

# Impact of the meson-exchange currents on magnetic dipole moments in odd near doubly magic nuclei

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Collaborators: R. Han, B. C. Backes, G. Danneaux, J. Dobaczewski,  
W. Jiang, and M. Kortelainen

Nuclear Physics Conference 2025  
April 24, 2025



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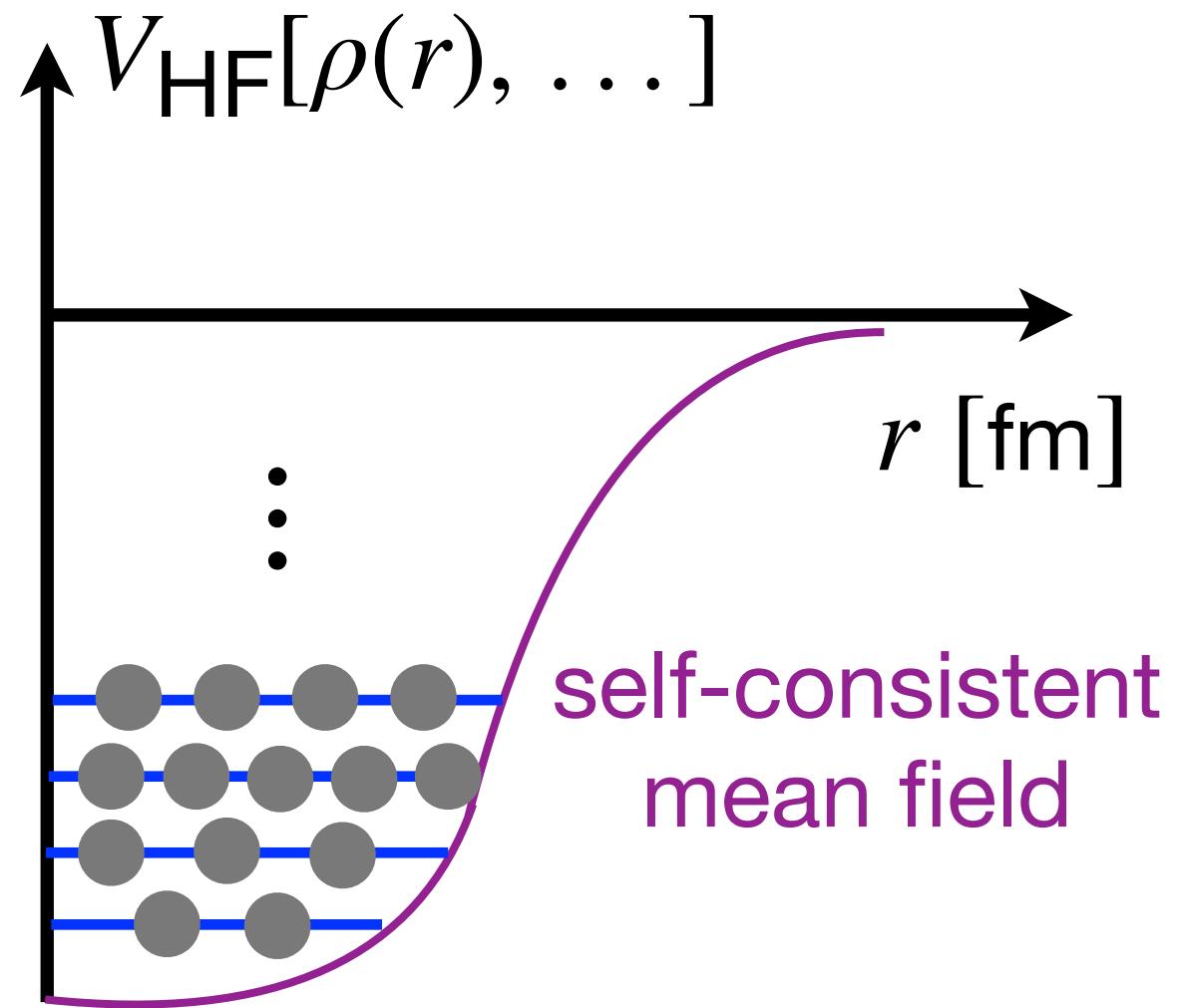
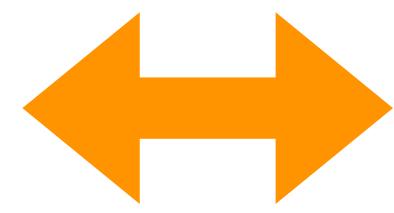
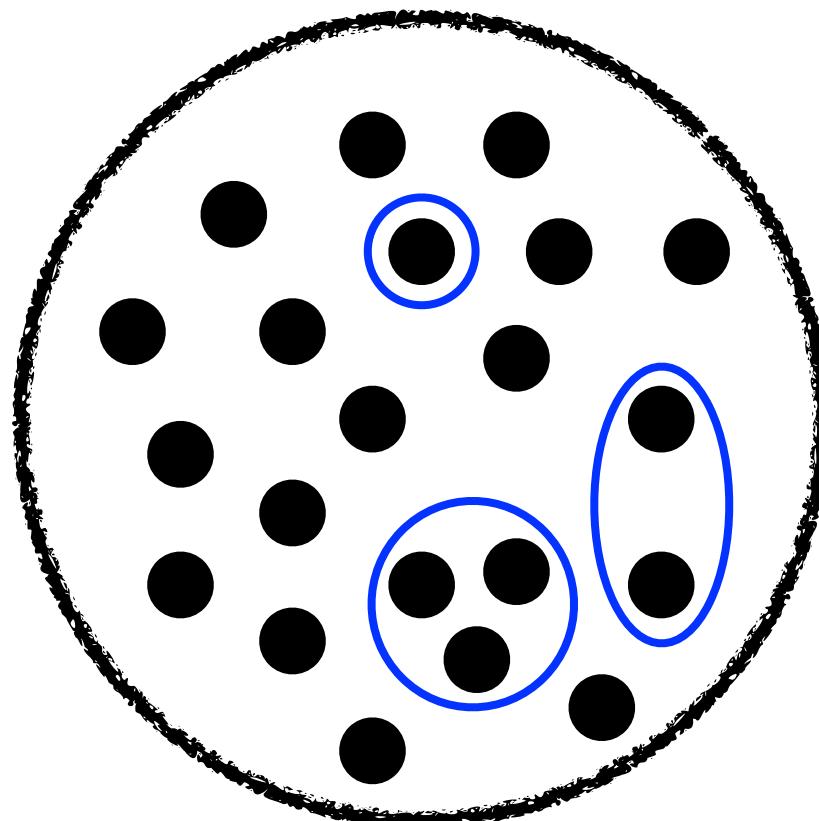
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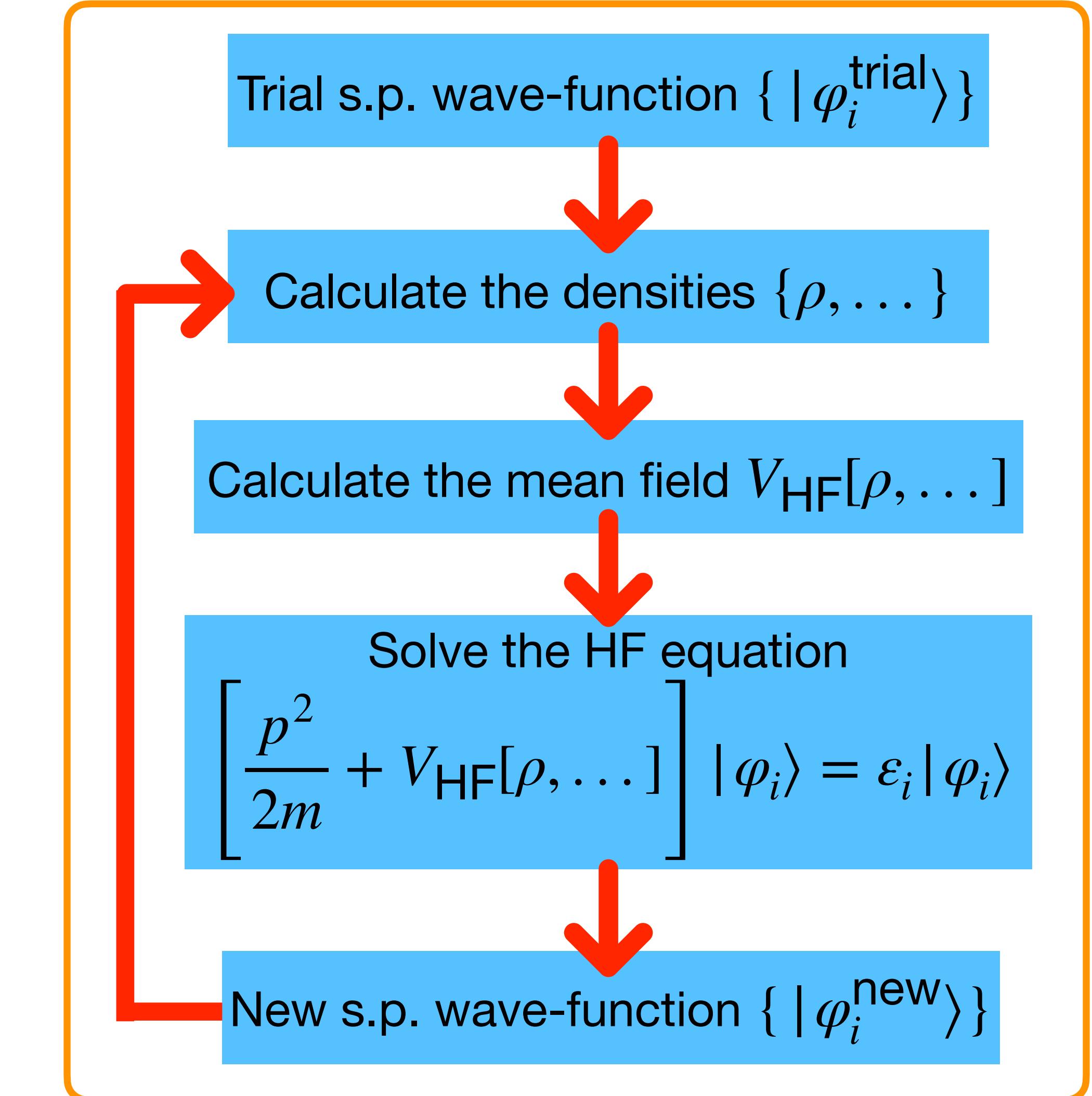
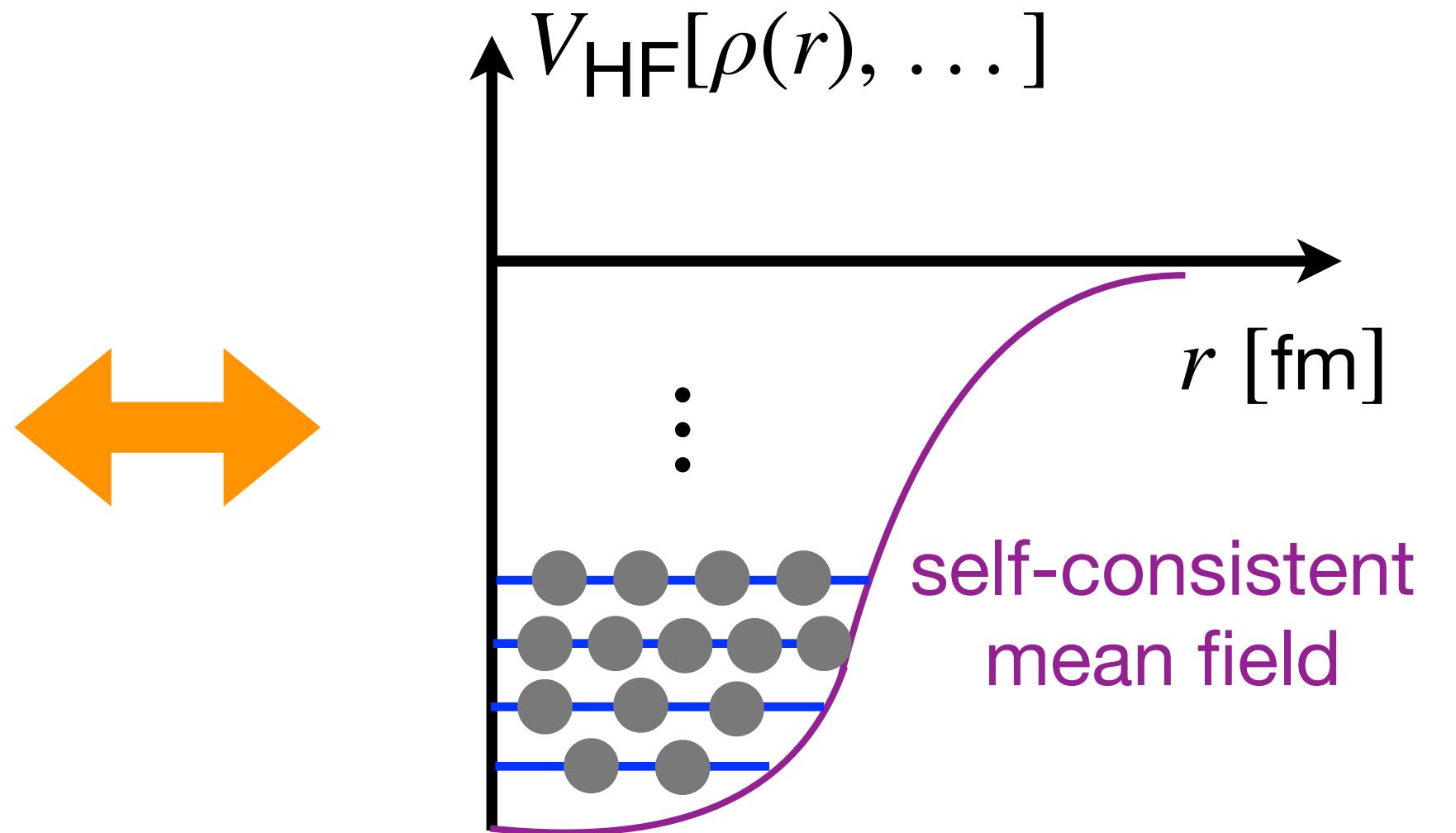
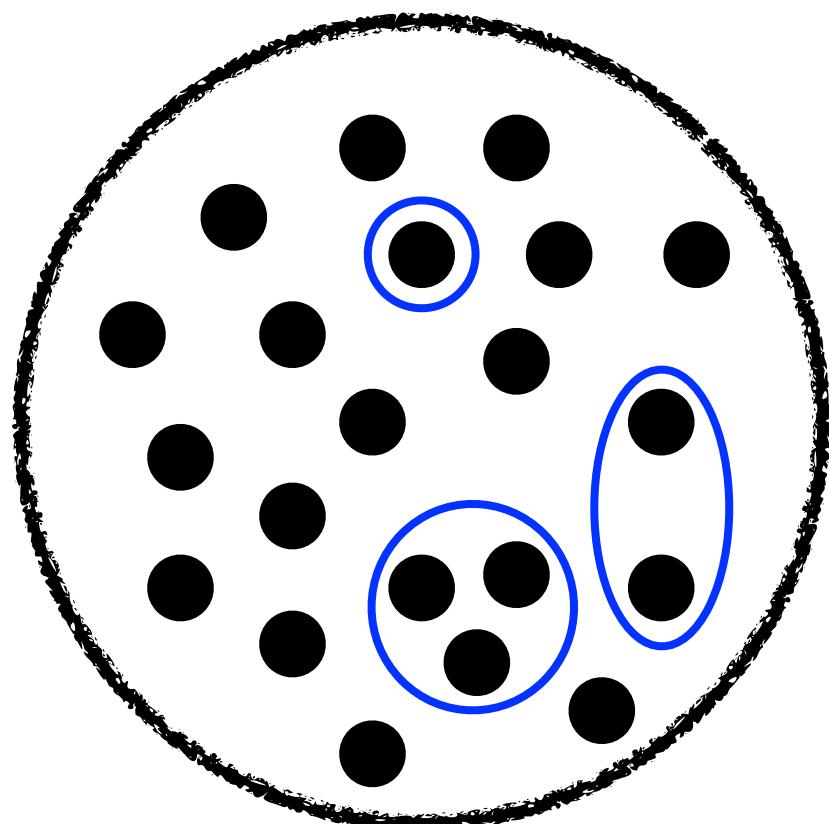
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# Nuclear density functional theory

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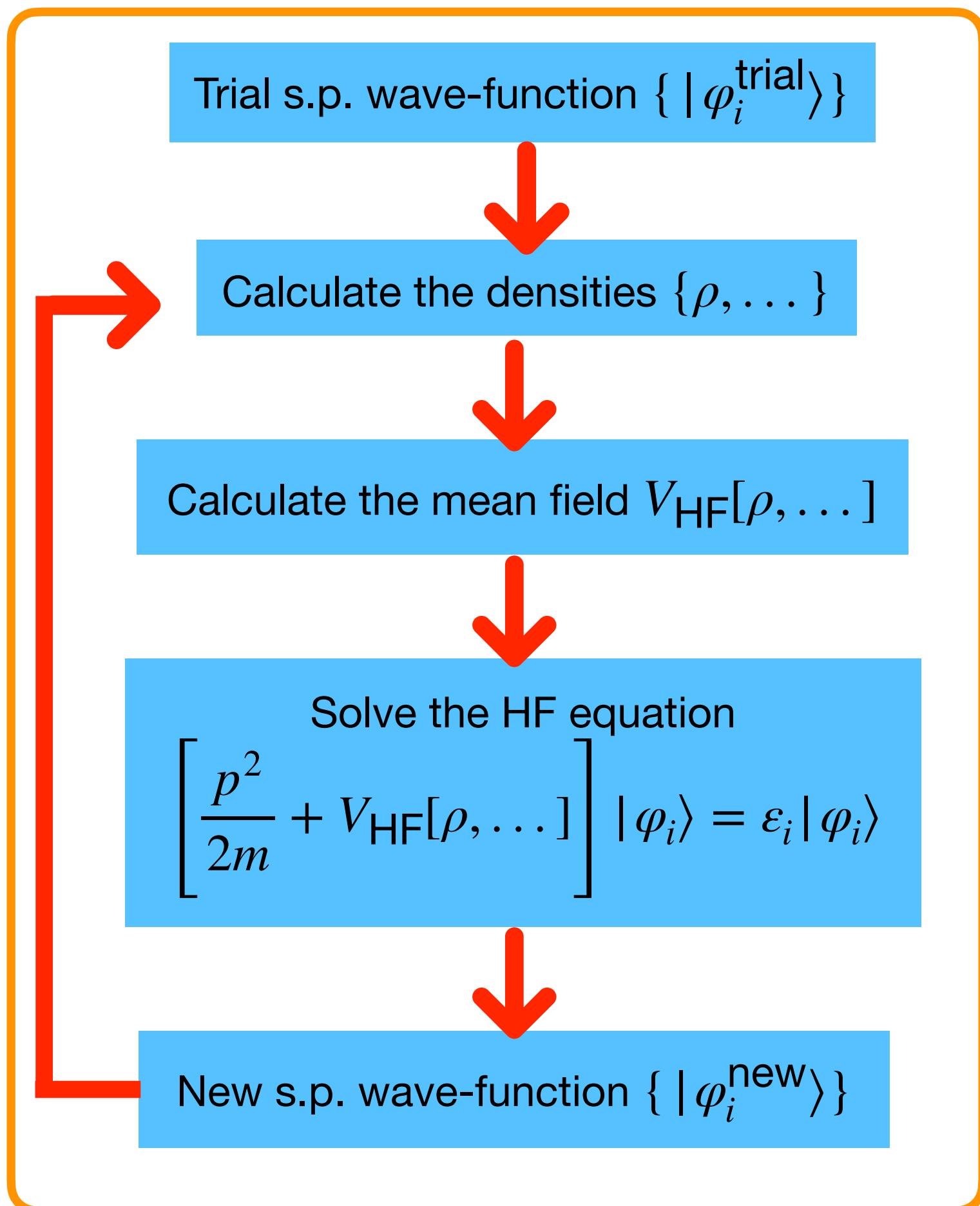
# Nuclear density functional theory



# Nuclear DFT methodology

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## Self-consistency



# Nuclear DFT methodology

## Self-consistency

Trial s.p. wave-function  $\{ |\varphi_i^{\text{trial}}\rangle \}$

Calculate the densities  $\{ \rho, \dots \}$

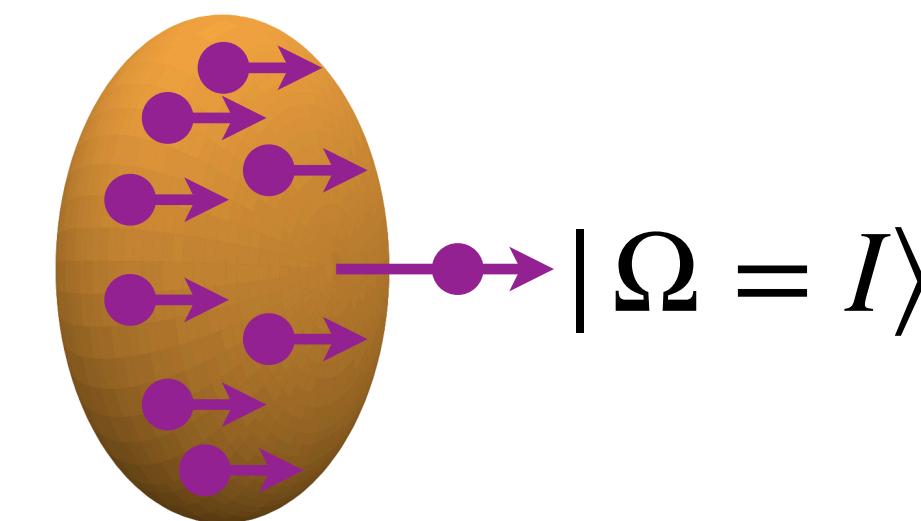
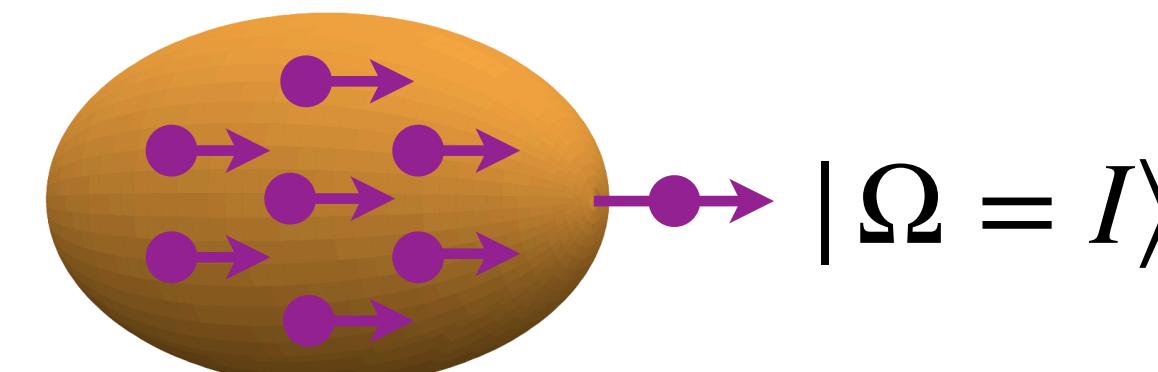
Calculate the mean field  $V_{\text{HF}}[\rho, \dots]$

Solve the HF equation

$$\left[ \frac{p^2}{2m} + V_{\text{HF}}[\rho, \dots] \right] |\varphi_i\rangle = \varepsilon_i |\varphi_i\rangle$$

New s.p. wave-function  $\{ |\varphi_i^{\text{new}}\rangle \}$

## Spin-core polarization



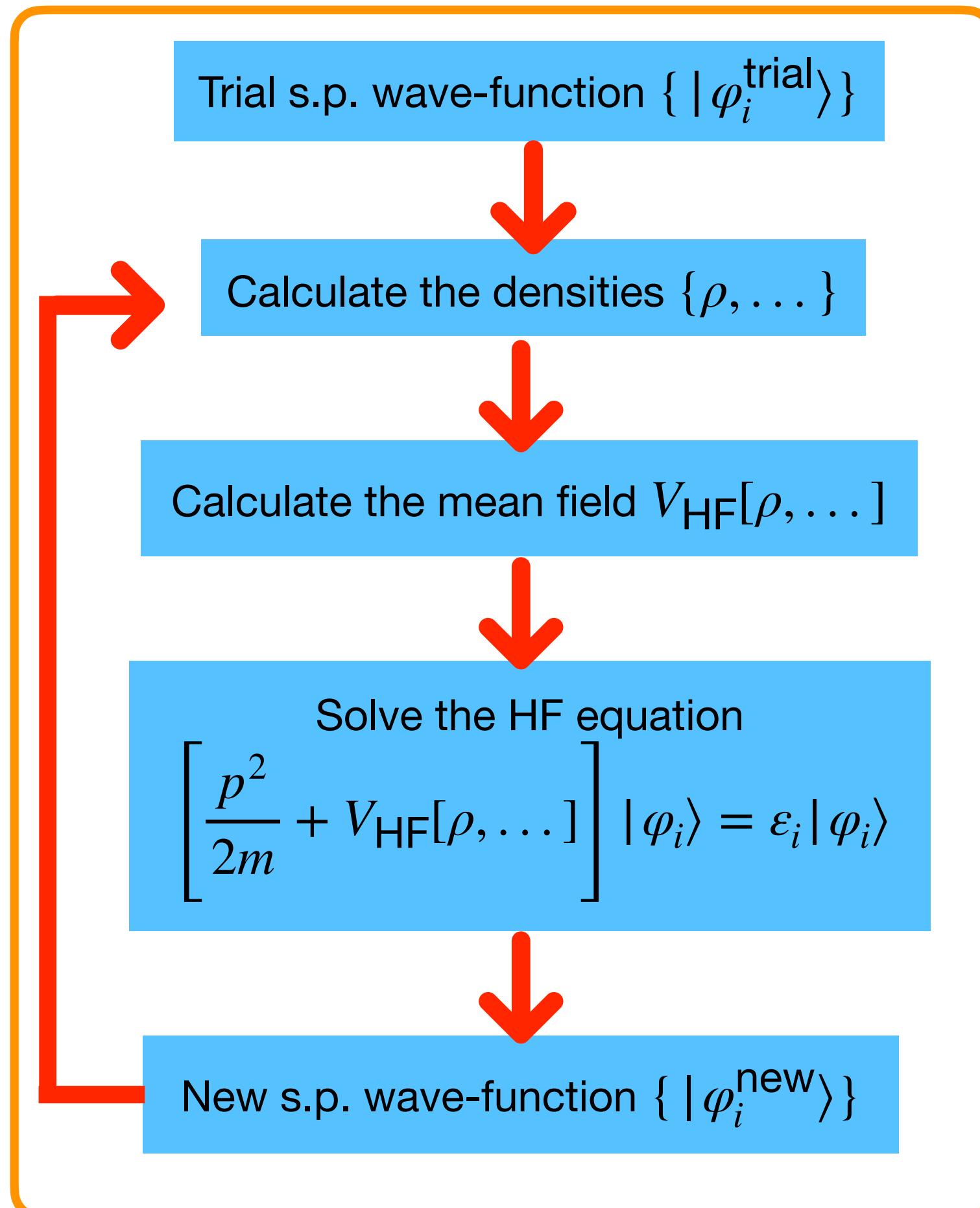
Landau parameter  $g'_0$

$$g'_0 = N_0 \left( 2C_1^S + 2C_1^T (3\pi^2 \rho_0/2)^{2/3} \right)$$

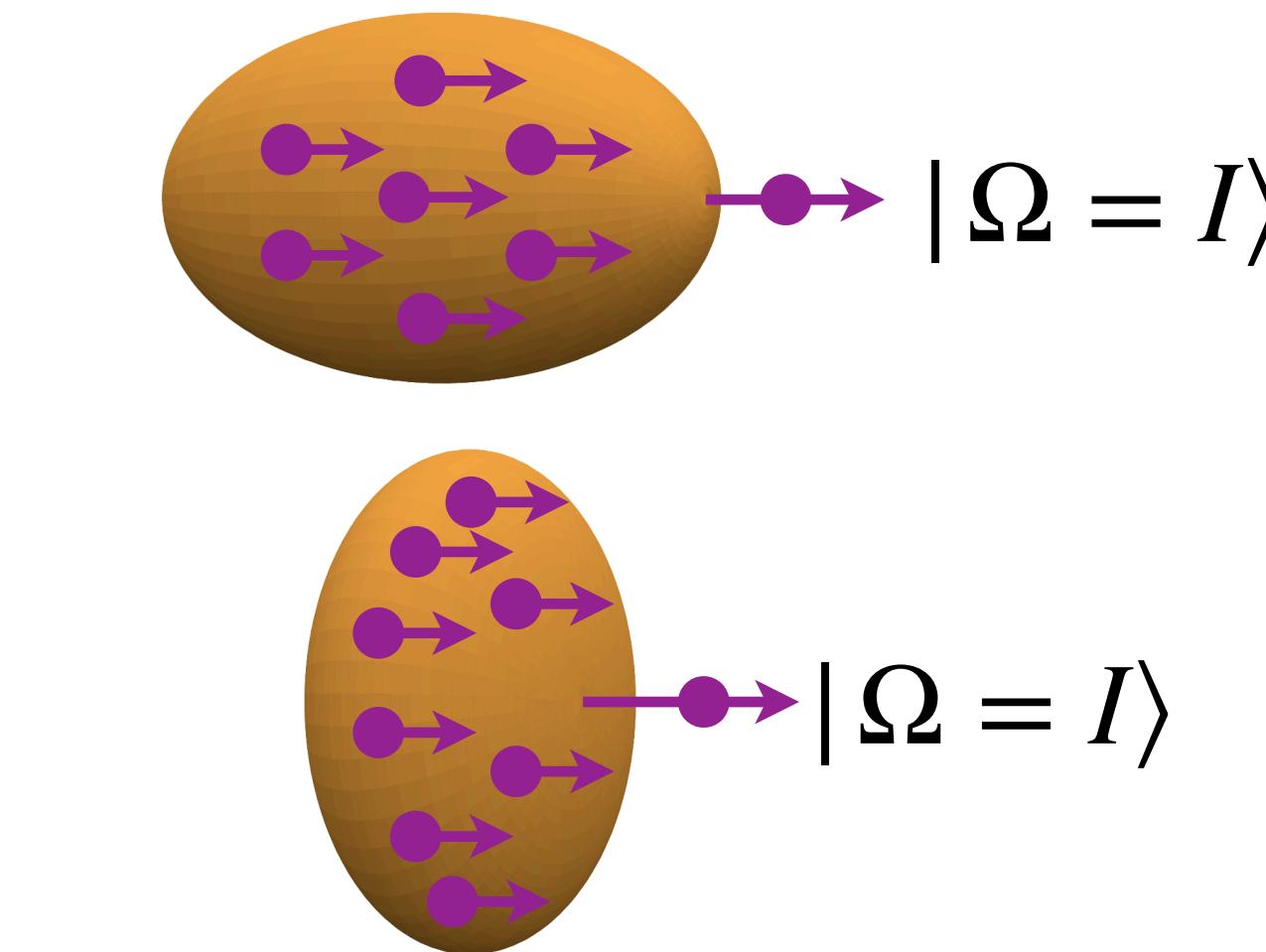
$$\frac{1}{N_0} \approx 150 \frac{m}{m^*} \text{ MeV} \cdot \text{fm}^3$$

# Nuclear DFT methodology

## Self-consistency



## Spin-core polarization

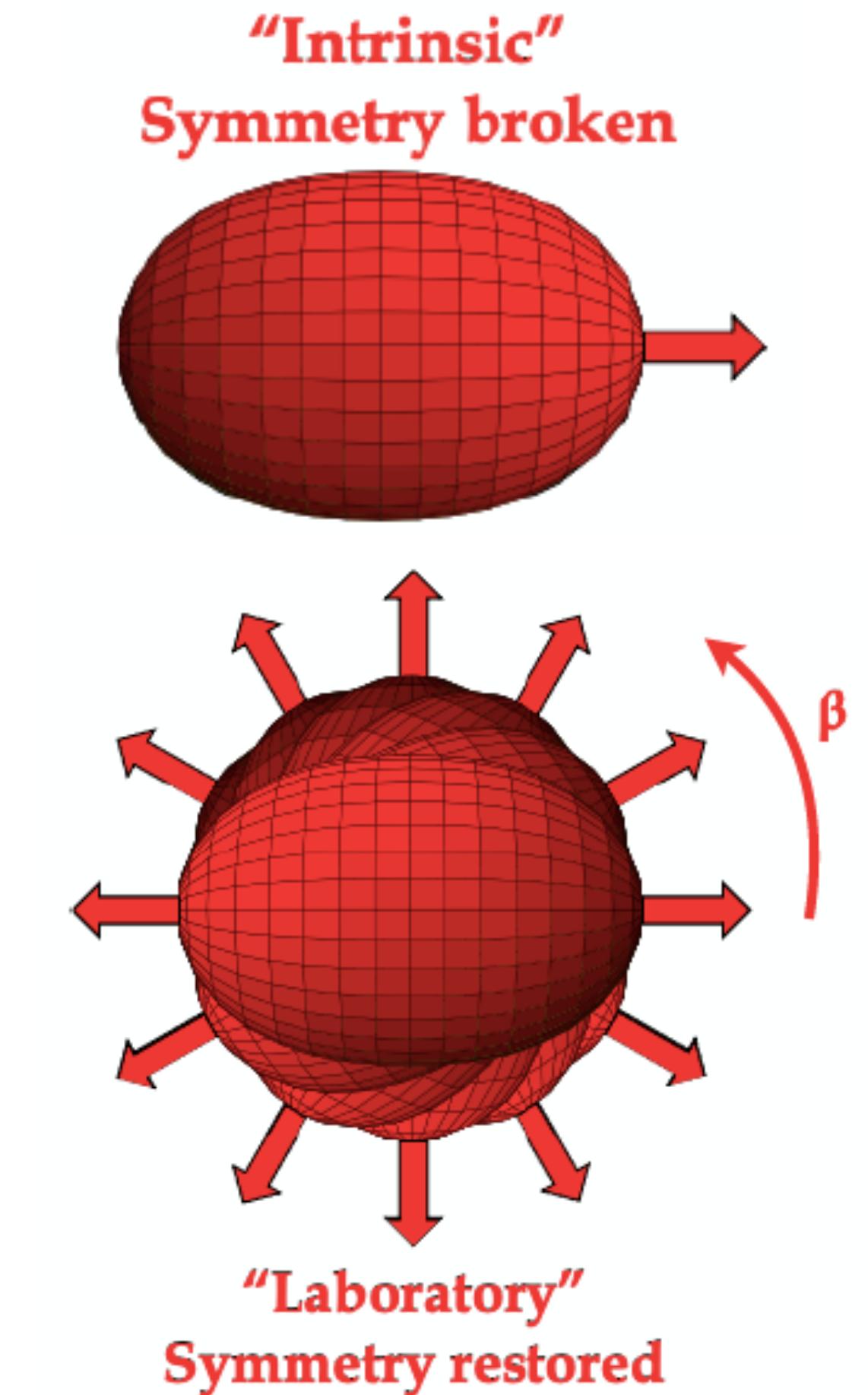


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## Symmetry restoration



# Nuclear magnetic dipole moments

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Spectroscopic magnetic dipole  $\mu$  moment is defined as:

$$\mu = \sqrt{\frac{4\pi}{3}} \langle II | \hat{M}_{10} | II \rangle,$$

P. Ring and P. Schuck, *The Nuclear Many-Body Problem*

$|II\rangle$  = angular momentum projected (AMP) states

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$$\hat{M}_{10} = \sqrt{\frac{3}{4\pi}} \mu_N \sum_{i=1}^A \left\{ g_s^{(i)} s_{zi} + g_\ell^{(i)} \ell_{zi} \right\},$$

$$g_s^{(i)} = g_s^{(\pi)}(g_s^{(\nu)}) = 5.59(-3.83), \quad g_\ell^{(i)} = g_\ell^{(\pi)}(g_\ell^{(\nu)}) = 1(0).$$

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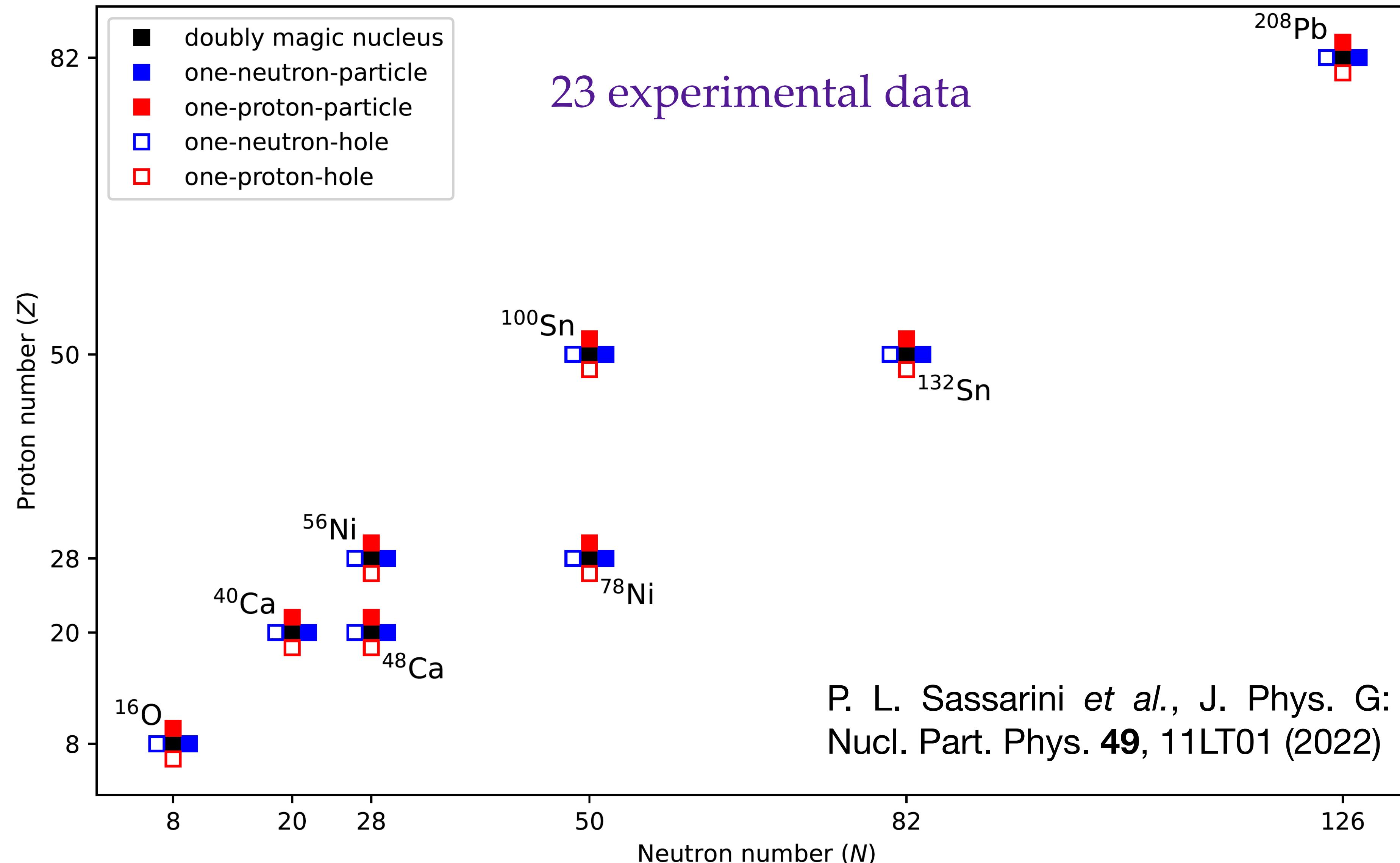
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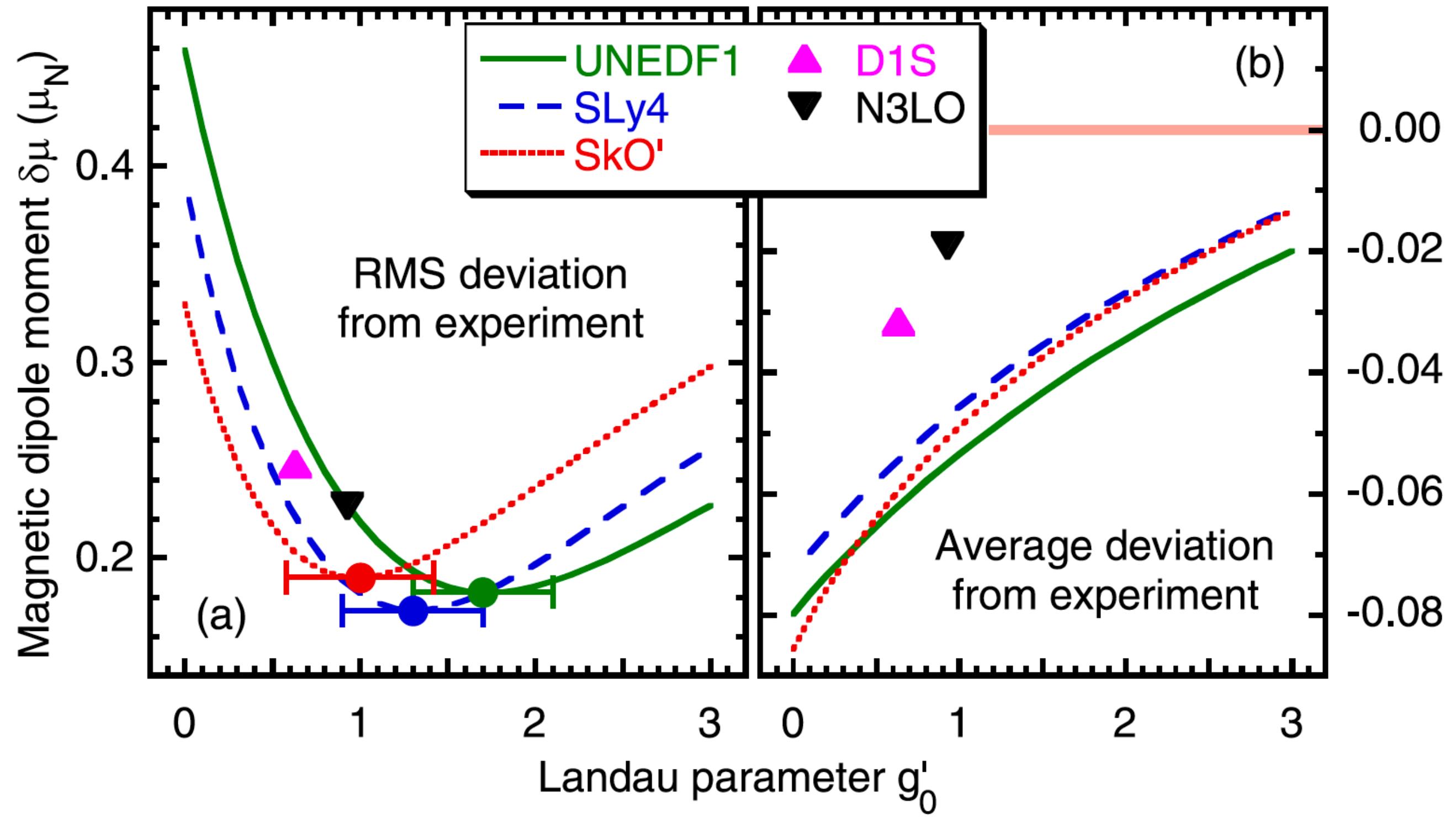
Since the self-consistent polarizations act in the full single-particle space, no effective  $g$ -factors are needed!

# Odd near doubly magic nuclei



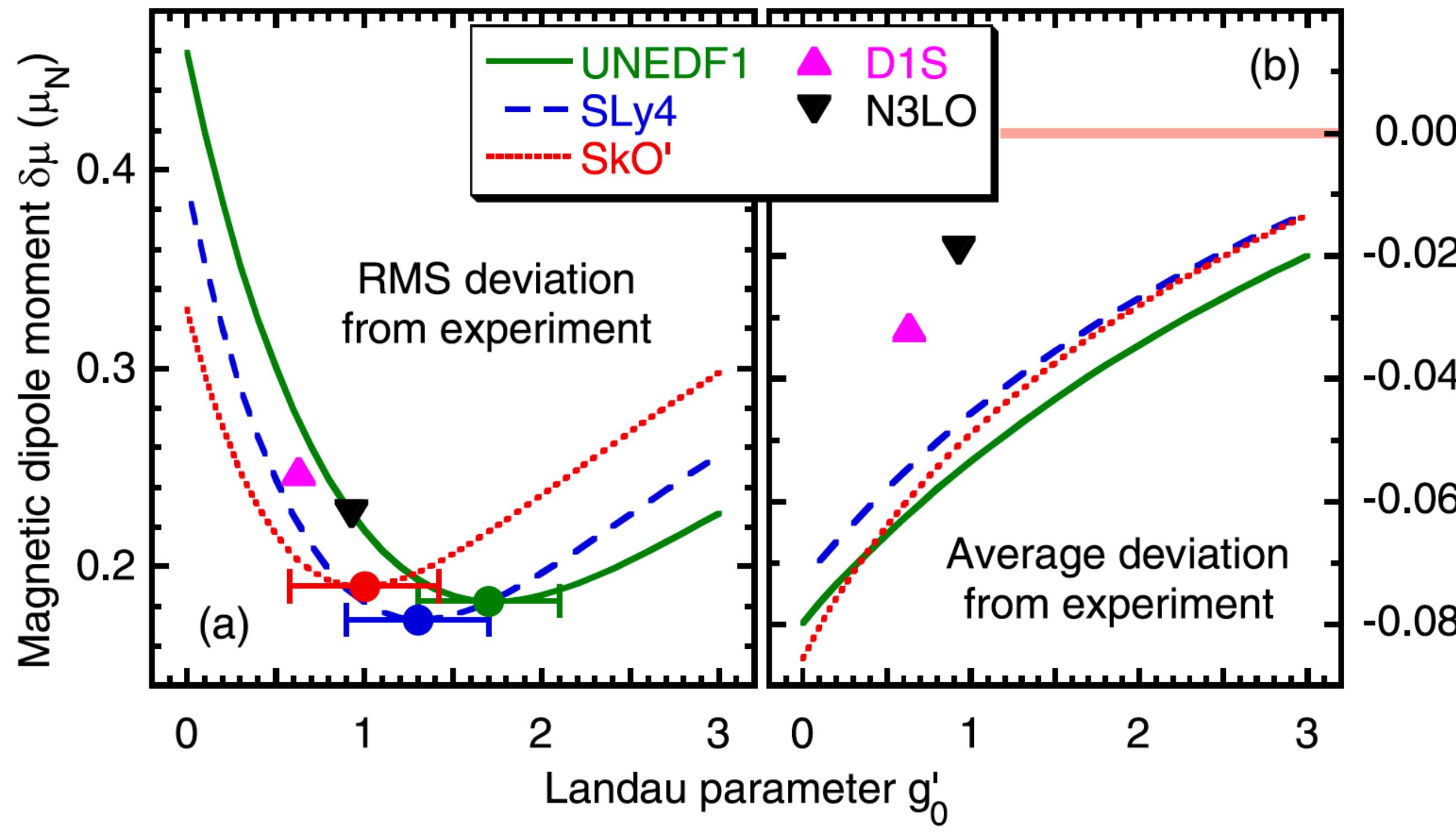
# Optimum Landau parameter $g'_0$

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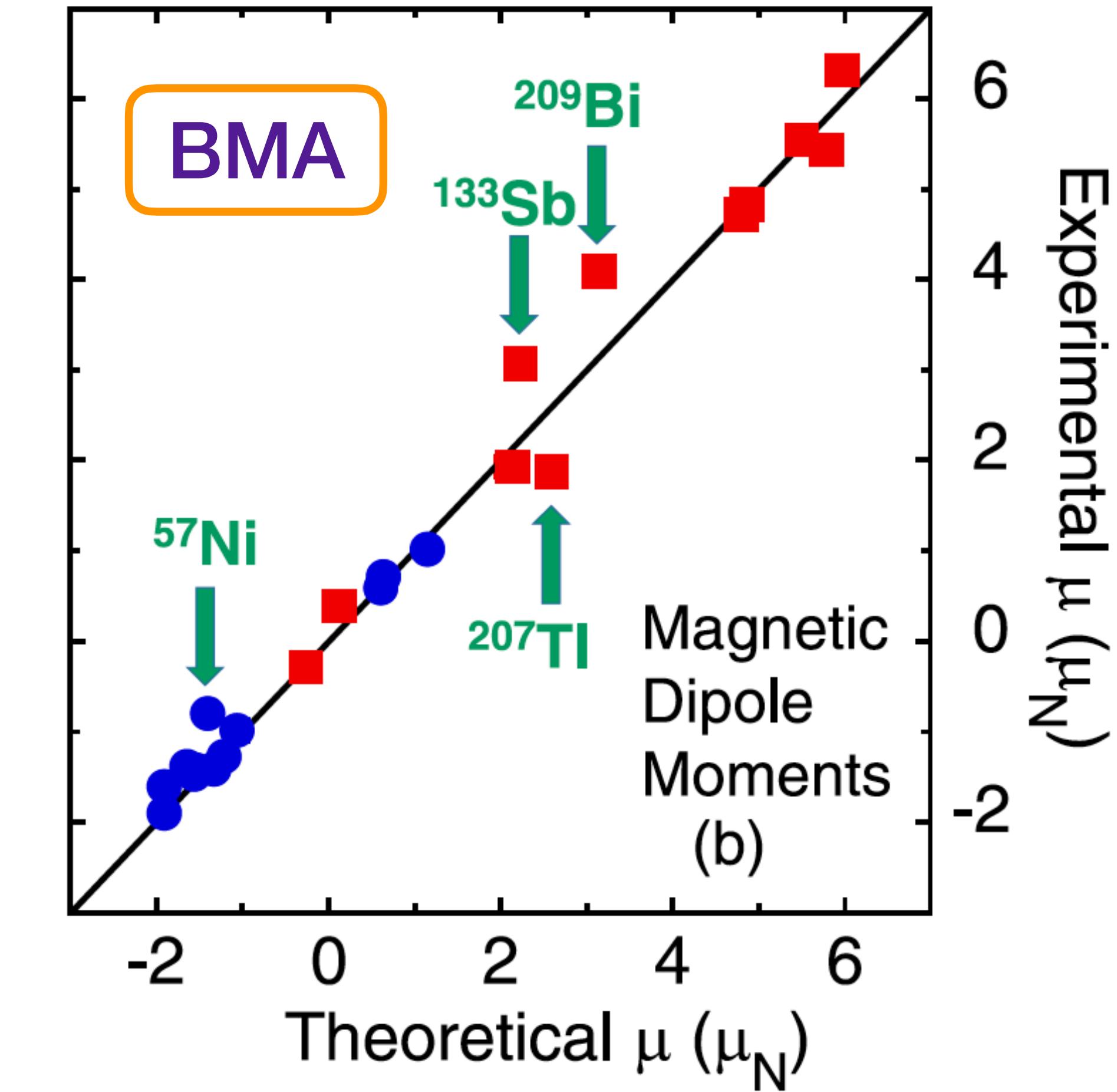


$$g'_0 = 1.7(4) \text{ (UNEDF1)}$$

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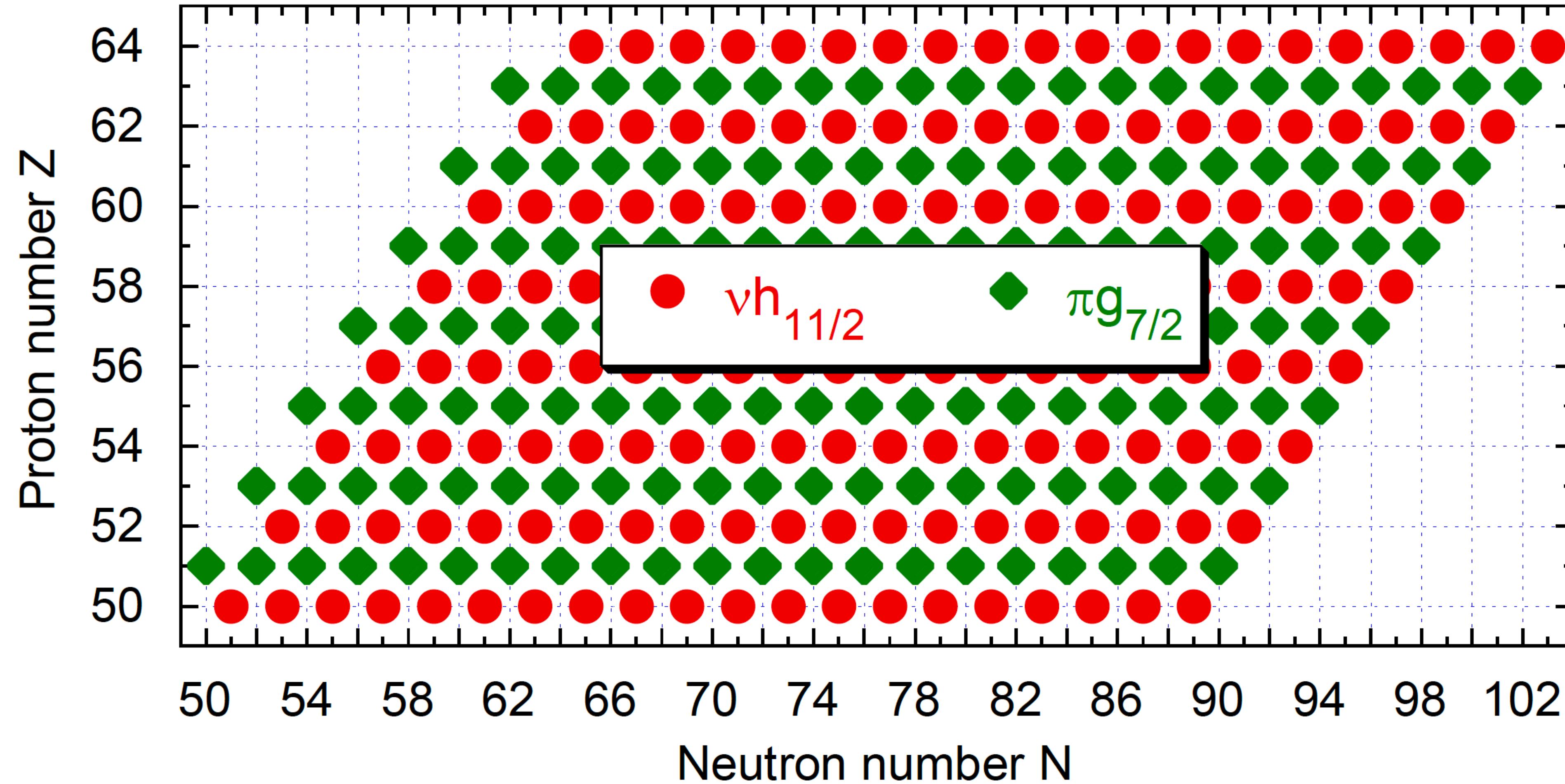


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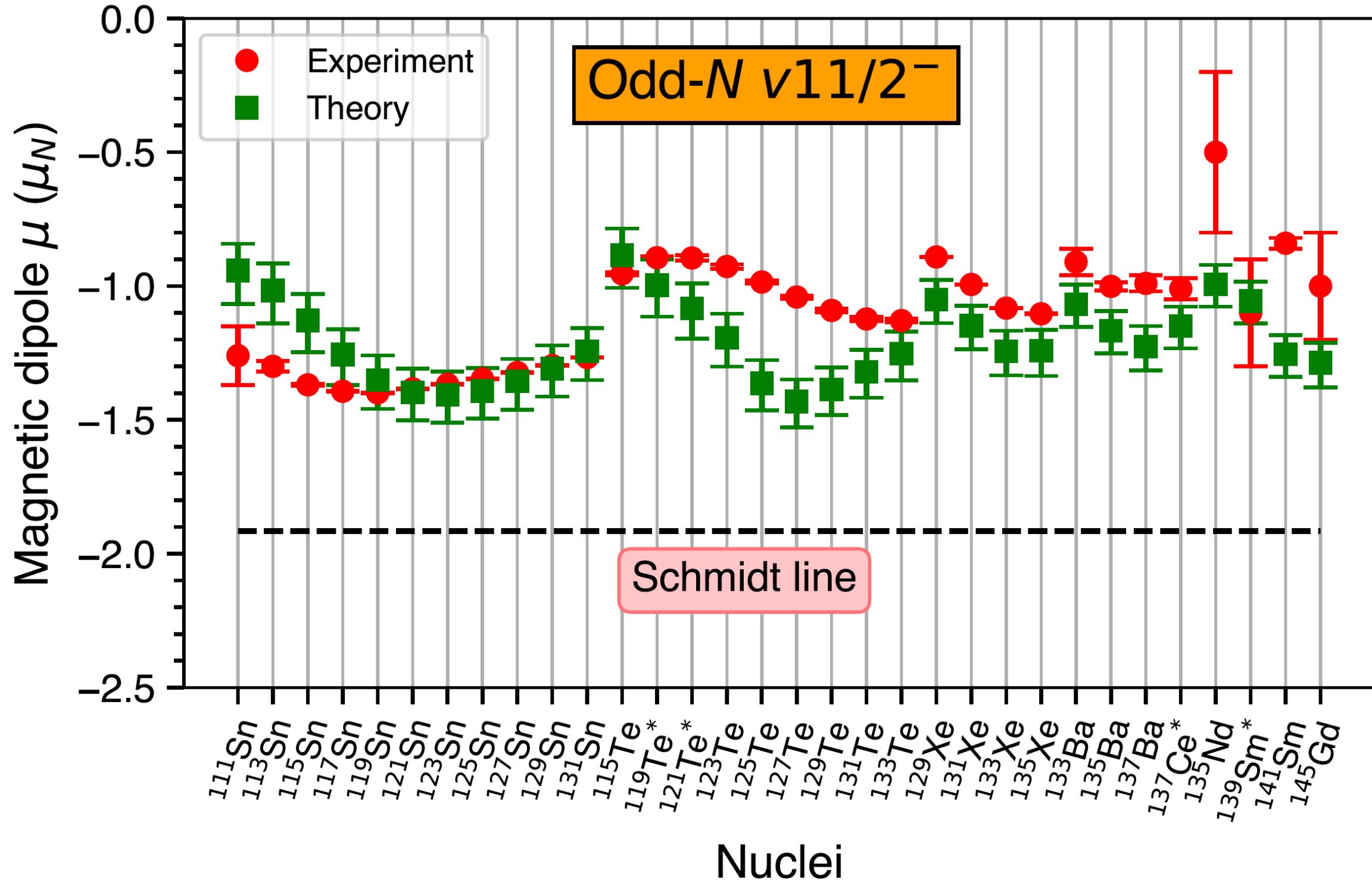


- Z-odd nuclei
- N-odd nuclei

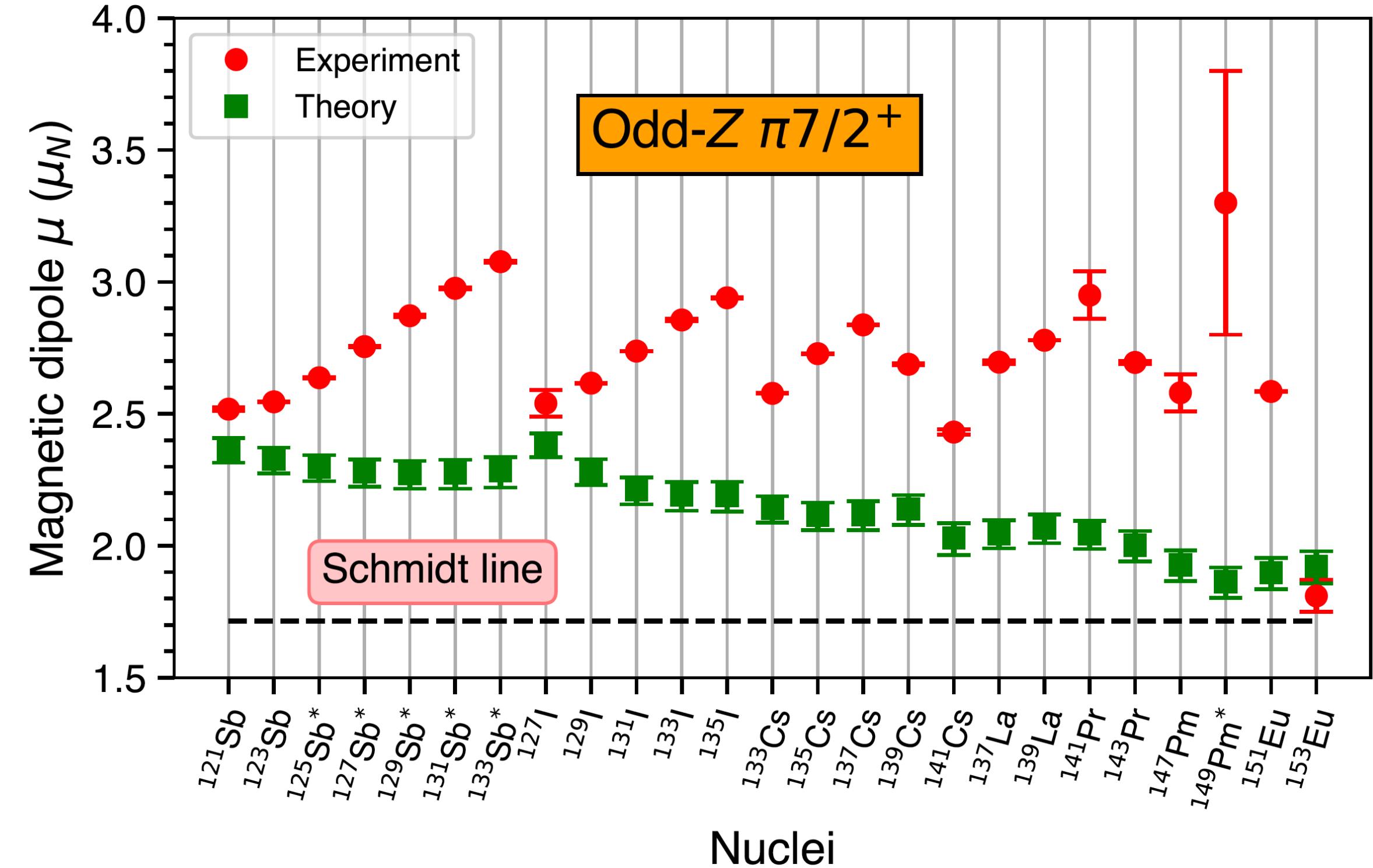
# Systematic nuclear DFT calculations: Sn-Gd



# Theory vs. experiment: Sn-Gd



$$g'_0 = 1.7(4) \text{ (UNEDF1)}$$



N. J. Stone, INDC, report INDC(NDS)-0658

N. J. Stone, INDC, report INDC(NDS)-0794

N. J. Stone, INDC, report INDC(NDS)-0816

Yordanov D. T. et al., Comm. Phys. 3, 107 (2020)

Lechner S. et. al., Phys. Lett. B 847, 138278 (2023)

# Meson-exchange contributions

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Magnetic moment operator:

$$\hat{\mu} = \sum_k^A \hat{\mu}_{1b,k} + \sum_{k < \ell}^A \hat{\mu}_{2b,k\ell} + \dots$$

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Two-body meson-exchange:

$$\hat{\mu}_{2b}^{\text{NLO}}(\mathbf{r}_1, \mathbf{r}_2) = \frac{1}{2} \int d^3 \mathbf{x} \, \mathbf{x} \times \mathbf{j}_{2b}^{\text{NLO}}(\mathbf{x}, \mathbf{r}_1, \mathbf{r}_2)$$

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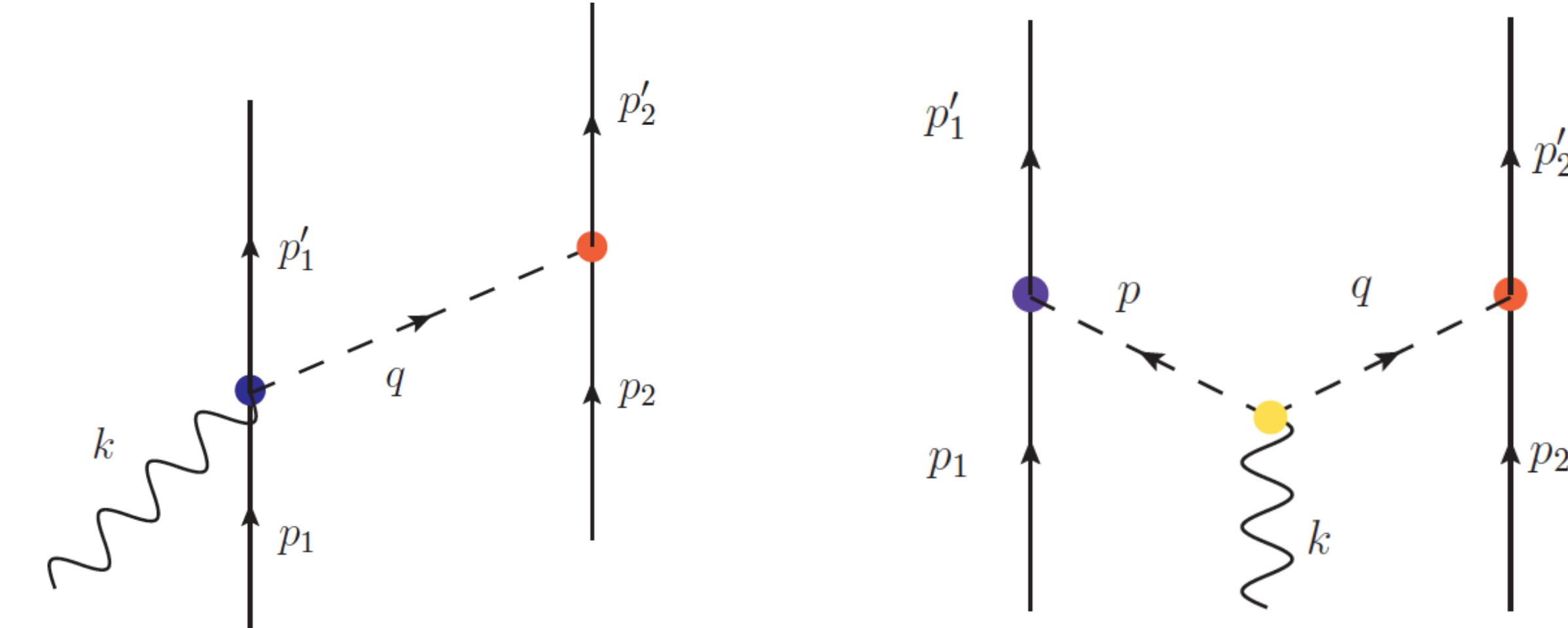
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Next-to-leading order (NLO):



R. Seutin, et al., PRC **108**, 054005 (2023)

T. Miyagi, et al., PRL **132**, 232503 (2024)

Seagull graph

Pion-in-flight graph

# NLO magnetic moment operators

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The NLO intrinsic and Sachs contributions to the magnetic moment operator are given by

$$\hat{\mu}_{2b}^{\text{NLO, int}}(\mathbf{r}) = -\frac{g_A^2 m_\pi}{32\pi F_\pi^2} (\hat{\boldsymbol{\tau}}_1 \times \hat{\boldsymbol{\tau}}_2)_z \left\{ \left( 1 + \frac{1}{m_\pi r} \right) [(\hat{\boldsymbol{\sigma}}_1 \times \hat{\boldsymbol{\sigma}}_2) \cdot \hat{\mathbf{r}}] \hat{\mathbf{r}} - (\hat{\boldsymbol{\sigma}}_1 \times \hat{\boldsymbol{\sigma}}_2) \right\} e^{-m_\pi r}$$

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and  $(g_A = 1.27; F_\pi = 92.3 \text{ MeV}; m_\pi = 138.039 \text{ MeV})$

$$\hat{\mu}_{2b}^{\text{NLO, Sachs}}(\mathbf{r}) = -\frac{1}{2} (\hat{\boldsymbol{\tau}}_1 \times \hat{\boldsymbol{\tau}}_2)_z V_{1\pi}(r) \mathbf{R}_{\text{NN}} \times \mathbf{r},$$

respectively, where

$$V_{1\pi}(r) = \frac{m_\pi^2 g_A^2}{48\pi F_\pi^2} \left[ \hat{S}_{12} \left( 1 + \frac{3}{m_\pi r} + \frac{3}{(m_\pi r)^2} \right) + \hat{\boldsymbol{\sigma}}_1 \cdot \hat{\boldsymbol{\sigma}}_2 \right] \frac{e^{-m_\pi r}}{r}.$$

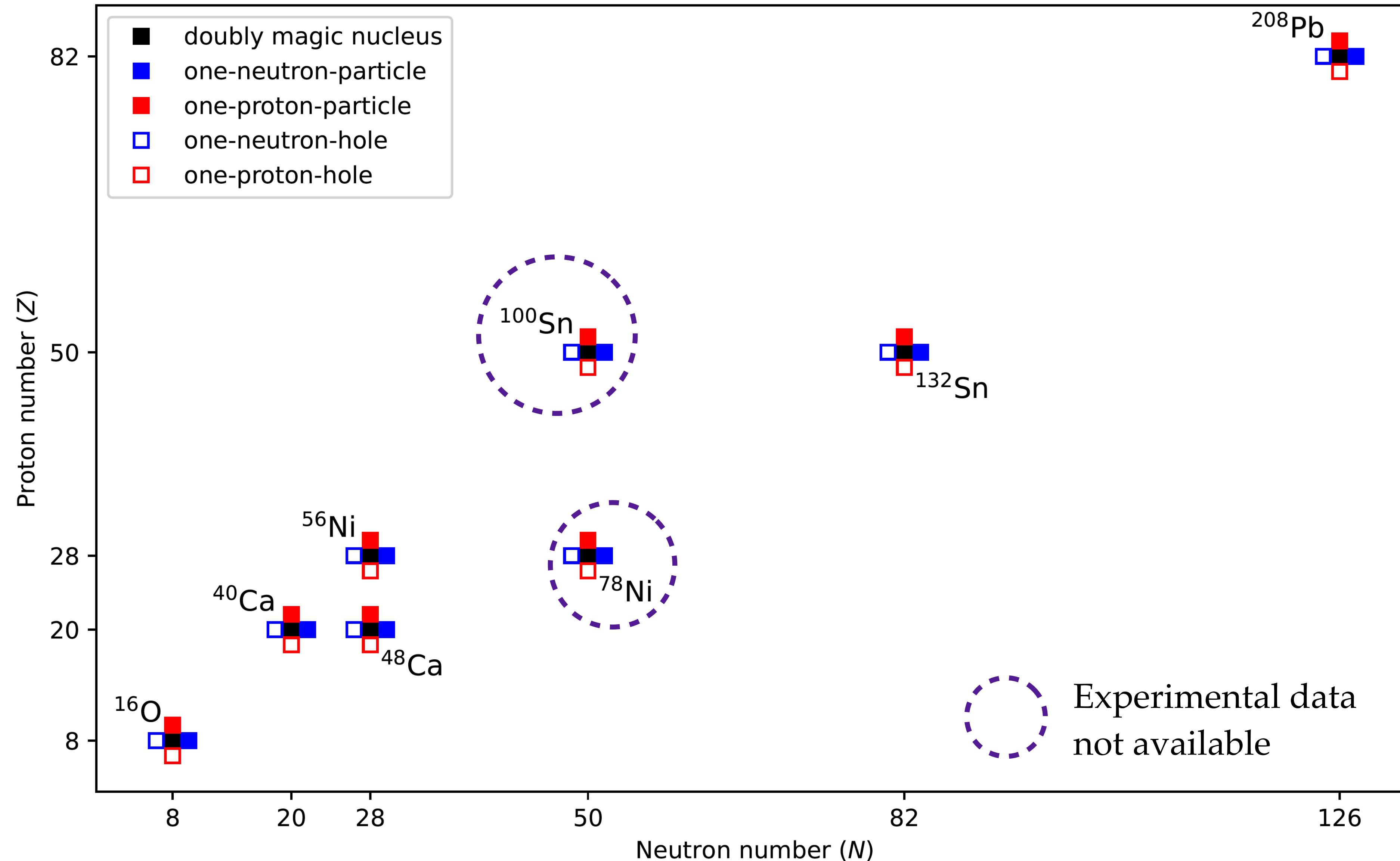
Tensor operator:

$$\hat{S}_{12} = 3 (\hat{\mathbf{r}} \cdot \hat{\boldsymbol{\sigma}}_1) (\hat{\mathbf{r}} \cdot \hat{\boldsymbol{\sigma}}_2) - \hat{\boldsymbol{\sigma}}_1 \cdot \hat{\boldsymbol{\sigma}}_2.$$

relative  
coordinate

COM

# Odd near doubly magic nuclei (two-body currents)



# Optimum Landau parameter $g'_0$ (two-body currents)

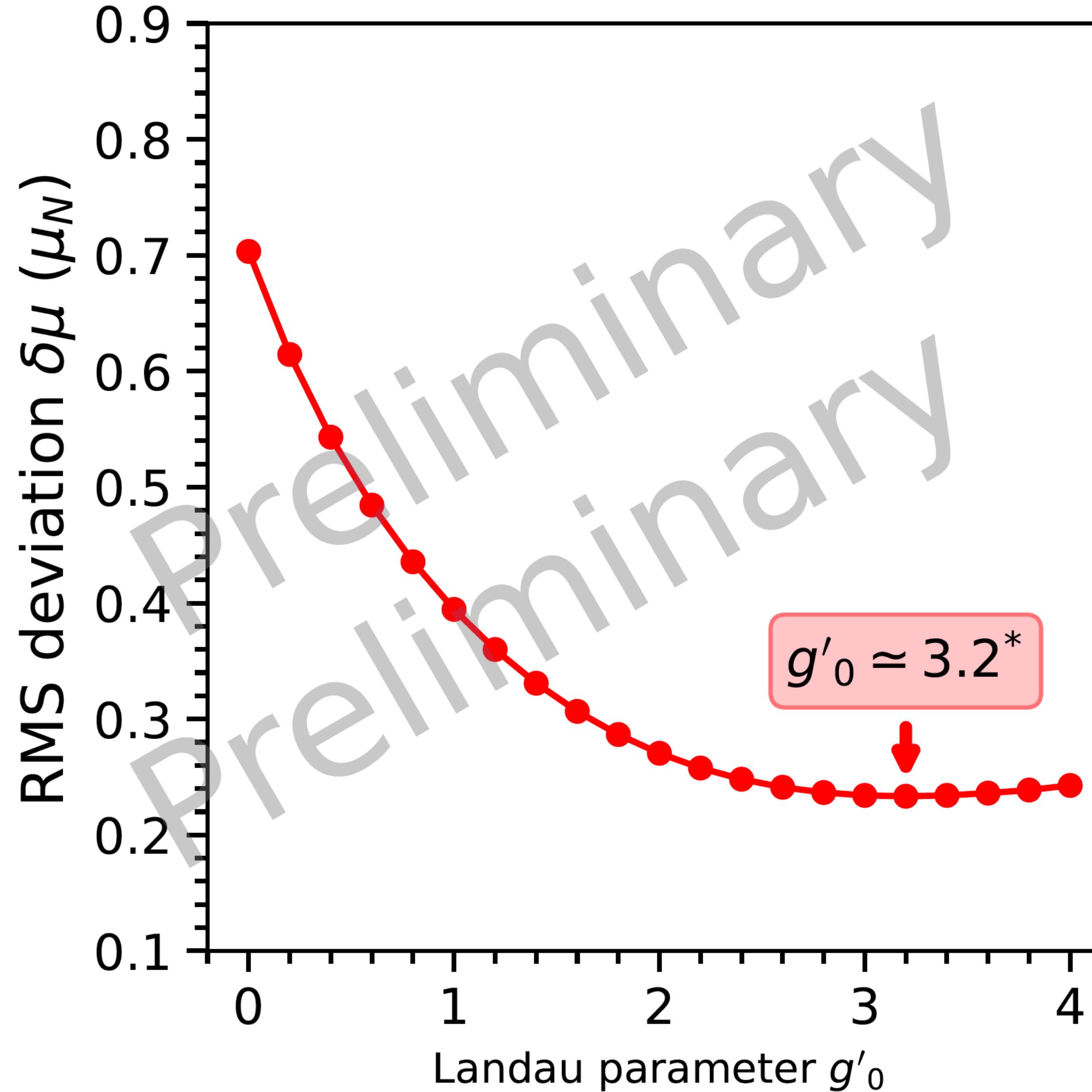
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$$\delta\mu(g'_0) = \sqrt{\frac{1}{N} \sum_{i=1}^N [\mu_{\text{calc}}(i, g'_0) - \mu_{\text{exp}}(i, g'_0)]^2}$$

$i$  = odd nucleus

$N$  = number of odd nuclei

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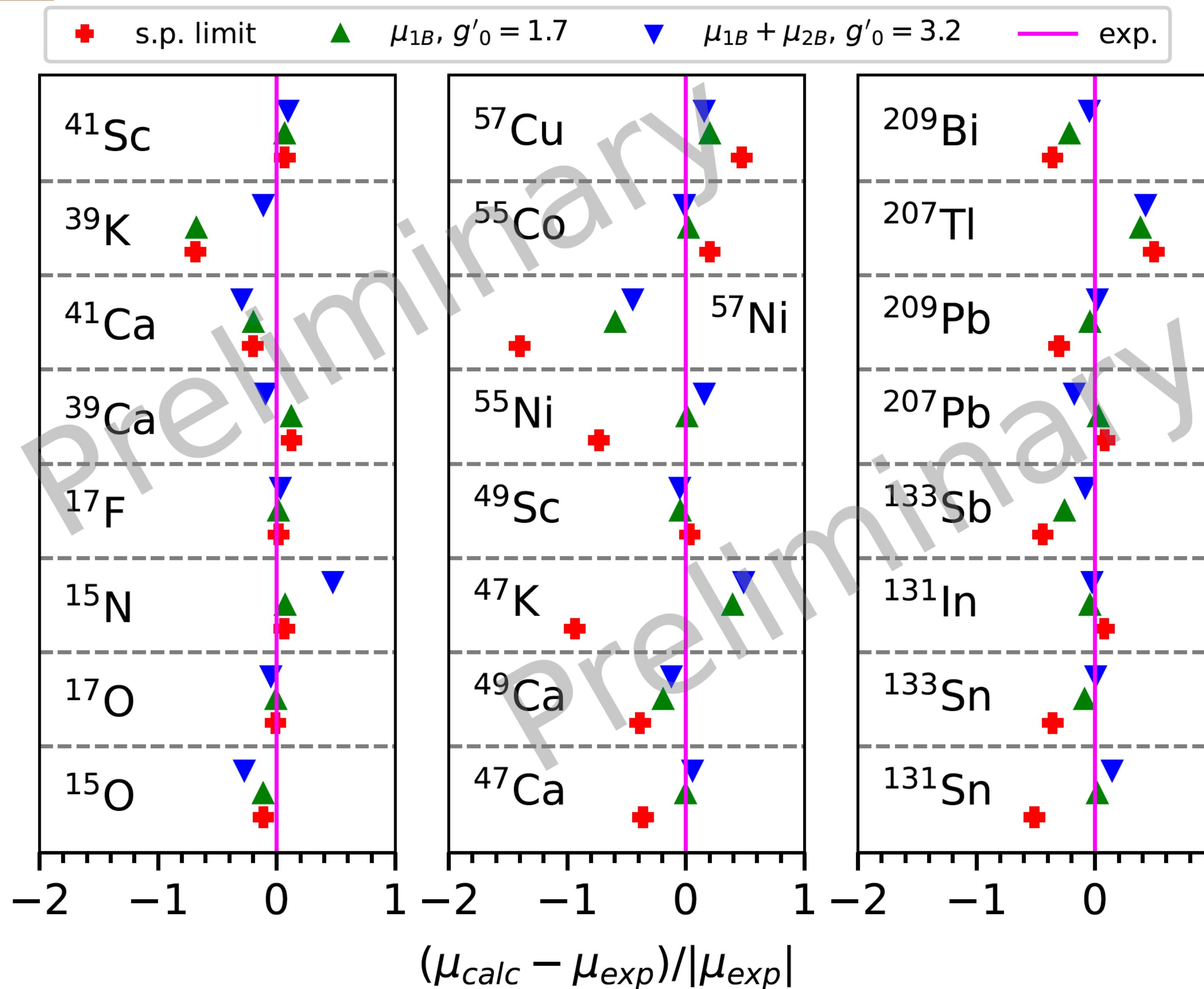
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(\*) 22 odd near doubly magic nuclei without Bayesian analysis (**for now!**).

H. Wibowo, et.al.,  
to be published

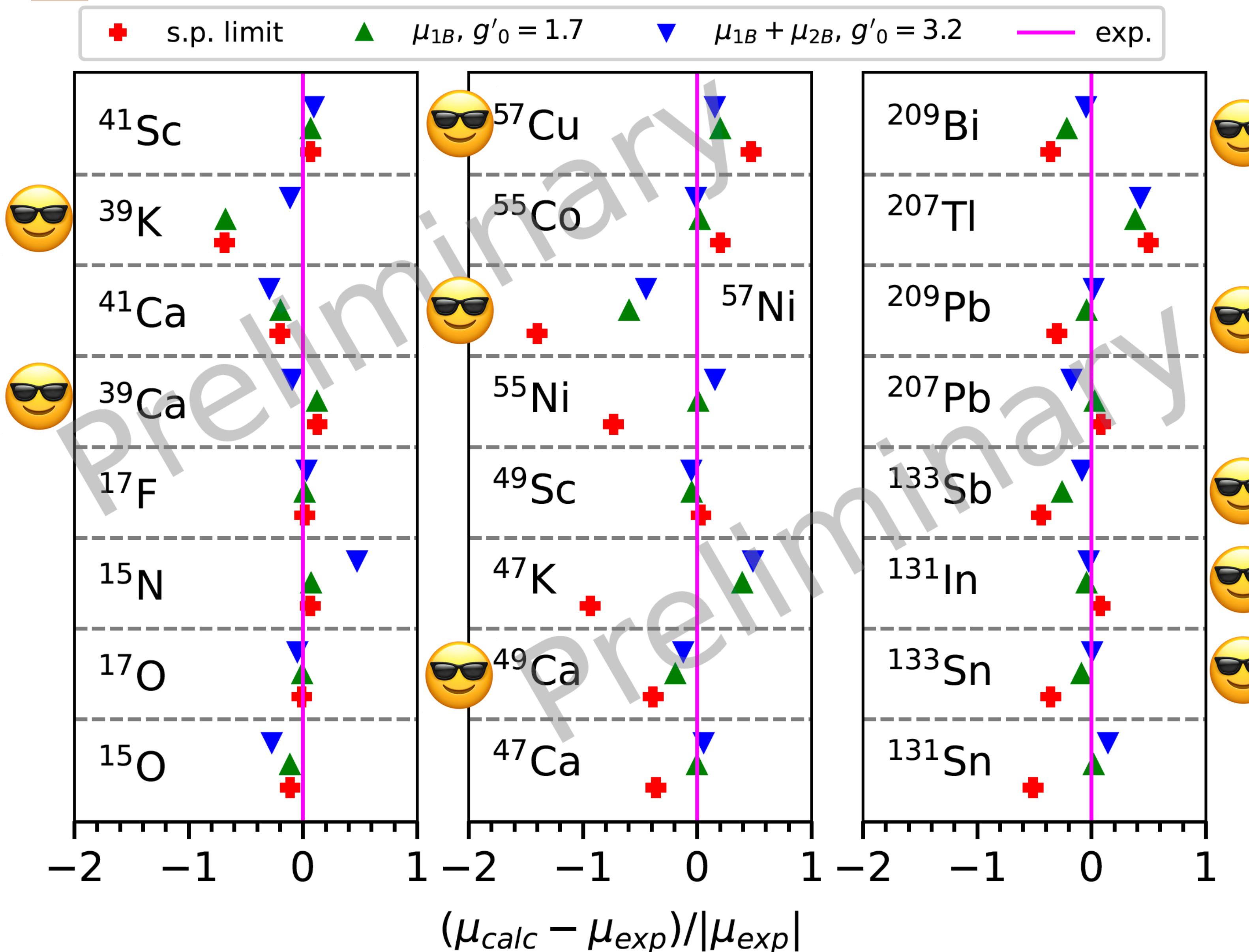
# Magnetic dipole moments: theory vs. experiment

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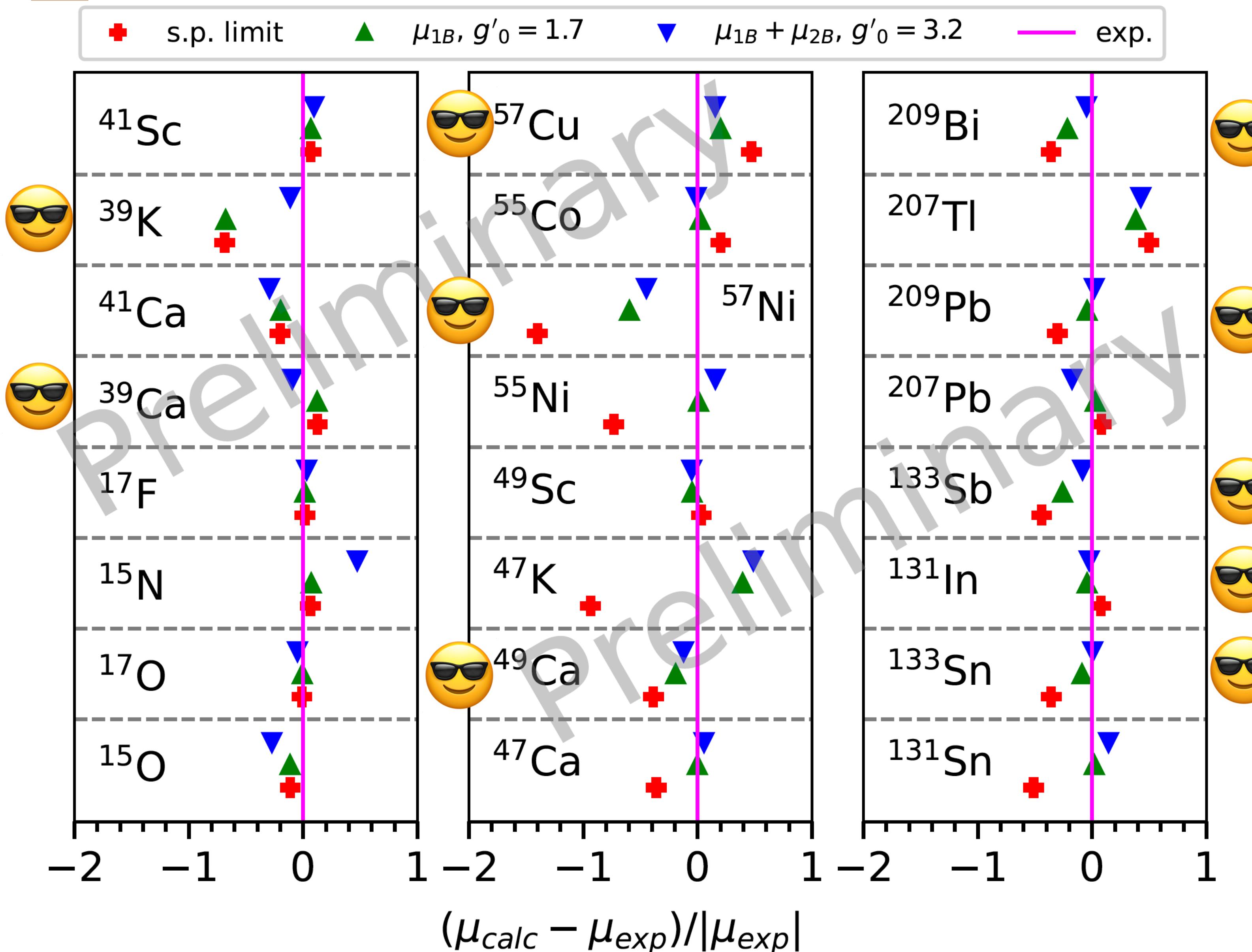
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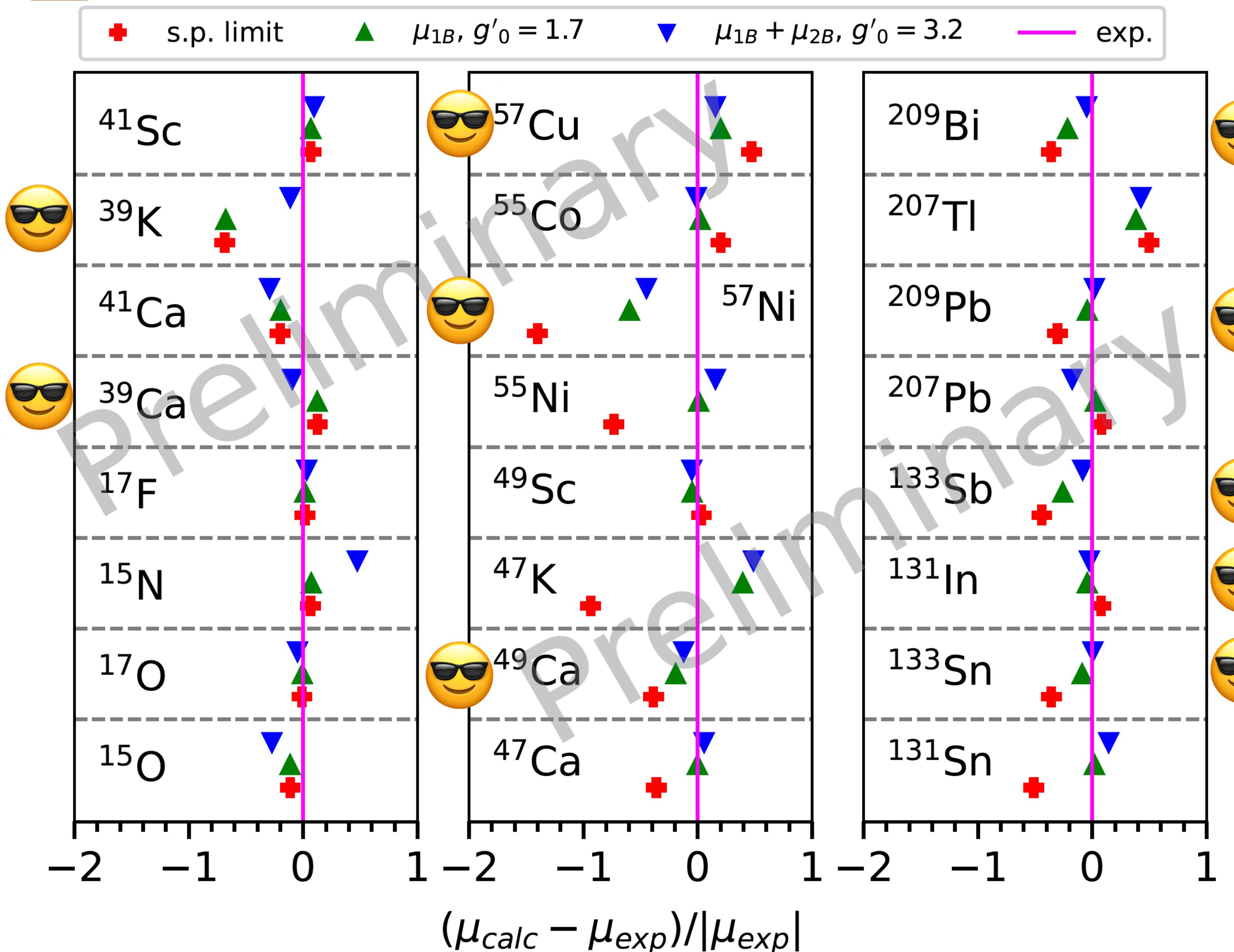
★ In **10/14** cases, the two-body-current corrections **improve / deteriorate** agreement with experimental data.

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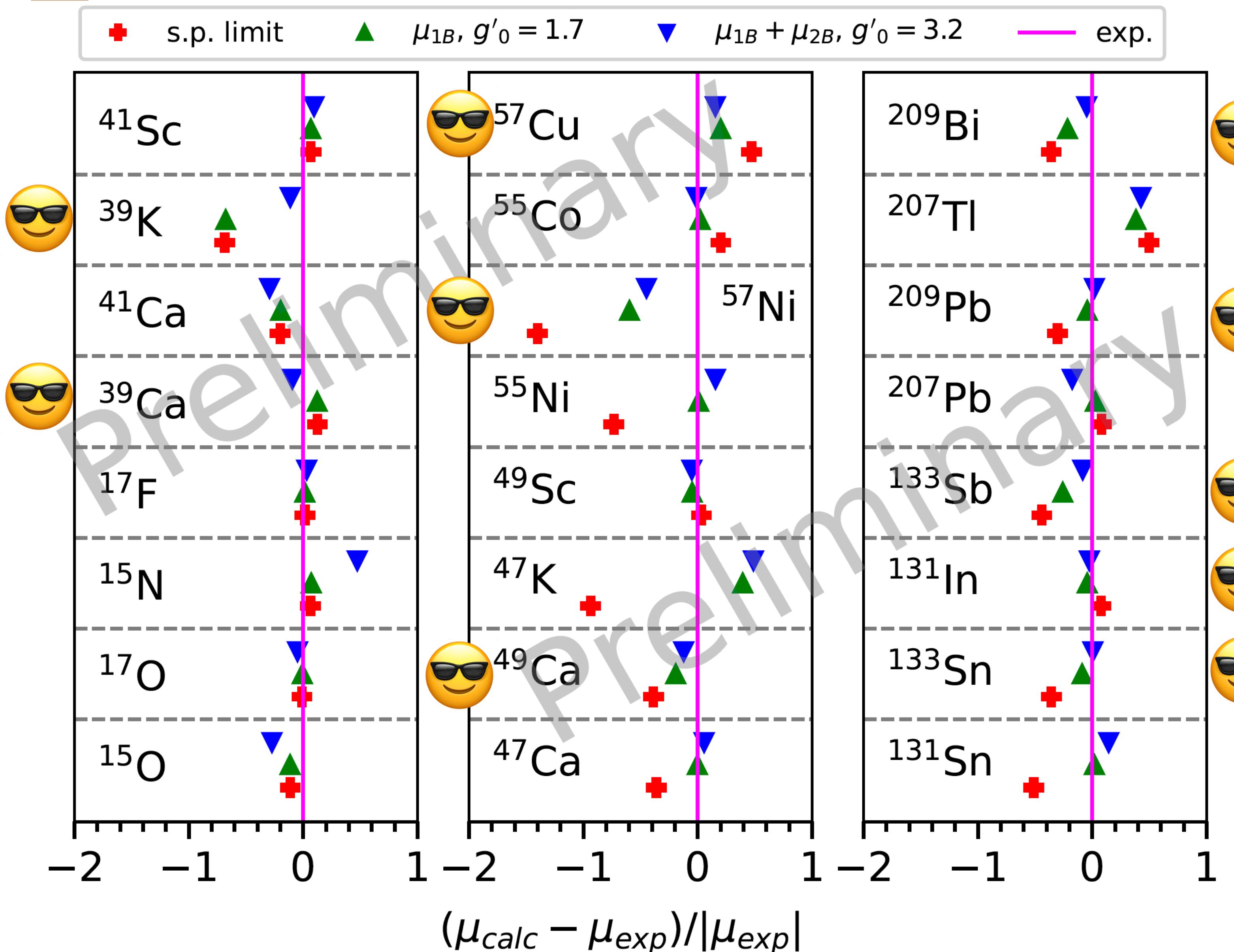
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★ The complete statistical analysis will follow.

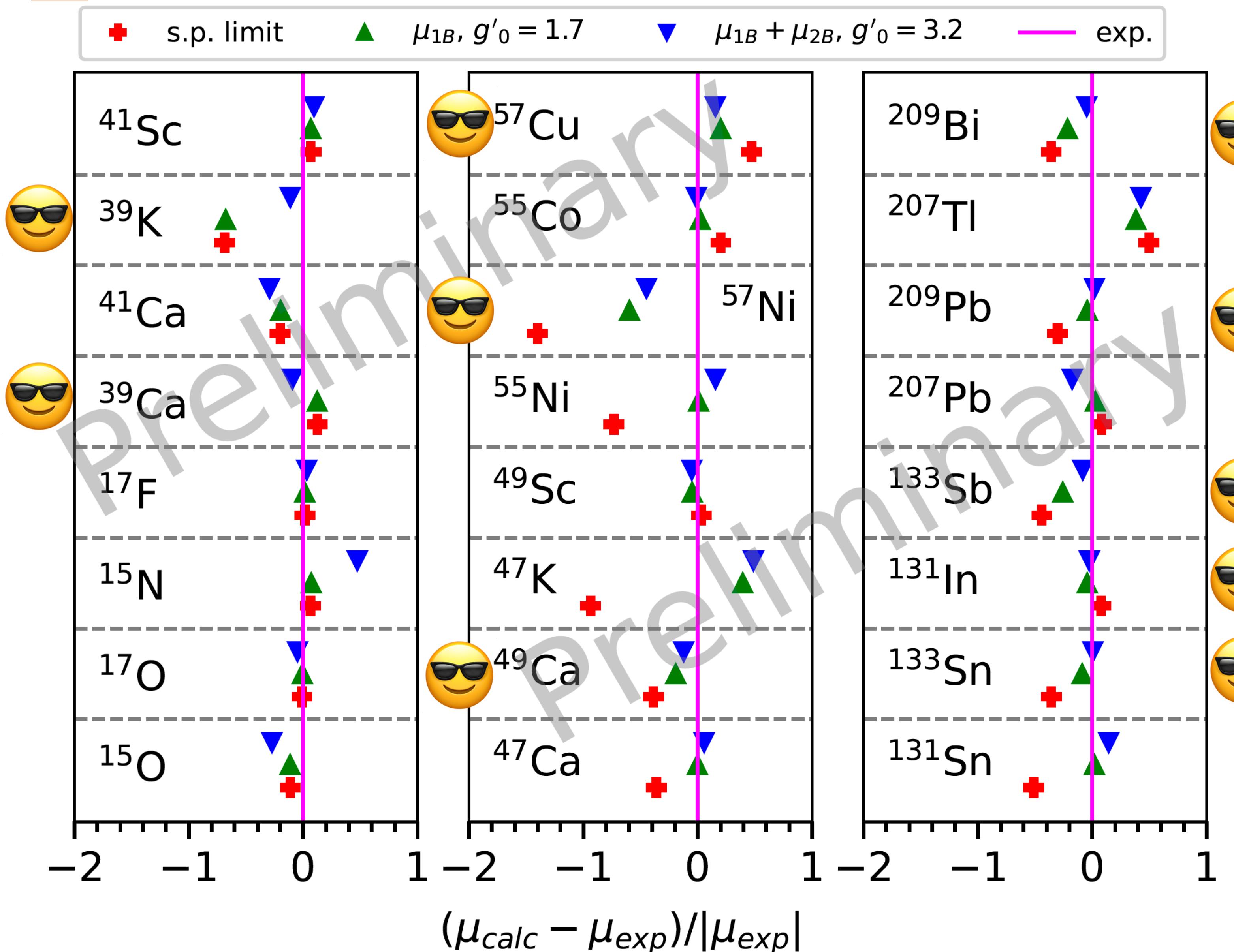
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# Magnetic dipole moments: theory vs. experiment



- ★ In **10/14** cases, the two-body-current corrections **improve/deteriorate** agreement with experimental data.
- ★ The complete statistical analysis will follow.
- ★ Results in deformed nuclei will follow.

# Summary and Outlook

## Summary

- ★ The obtained value of the optimum Landau parameter  $g'_0$  is around 3.2.
- ★ The inclusion of the meson-exchange currents improves the predictions of magnetic dipole moments in  $^{39}\text{K}$ ,  $^{39}\text{Ca}$ ,  $^{57}\text{Cu}$ ,  $^{57}\text{Ni}$ ,  $^{49}\text{Ca}$ ,  $^{209}\text{Bi}$ ,  $^{209}\text{Pb}$ ,  $^{133}\text{Sb}$ ,  $^{131}\text{In}$ , and  $^{133}\text{Sn}$ .

## Outlook

- ★ Bayesian analysis to determine precisely the optimum Landau parameter  $g'_0$ .
- ★ Systematic nuclear DFT calculations of magnetic dipole moments across the nuclear chart with the inclusion of meson-exchange currents.

# Acknowledgements

We acknowledge the support from a **Leverhulme Trust Research Project Grant**. This work was partially supported by the **STFC Grant** Nos. ST/P003885/1 and ST/V001035/1 and by the **Polish National Science Centre** under Contract No. 2018/31/B/ST2/02220. We acknowledge the **CSC-IT Center for Science Ltd., Finland**, for the allocation of computational resources. This project was partly undertaken on the **Viking Cluster**, which is a high performance compute facility provided by the **University of York**. We are grateful for computational support from the **University of York High Performance Computing service, Viking and the Research Computing team**.