Scattering Rate from the Structure Function of Superfluid ³He

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Signal Event Rate

We want a signal rate that takes into account the many-body physics of superfluid 3 He. A DM interaction potential $\mathcal{V}_{int.}(x) = (\mathcal{V}_{Atom} * n)(x)$, gives

$$\begin{split} \frac{dR}{d\omega} &= \frac{\rho_{\chi}}{\rho_{T} m_{\chi}} \frac{\pi \bar{\sigma}}{\mu^{2}} \left\langle \int \frac{d^{3}q}{(2\pi)^{3}} F(q) \underbrace{\frac{2\pi}{V} \sum_{f} |\langle f | \hat{n}_{T}(-q) | i \rangle|^{2} \delta(\Delta E_{fi} - \omega_{q})}_{S(q, \omega_{q})} \delta(\omega - \omega_{q}) \right\rangle_{v} \\ &= \frac{\rho_{\chi}}{\rho_{T} m_{\chi}} \frac{\bar{\sigma}}{4\pi \mu^{2}} \int q \, dq \left[S(q, \omega) F(q) \, \eta(v_{\min}(q, \omega)) \right] \end{split}$$

Signal Event Rate

We want a signal rate that takes into account the many-body physics of superfluid ³He. A DM interaction potential $V_{\text{int.}}(x) = (V_{\text{Atom}} * n)(x)$, gives

$$\begin{split} \frac{dR}{d\omega} &= \frac{\rho_{\chi}}{\rho_{T} m_{\chi}} \frac{\pi \bar{\sigma}}{\mu^{2}} \left\langle \int \frac{d^{3}q}{(2\pi)^{3}} F(q) \underbrace{\frac{2\pi}{V} \sum_{f} |\langle f | \hat{n}_{T}(-q) | i \rangle|^{2} \delta(\Delta E_{fi} - \omega_{q})}_{S(q, \omega_{q})} \delta(\omega - \omega_{q}) \right\rangle_{v} \\ &= \frac{\rho_{\chi}}{\rho_{T} m_{\chi}} \frac{\bar{\sigma}}{4\pi u^{2}} \int q \, dq \left[S(q, \omega) F(q) \, \eta(v_{\min}(q, \omega)) \right] \end{split}$$

The Classical Gas has

$$S(q,\omega) = 2\pi \frac{N_A}{V} \delta\left(\frac{q^2}{2m_N} - \omega\right), \text{ s.t. } \frac{dR}{d\omega} = \frac{\rho_\chi \bar{\sigma}}{2m_\chi \mu^2} \eta(v_{\min}(q,\omega)) F(q) \bigg|_{q=\sqrt{2m_N\omega}}$$

Coherence Factor

The DM sees the atoms; inital and final states are described in terms of quasiparticles.

$$\implies \hat{n}_{-q} \supset \sum_{c,d} U_{p-q,\uparrow c}^* V_{p,\uparrow d} \, \hat{\alpha}_{p-q,c}^\dagger \hat{\alpha}_{-p,d}^\dagger + U_{p-q,\downarrow c}^* V_{p,\downarrow d} \, \hat{\alpha}_{p-q,c}^\dagger \hat{\alpha}_{-p,d}^\dagger$$

For the s-wave pairing, $\Delta_{\mathbf{k}} = \Delta_0(i\sigma_2)$, we reproduce the result of Hochberg et al. 2023

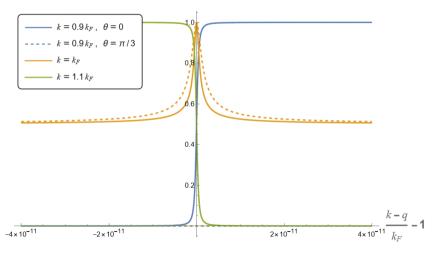
$$|\langle f; -\boldsymbol{p} - \boldsymbol{q}, \boldsymbol{p} | \hat{n}_{-q} | i \rangle|^2 = 1 - \frac{\xi_p \xi_{-p-q}}{E_p E_{-p-q}} + \frac{\Delta^2}{E_p E_{-p-q}}$$

For the B-phase, p-wave pairing, $\Delta_{\mathbf{k}} = \Delta_0 \, \hat{\mathbf{k}} \cdot \boldsymbol{\sigma} \, (i\sigma_2)$,

$$|\langle f; -\boldsymbol{p} - \boldsymbol{q}, \boldsymbol{p} | \hat{n}_{-\boldsymbol{q}} | i \rangle|^2 = 1 - \frac{\xi_{\boldsymbol{p}} \xi_{-\boldsymbol{p}-\boldsymbol{q}}}{E_{\boldsymbol{p}} E_{-\boldsymbol{p}-\boldsymbol{q}}} + \frac{\Delta^2}{E_{\boldsymbol{p}} E_{-\boldsymbol{p}-\boldsymbol{q}}} \frac{\boldsymbol{p} \cdot (-\boldsymbol{p} - \boldsymbol{q})}{|\boldsymbol{p}| |\boldsymbol{p} + \boldsymbol{q}|}.$$

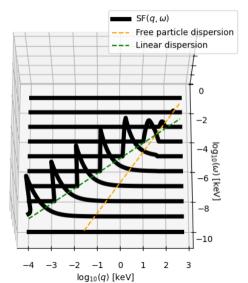
Coherence Factor





Structure Function

- From the Coherence factor we can evaluate the structure function and find the differential scattering rate and power injection.
- This reproduces the classical structure function we have the usual scattering at high-q (near the threshold).
- The superfluid gap Δ provides an absolute cutoff corresponding to $m_\chi \simeq 10^{-4} {\rm eV}.$



Scope

- We have a more complete description of nuclear scattering signal that verifies the classical limit.
- This can be used for various DM scattering interactions. Would need to be extended for DM absorption rate.
- This can be built upon to be used alongside the copper heating calculations to investigate sensitivity 'sub-threshold' masses through power injection.

