the internal and external degrees of freedom of a single trapped atom, which is first laser cooled to the zero-point energy. Decoherence effects are identified for the operation, and the possibility of extending the system to more qubits appears promising.

PACS numbers: 89.80.+h, 03.65.-w, 32.80.Fj

We report the first demonstration of a fundamental quantum logic gate that operates on prepared quantum states. Following the scheme proposed by Cirac and computers and CN gates involving a dipole-dipole interaction between quantum dots or atomic nuclei [6,7,11,12] may suffer from decoherence efforts. The light shifts on

PRL **100**, 113003 (2008)      PHYSICAL REVIEW LETTERS      week ending 21 MARCH 2008

**Rabi Oscillations between Ground and Rydberg States with Dipole-Dipole Atomic Interactions**

T. A. Johnson, E. Urban, T. Henage, L. Isenhower, D. D. Yavuz, T. G. Walker, and M. Saffman  
*Department of Physics, University of Wisconsin, 1150 University Avenue, Madison, Wisconsin 53706, USA*  
(Received 2 November 2007; published 19 March 2008)

We demonstrate Rabi oscillations of small numbers of <sup>87</sup>Rb atoms between ground and Rydberg states with  $n \leq 43$ . Coherent population oscillations are observed for single atoms, while the presence of two or more atoms decoheres the oscillations. We show that these observations are consistent with van der Waals interactions of Rydberg atoms.

DOI: 10.1103/PhysRevLett.100.113003      PACS numbers: 32.80.Rm, 03.67.-a

Atoms in highly excited Rydberg states with principal quantum number  $n \gg 1$  have very large transition dipole moments which scale as  $d \propto na^2$  with  $a$  the electron-magnetic interaction length. The experiment starts by loading a far-off-resonance optical trap (FORT) from a <sup>87</sup>Rb vapor cell magnetooptical trap (MOT) as described in our recent Letter [7].

PHYSICAL REVIEW A **68**, 032316 (2003)

**Experimental controlled-NOT logic gate for single photons in the coincidence basis**

T. B. Pittman, M. J. Fitch, B. C. Jacobs, and J. D. Franson  
*Applied Physics Laboratory, Johns Hopkins University, Laurel, Maryland 20723, USA*  
(Received 2 April 2003; published 26 September 2003)

We report a proof-of-principle demonstration of a probabilistic controlled-NOT gate for single photons. Single-photon control and target qubits were mixed with a single ancilla photon in a device constructed using only linear optical elements. The successful operation of the controlled-NOT gate relied on post-selected three-photon interference effects, which required the detection of the photons in the output modes.

DOI: 10.1103/PhysRevA.68.032316      PACS number(s): 03.67.Lx, 42.50.Dv, 42.65.Lm

There has been considerable interest in a linear optics approach to quantum computing [1,2], in which probabilistic two-qubit logic operations are implemented using linear optical elements and single-photon detectors. The first demonstration of a controlled-NOT gate for single photons was shown in Fig. 1(a). Here the logical value of each of the qubits is represented by the polarization state of a single photon, where a horizontal polarization state  $|H\rangle$  represents a value of 0 and a vertical polarization state  $|V\rangle$  represents a value of 1.

1995 – Trapped-ion qubit demonstrated

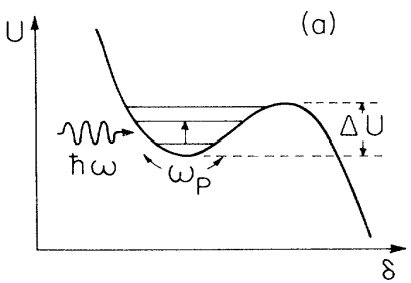
2003 – Photonics qubit demonstrated

2008 – Neutral atoms qubit demonstrated

“Natural” qubits



National Quantum Computing Centre



An electric circuit acting as an *artificial atom*

... so quantum computing?

Qubit is building block of quantum computer

Isolate and manipulate the occupation of two energy levels  
→ Qubit

letters to nature

**Coherent control of macroscopic quantum states in a single-Cooper-pair box**

Y. Nakamura\*, Yu. A. Pashkin† & J. S. Tsai\*

\* NEC Fundamental Research Laboratories, Tsukuba, Ibaraki 305-8051, Japan  
† CREST, Japan Science and Technology Corporation (JST), Kawaguchi, Saitama 332-0012, Japan

A nanometre-scale superconducting electrode connected to a reservoir via a Josephson junction constitutes an artificial two-level electronic system: a single-Cooper-pair box. The two levels consist of charge states (differing by  $2e$ , where  $e$  is the electronic charge) that are coupled by tunnelling of Cooper pairs through the junction. Although the two-level system is macroscopic,

**RESEARCH ARTICLES**

**Coherent Manipulation of Coupled Electron Spins in Semiconductor Quantum Dots**

J. R. Petta,<sup>1</sup> A. C. Johnson,<sup>1</sup> J. M. Taylor,<sup>1</sup> E. A. Laird,<sup>1</sup> A. Yacoby,<sup>2</sup> M. D. Lukin,<sup>1</sup> C. M. Marcus,<sup>1</sup> M. P. Hanson,<sup>3</sup> A. C. Gossard<sup>3</sup>

We demonstrated coherent control of a quantum two-level system based on two-electron spin states in a double quantum dot, allowing state preparation, coherent manipulation, and projective readout. These techniques are based on rapid electrical control of the exchange interaction. Separating and later recombining a singlet spin state provided a measurement of the spin dephasing time,  $T_2^*$ , of  $\sim 10$  nanoseconds, limited by hyperfine interactions with the gallium arsenide host nuclei. Rabi oscillations of two-electron spin states were demonstrated, and spin-echo pulse sequences were used to suppress hyperfine-induced dephasing. Using these quantum control techniques, a coherence time for two-electron spin states exceeding 1 microsecond was observed.

“Engineered” qubits

2025 Nobel prize recognizes that an electrical circuit can behave as an artificial atom

→ Foundational to “engineered” qubits being possible



# UK-NQTP

## (Quantum Mission 1)

### Key components

- 10-year goal with key milestone at 2035
- UK-based system with compute capabilities that should enable application

### Not specified

- Modality – Hardware agnostic
- Fidelities – agnostic to specifics of error-mitigation, specifics of error-correction
- Qubit count – agnostic to specifics of error-correction



Policy paper

### National Quantum Strategy Missions

Updated 22 November 2023

Quantum computing specifically part of IS-8 in UK Industrial Strategy 2025

By **2035**, achieving quantum advantage at scale through reaching  **$10^{12}$  (a trillion) quantum operations**, enabling **applications** such as optimising the production of clean hydrogen.

By 2032, demonstrating large-scale error correction capabilities with  **$10^9$  (a billion) quantum operations**, with including accelerated drug discovery. applications


By 2028, extending beyond the NISQ-era with  **$10^6$  (million) quantum operations**, which will enable the exploration of applications associated with the simulation of chemical processes, helping to improve catalyst design for example.



National Quantum  
Computing Centre



# UKRI Quantum Computing



National Quantum Computing Centre

About us Technology & research Infrastructure Engage with us Learn with us Careers News & events Resources

## Integrating Cat-Qubit Quantum Computers into Standard HPC Workflow and Job Scheduling

New access to incorporate cat-qubit technology

NQCC focus – enabling government to act as an informed customer / funder of quantum computing

R&D on quantum computing

- Neutral Atoms, Superconducting, Trapped Ions
- Software, HPC
- Application & Cloud

Quantum readiness

- Training, skills, user-onboarding
- Ethical use of quantum

To celebrate the milestone, the National Quantum Computing Centre will mark this year through a range of activities, including events, interviews and articles. The aim is to stay connected with the quantum community, inspire people from all backgrounds to engage with quantum science and technology, and to advance quantum computing for societal benefit.

## Integrating Cat-Qubit Quantum Computers into Standard HPC Workflow and Job Scheduling

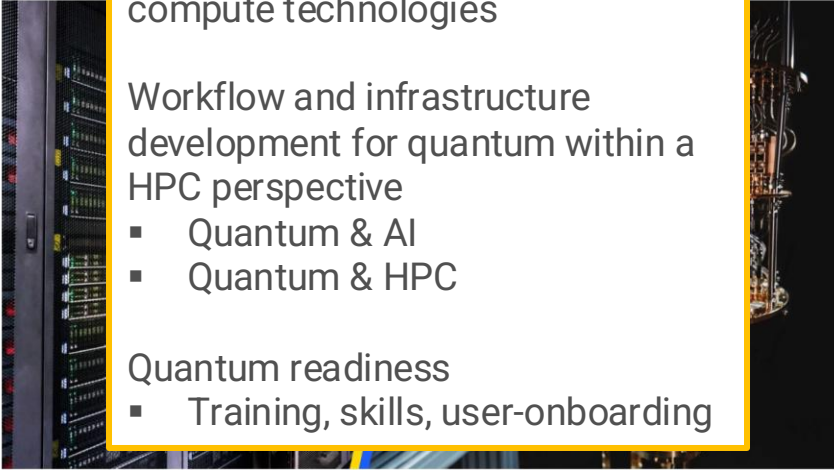
Hartree focus – Business transformation & integration of new compute technologies

Workflow and infrastructure development for quantum within a HPC perspective

- Quantum & AI
- Quantum & HPC

Quantum readiness

- Training, skills, user-onboarding



[Alice & Bob](#), a global leader in the race for fault-tolerant quantum computing, today announced the software integration of their future QPUs with high performance computing (HPC) environments via SLURM, the world's most widely used workload management system, present in 60% of the world's top supercomputers.

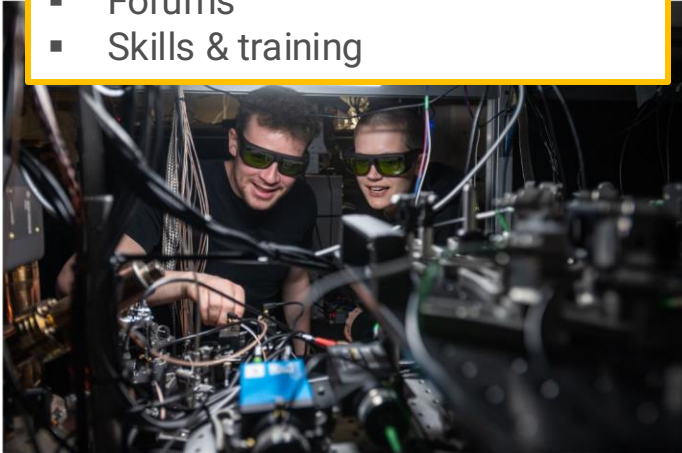
Developed in collaboration with the Science and Technology Facilities Council's (STFC) Hartree Centre through the [Hartree National Centre for Digital Innovation \(HNC DI\)](#), this integration makes Alice & Bob's technology compatible with existing HPC workflows, marking the first time cat-qubit technology is adapted for seamless scheduling and execution within standard HPC resource management frameworks.

## QCi3: Hub for Quantum Computing via Integrated and Interconnected Implementations

Led by: Professor Dominic O'Brien, University of Oxford

QCi3 focus – Academic research hub – network of universities – underpinning science and applied research

- Working groups
- Forums
- Skills & training



Trapped Ions laboratory at the University of Oxford

# Summary

When does something behave quantum -> question of whether an electrical circuit (something macroscopic) could behave as a quantum object

Evidence found in electrical circuit demonstrating non-classical escape events i.e. **macroscopic quantum tunnelling**

Evidence found in electrical circuit demonstrating resonant enhancement of escape events i.e. **quantised energy levels**

Collectively – this means the electrical circuit acts an **artificial atom**

\* Not all Quantum Computers based on macroscopic devices  
– see instead AMO family (photronics, neutral atoms, Trapped Ions)

Quantum computing is part of IS-8 UK Industrial Strategy of 2025

- National Quantum Technologies Program: Quantum Mission 1 =  $10^{12}$  quantum operations by 2035 + quantum as a part of the digital sector
- UKRI activities in quantum computing -> NQCC, Hartree Centre, QCi3 hub (+ many more)

\* Not all Quantum Technologies are  
Quantum Computing!

# Nobel prize in Physics 2025

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Professor Emeritus at UC Berkeley

Michel H. Devoret (born 1953)  
Professor Emeritus at Yale  
Professor at UC Santa Barbara  
Chief Scientist at Google AI

John M. Martinis (born 1958)  
Professor Emeritus at UC Santa Barbara  
Team-lead Google AI  
Co-Founder & CTO at QoLab

Thanks for listening!

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Measurements of Macroscopic Quantum Tunneling out of the Zero-Voltage State of a Current-Biased Josephson Junction

Michel H. Devoret,<sup>(a)</sup> John M. Martinis, and John Clarke  
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(Received 26 July 1985)

The escape rate of an underdamped ( $Q \approx 30$ ), current-biased Josephson junction from the zero-voltage state has been measured. The relevant parameters of the junction were determined *in situ* in the thermal regime from the dependence of the escape rate on bias current and from resonant activation in the presence of microwaves. At low temperatures, the escape rate became independent of temperature with a value that, with no adjustable parameters, was in excellent agreement with the zero-temperature prediction for macroscopic quantum tunneling.

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The observation of macroscopic quantum tunneling is regarded as a test of whether quantum mechanics is valid for macroscopic variables, a fundamental question<sup>1</sup> that has only recently been addressed experimentally. The necessary conditions for the observation of macroscopic quantum tunneling can be realized in the current-biased Josephson tunnel junction, where the phase difference between the two superconductors is the macroscopic variable, and the tunneling occurs from the zero-voltage state to the nonzero-voltage state. Previous experiments on a current-biased Josephson junction<sup>2-4</sup> or on a superconducting ring interrupted by a Josephson junction<sup>5-7</sup> have yielded results that have been interpreted as being consistent with the theoretical predictions for macroscopic quantum tunneling. In this Letter, we present results of

In the zero-voltage state, the plasma frequency  $\omega_p/2\pi$  of small oscillations of the particle at the bottom of the well is  $\omega_p = (2\pi I_0/C\Phi_0)^{1/2}[(1 - (I/I_0)^2)^{1/4}]^{1/4}$ , while the damping factor is  $Q = \omega_p RC$ .

In the thermal regime ( $k_B T \gg \hbar\omega_p$ ), the escape of the particle from the well occurs via thermal activation at a rate<sup>10</sup>

$$\Gamma_t = a_t (\omega_p/2\pi) \exp(-\Delta U/k_B T), \tag{1}$$

where  $a_t = 4/[1 + Qk_B T/1.8\Delta U]^{1/2} + 1$ <sup>2</sup> is of the order of unity in our experiment,  $k_B$  is Boltzmann's constant, and  $T$  is the temperature. In the quantum regime ( $k_B T \ll \hbar\omega_p$ ), to lowest order in  $1/Q$  the escape is predicted to occur via macroscopic quantum tunneling at a rate<sup>11</sup>